

X-ray Photon Correlation Spectroscopy

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Outline

- **Introduction**

- Why (opportunities for mesoscale science) and How (coherence and speckles)
- Speckle fluctuations, dynamics
- Speckle Statistics

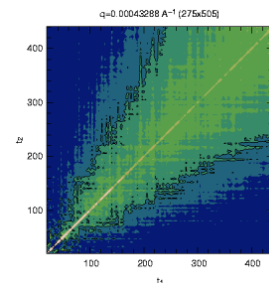
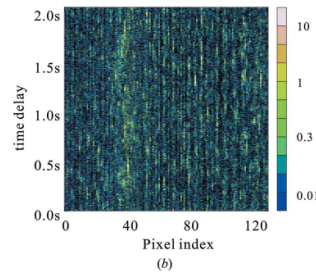
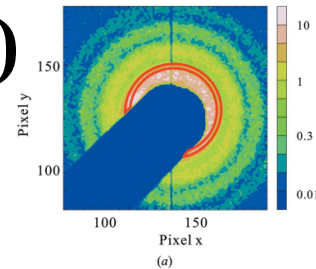
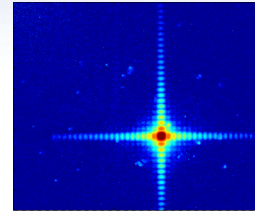
- **X-ray Photon Correlation Spectroscopy (XPCS)**

- Time autocorrelation functions, equilibrium dynamics
- Signal-to-Noise
- Two-time correlation functions, non-equilibrium dynamics
- Higher order correlation functions, dynamical heterogeneities
- X-ray Speckle Visibility Spectroscopy
- A mini user guide to XPCS

- **XPCS examples**

- Dynamics of concentrated hard-sphere suspensions. Is there a colloidal glass transition?
- “Anomalous” relaxations in “jammed” systems

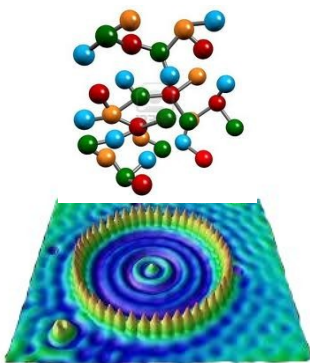
- **Conclusions**



The Next “Big Thing”

- Opportunities for “Mesoscale Science” DOE BESAC report Sept 2012 <http://www.meso2012.com>

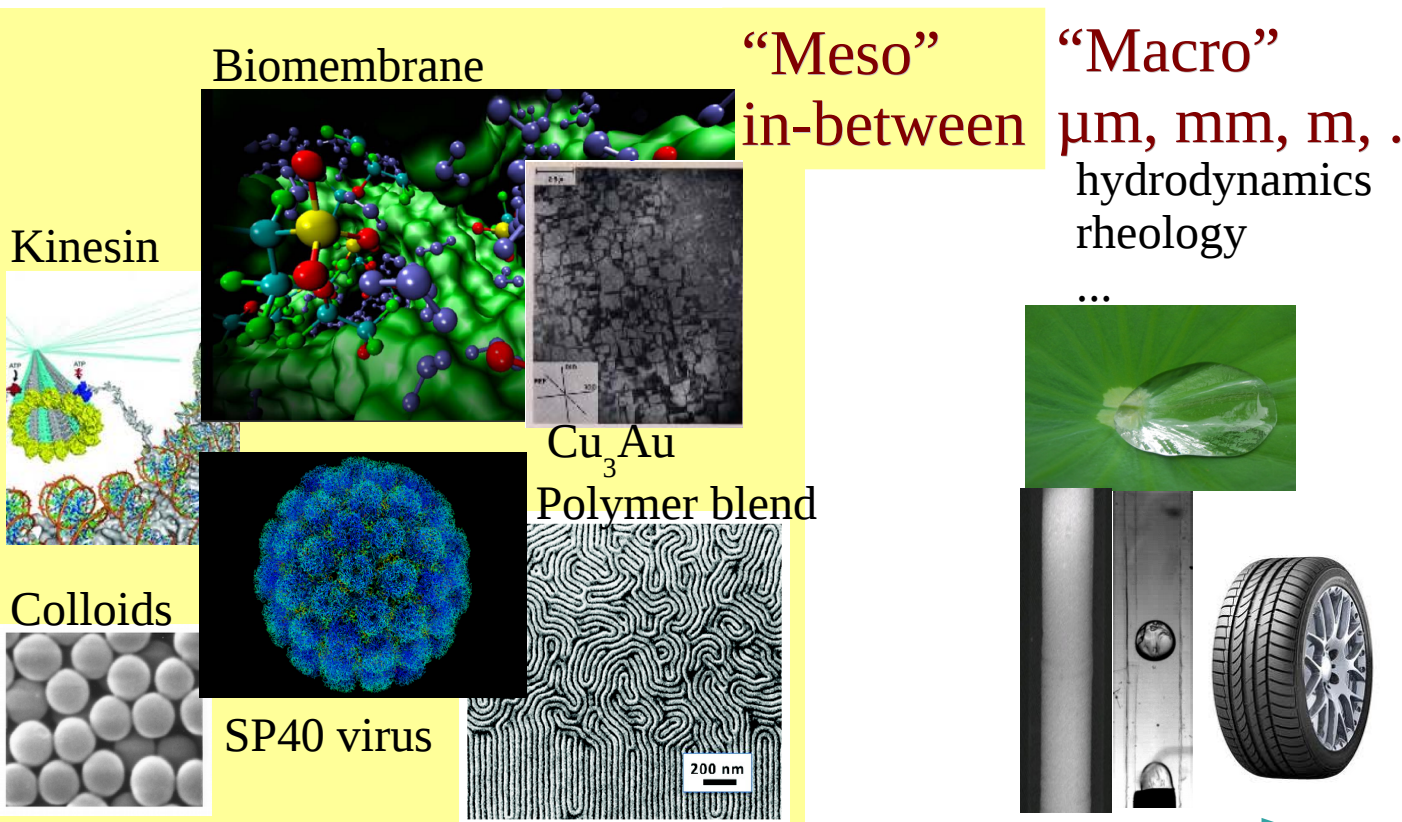
“Nano” nm



Reductionist Science
“Theory of Everything”

$$i\hbar \frac{\partial}{\partial t} \Psi = \hat{H} \Psi$$

Not practical....

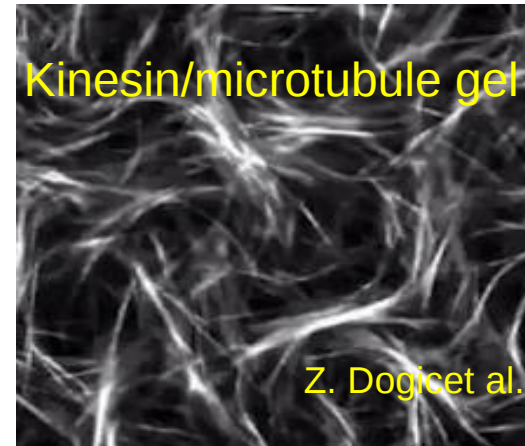
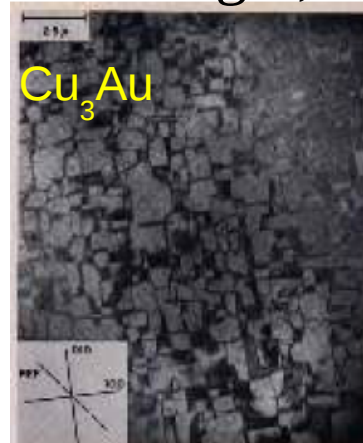
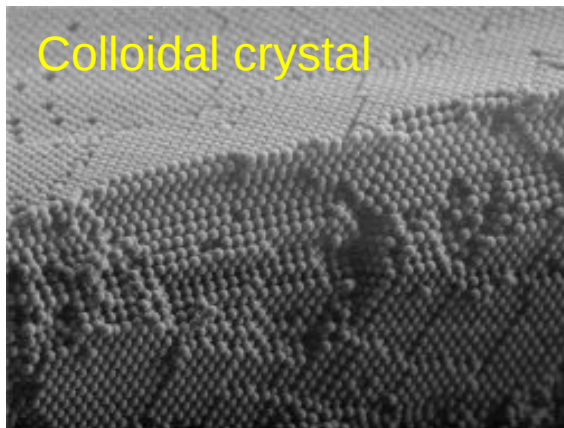


10^{-10} 10^{-9} (1nm) 10^{-8} 10^{-7} length [m] 10^{-6} ($1\mu\text{m}$) 10^{-5} 10^{-4} 10^{-3} (1 mm)

“More is Different”

P.W. Anderson, *Science* **177**, 393 (1972)

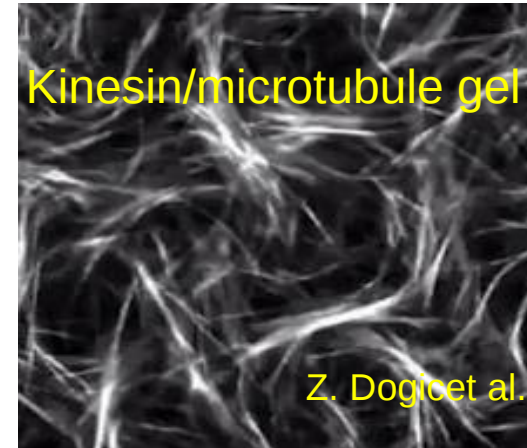
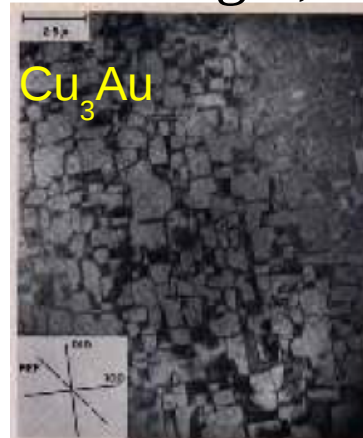
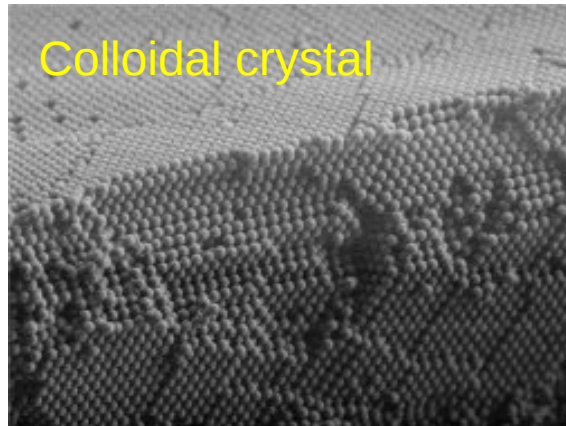
- Most *macroscopic properties of complex disordered materials emerge at the mesoscale* (nm to μm):
 - Mesoscale structure: defects, grain size, macromolecule shape/size, entanglement length, ...



“More is Different”

P.W. Anderson, *Science* 177, 393 (1972)

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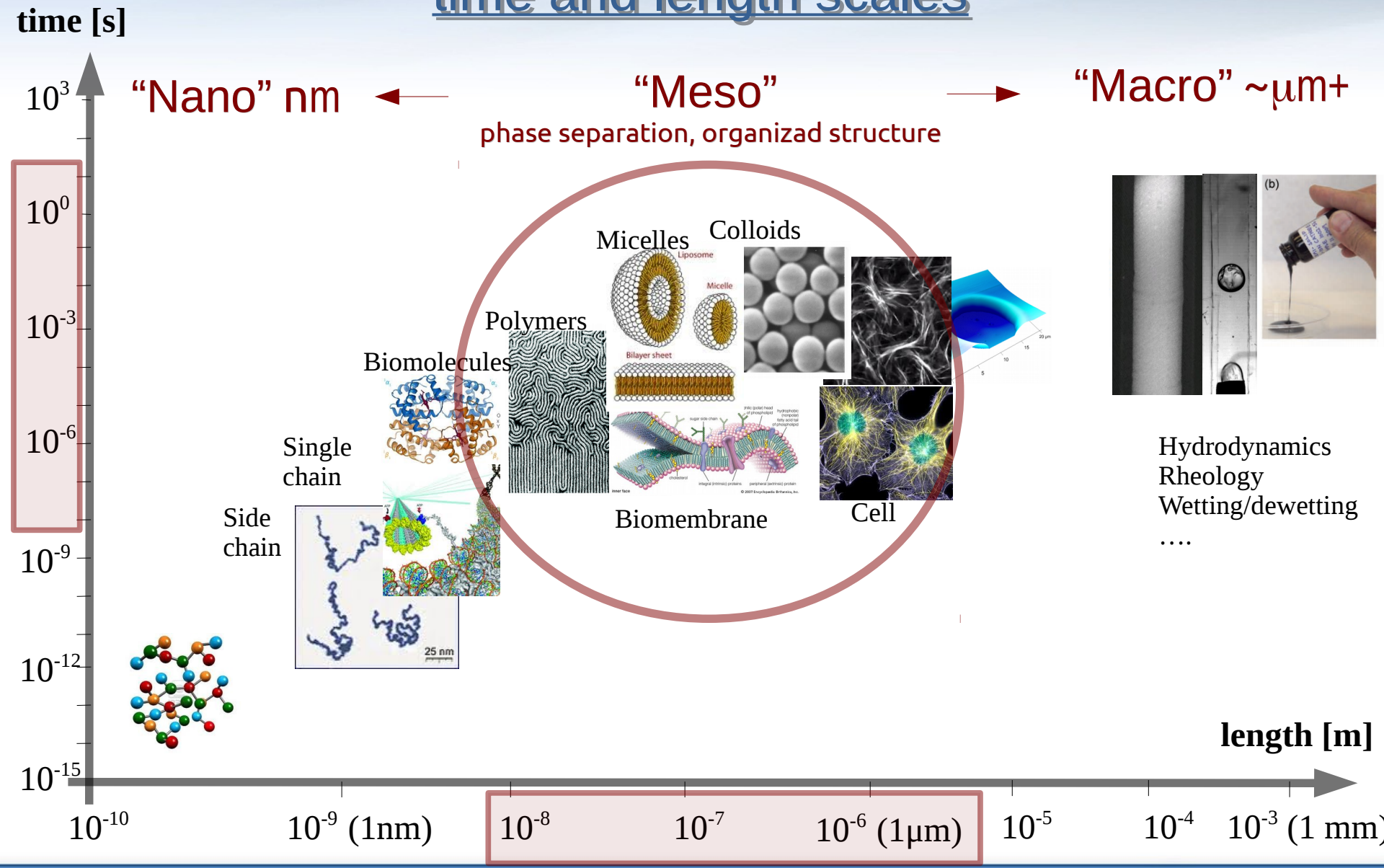
But things are not static !

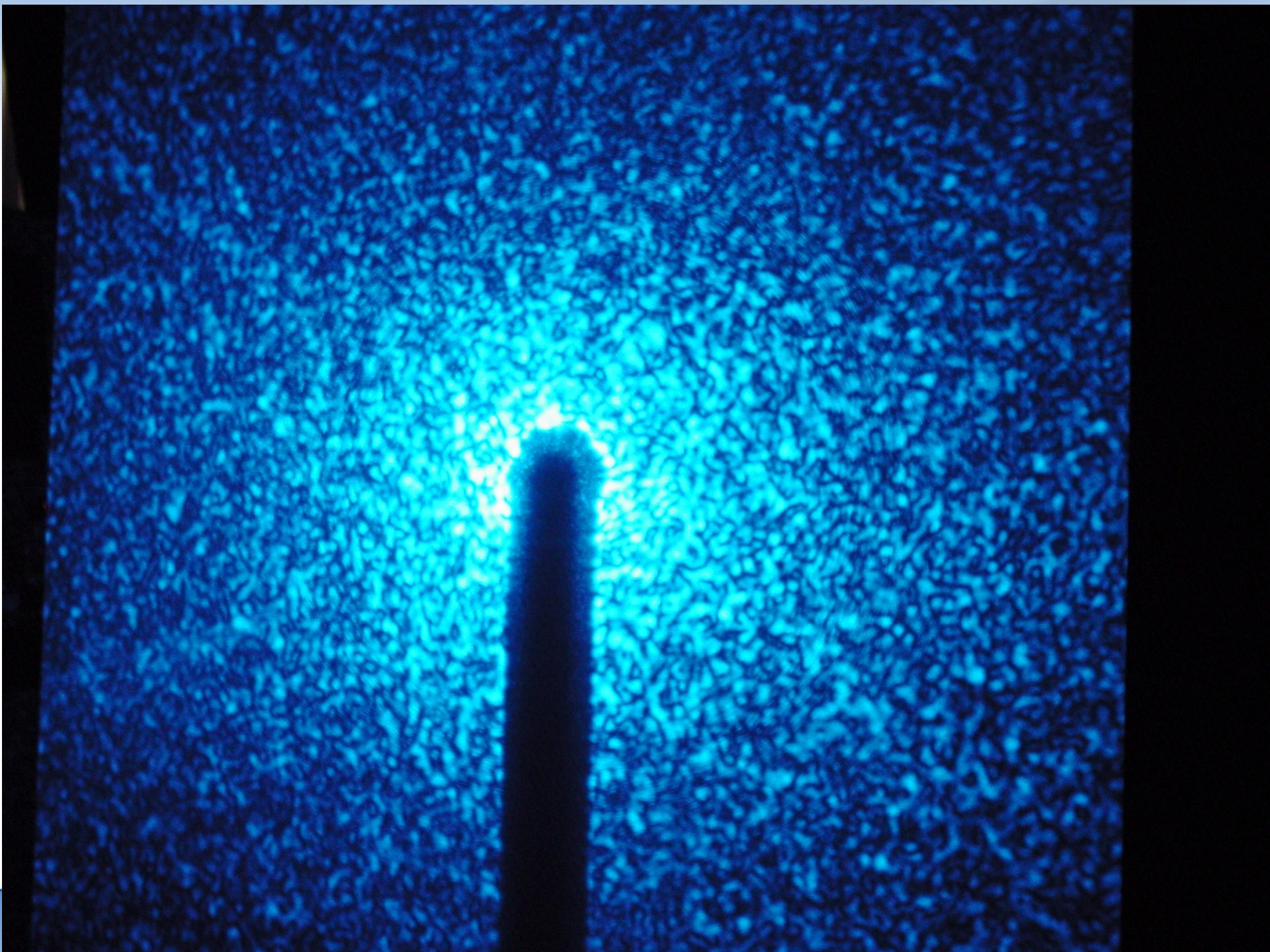
- Mesoscale Dynamics

Z. Dogic (Brandeis Univ.)
Dynamics of bundled active networks



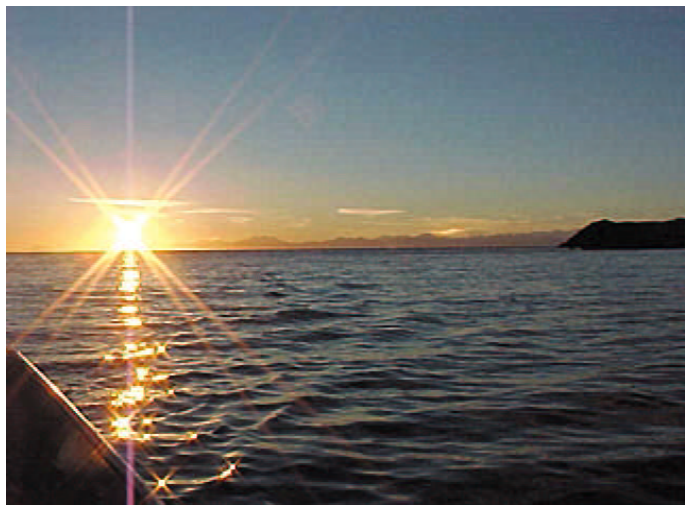
Dynamics of Materials (soft- and bio-): time and length scales





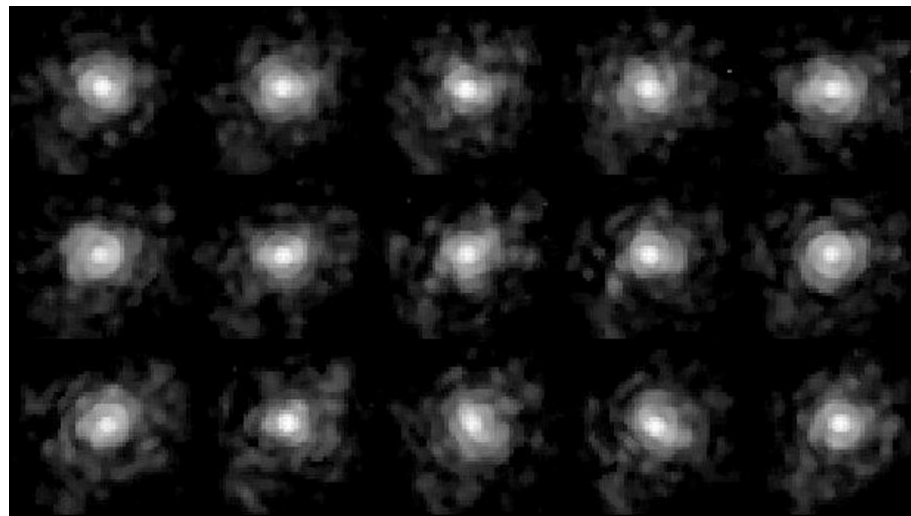
Speckles

Sunset in Alaska



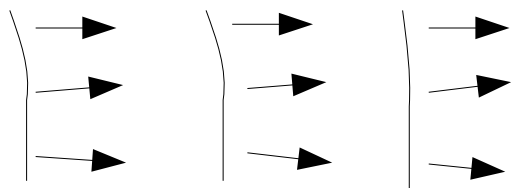
From S. Mochrie

Images of a Stars in a Telescope



J. Codona

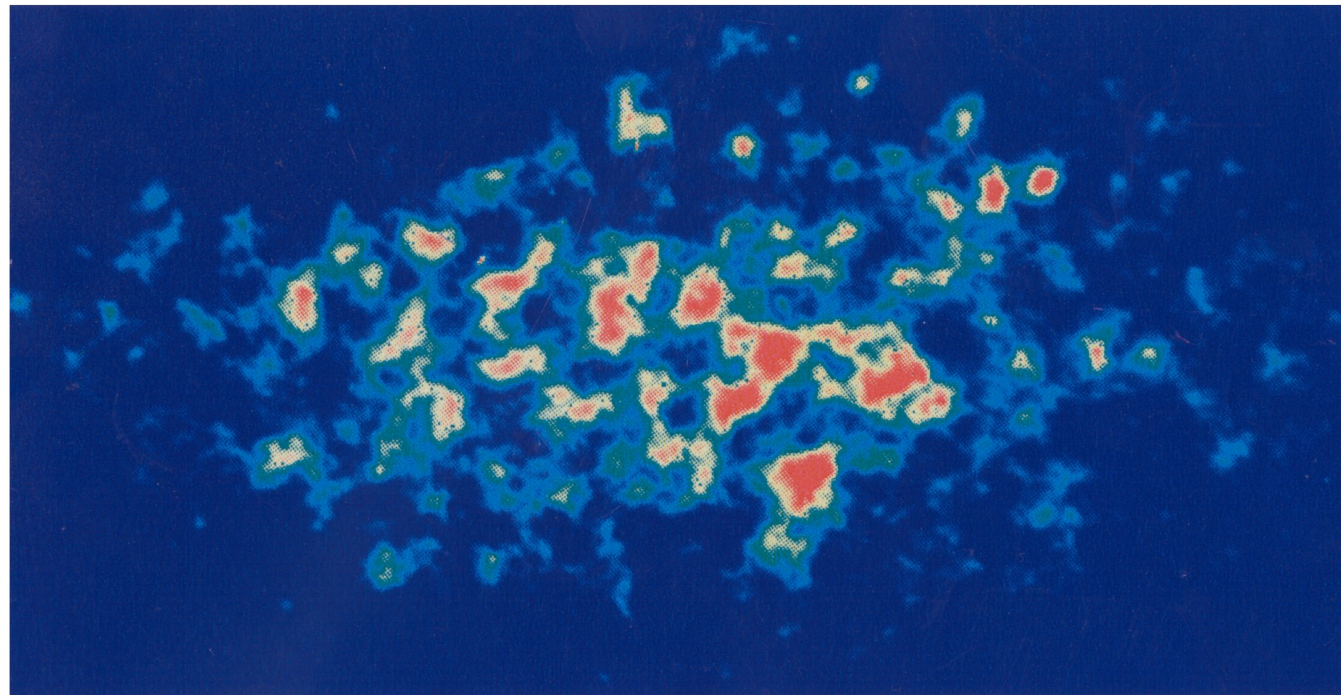
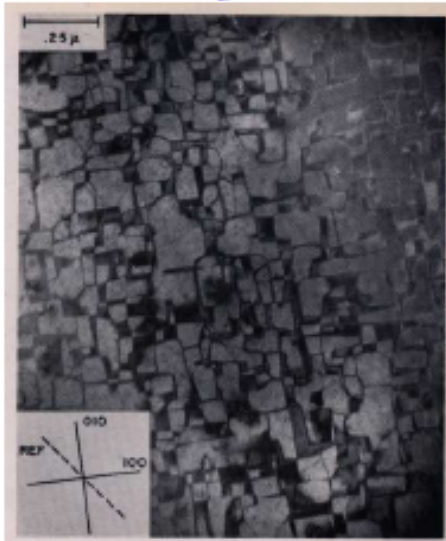
- Stars (far away) = nearly coherent “point-like” sources + fluctuations



Speckles with (partially) coherent X-rays

Speckles from Cu_3Au

Cu_3Au

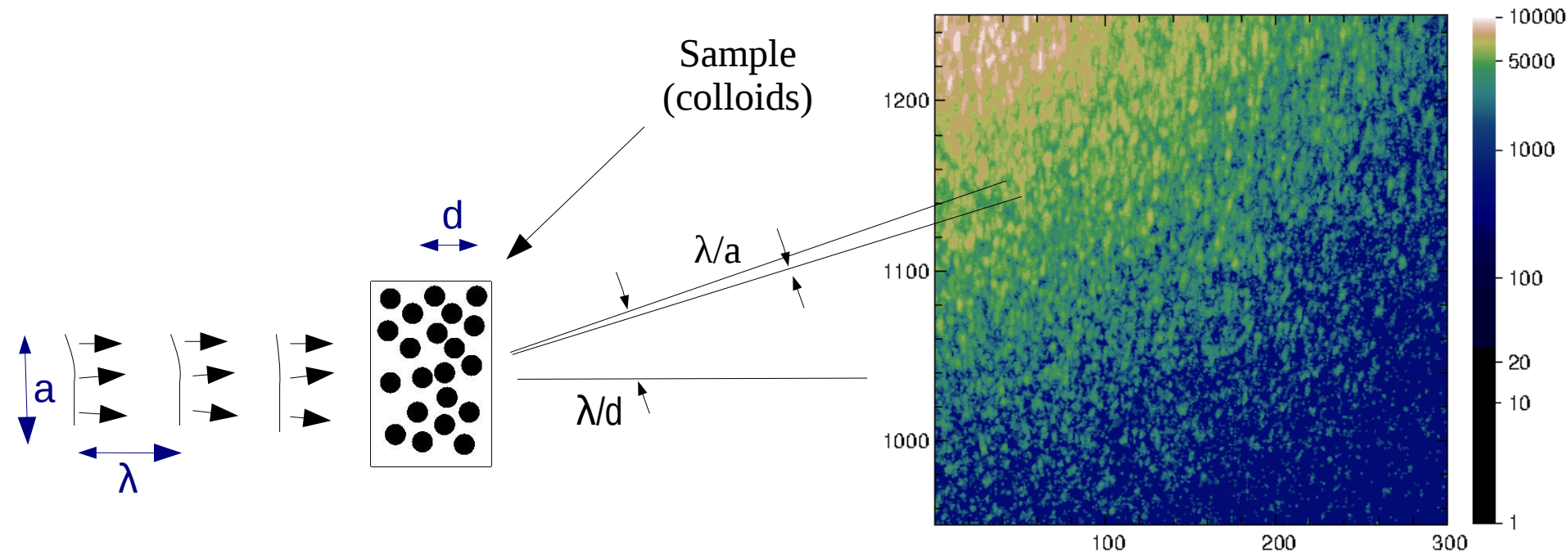


Recorded at X25, NSLS on Kodak film

M. Sutton, *et al.* *Nature* **352**, 608 (1991)

Speckles with (partially) coherent X-rays

Speckles from colloidal suspensions

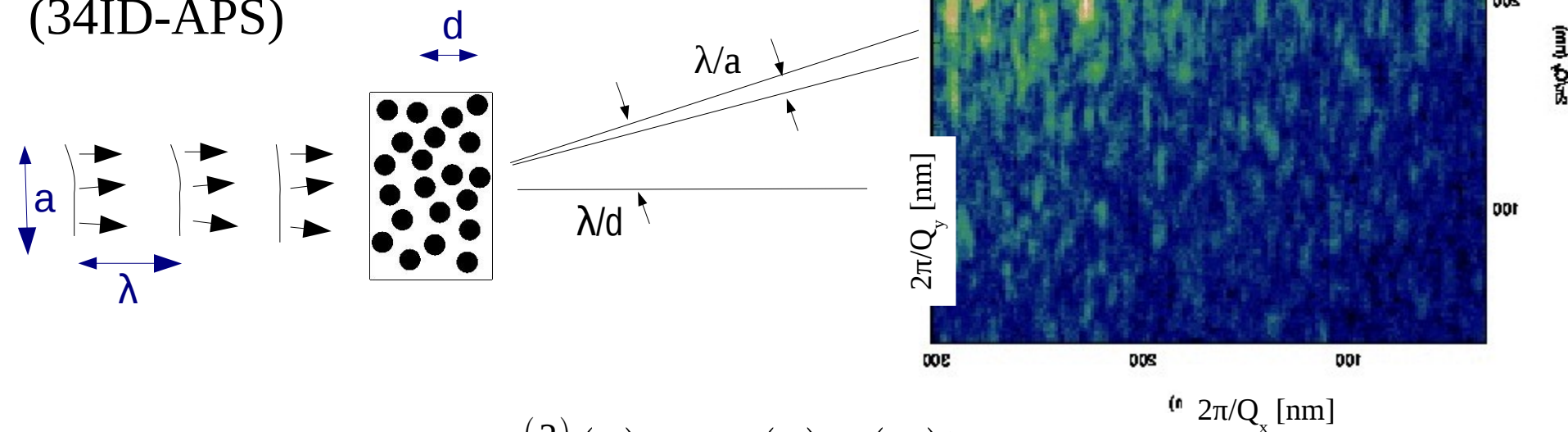


Measured at 34ID with a CCD detector

Speckle Fluctuations & Dynamics

- At high brightness light sources (APS, ESRF, Petra-III, NSLS-II ...) it is possible to measure dynamics by recording “speckle movies”

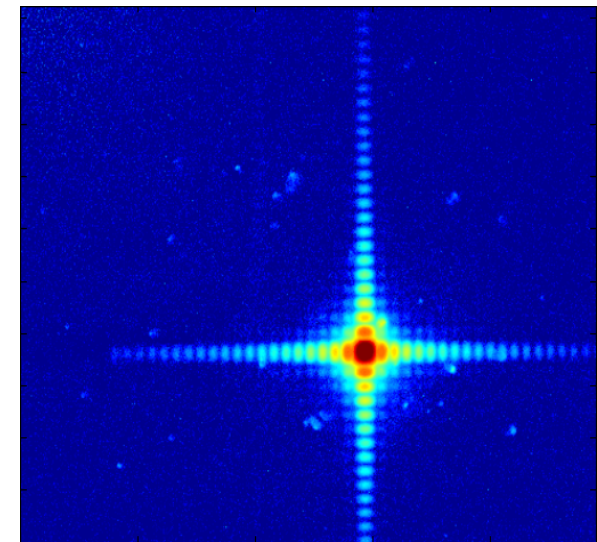
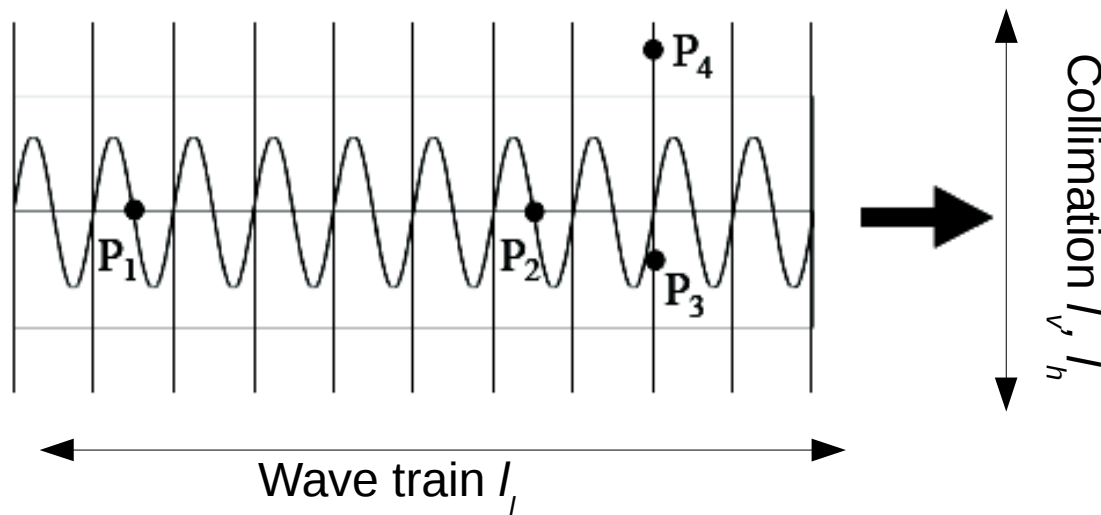
Partially Coherent X-rays (34ID-APS)



$$g^{(2)}(t) \propto \langle I(t) I(0) \rangle$$

Mini-introduction to coherence

- Coherence = ability to create interference fringes w. good contrast
 - i.e. exists within a region where the phase difference between any pair of points is well defined and constant in time
 - Transverse coherence: $\Delta\Phi(P3:P4)$
 - Longitudinal(temporal) coherence: $\Delta\Phi(P1:P2)$



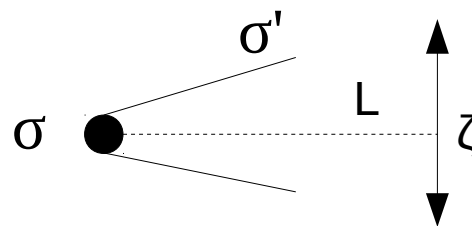
Malcolm Howells, Lecture Notes, ESRF 2007

L. Wiegart, CHX, NSLS-II

Transverse coherence

- Ideal *coherent* (Gaussian) source:
 - a source cannot be arbitrarily small and arbitrarily well collimated at the same time (diffraction limit)

$$\sigma \cdot \sigma' \simeq \frac{\lambda}{4\pi}$$



- A transverse coherence length (@ distance L from the source) can then be defined as:

$$l_{h,v} = \frac{\lambda L}{4\pi \sigma_{h,v}}$$

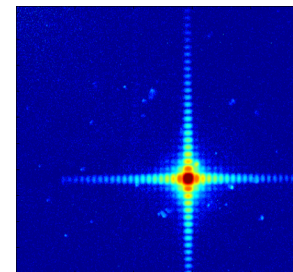
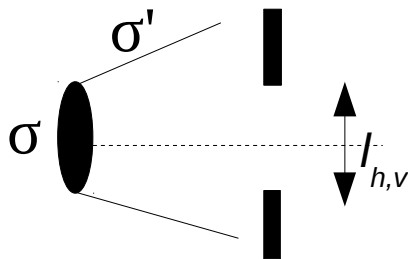
Transverse coherence

- Real Source:

- The degree of coherence is determined by the phase space volume $\sigma\sigma'$; “Heisenberg's inequality”:

$$\sigma \cdot \sigma' \geq \frac{\lambda}{4\pi}$$

- “Liouville's theorem”: the phase space is conserved by propagation, (ideal) crystal optics, (ideal) focusing, etc.
- To obtain a more coherent beam (at the expense of flux!), the phase space can be limited/reduced by collimation (a set of slits)



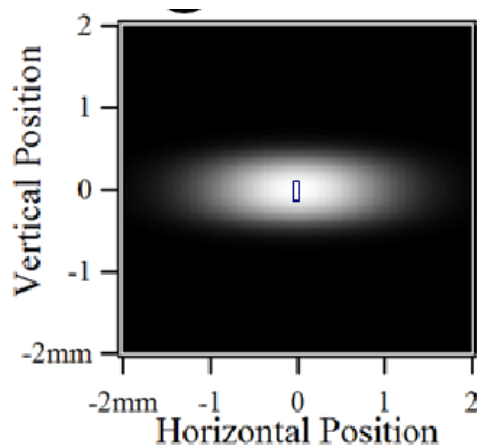
Coherence of (NSLS-II) Sychrotron Sources

- Real Source:

- Number of coherent modes:

$$\sigma \cdot \sigma' = N \frac{\lambda}{4\pi}, N \geq 1$$

- E.g. IVU20 undulator source at CHX, NSLS-II



E (keV)	6	8	10	12	16
σ_h (μm)	34.3	34.2	34.1	34.2	34.2
σ_h' (μrad)	18.3	18.3	18.0	18.2	18.2
σ_v (μm)	8.8	8.0	7.5	7.6	7.4
σ_h' (μrad)	8.5	8.2	7.7	8.1	8.0
M_h	38.2	50.7	62.2	75.7	94.6
M_v	4.5	5.3	5.8	7.5	9.0

Longitudinal coherence

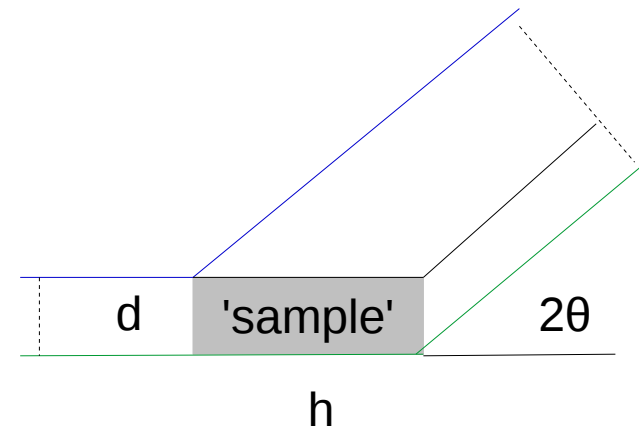
- Longitudinal (temporal) coherence:

$$\frac{\delta \lambda}{\lambda} \approx \frac{1}{N}, l_l = \lambda N \quad l_l \approx \frac{\lambda^2}{\delta \lambda}$$

- Experimental requirement:
max optical path diff. $< l_l$

- In a transmission geometry
 - Sample thickness h , beam size d

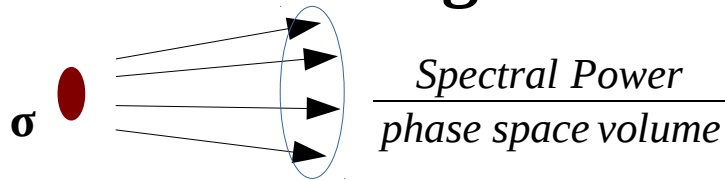
$$h \sin^2(2\theta) + d \sin(\theta) \leq l_l$$



A. Madsen, A. Fluerasu, B. Ruta, Structural Dynamics of Materials probed by X-ray Photon Correlation Spectroscopy, Springer, 2014

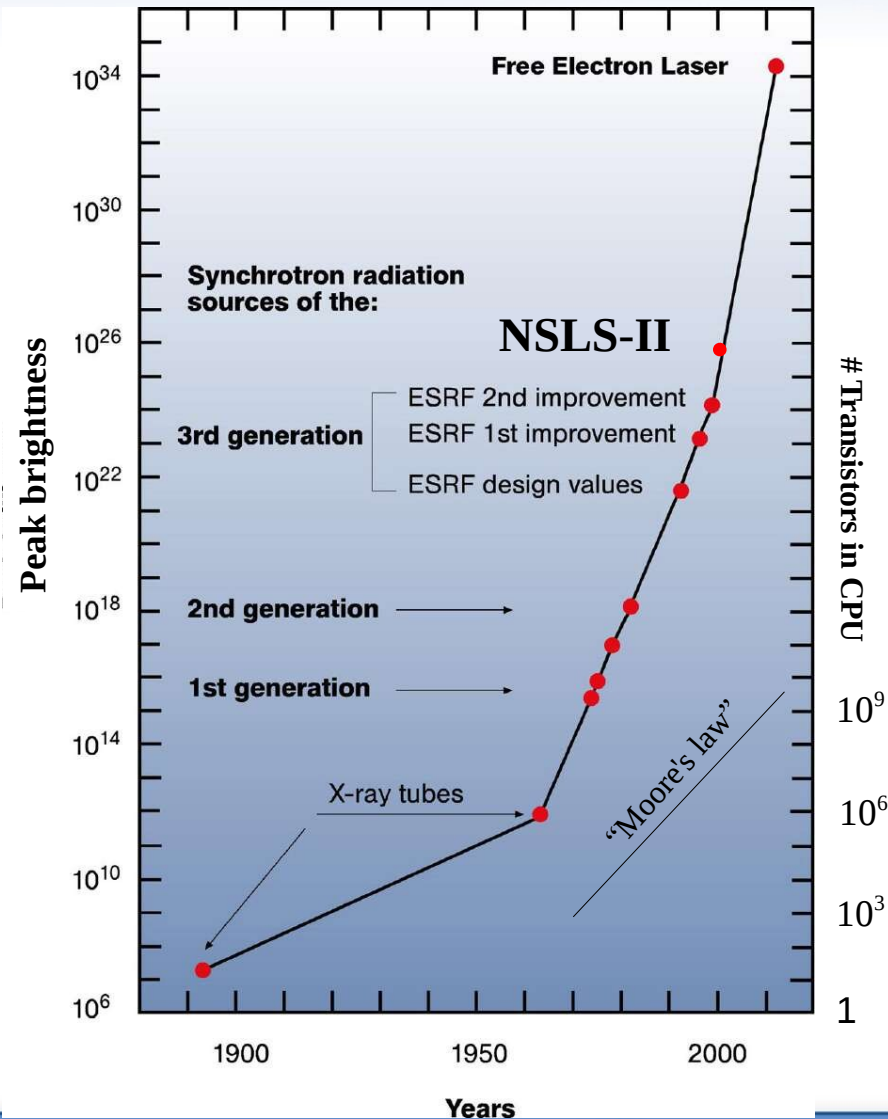
Synchrotron Source Brightness

- Key for XPCS: **Brightness**



Brightness=Coherence
increased faster than Moore's law!!

- Coherent Flux $I \propto B \lambda^2$
- CHX, NSLS-II (~10 keV)
 $B \sim 10^{21} \text{ ph/s/\%bw/mm}^2/\text{mrad}^2$
 $I \sim 10^{11} \text{ ph/s}$



Correlation Functions

- Coherence → measures dynamics

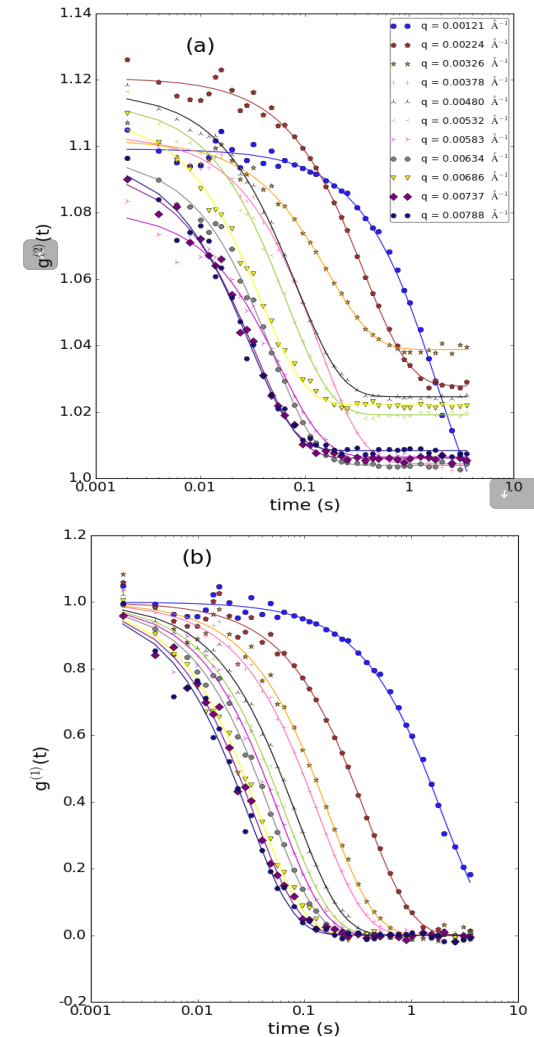
$$\langle I(q,t)I(q,t+\delta t) \rangle = \langle I(q) \rangle^2 + \beta(q)(\dots) |S(q,t)|^2$$

- Intensity autocorrelation function, dynamic structure factor & Siegert relationship:

$$g^{(2)}(q,t) = \frac{\langle I(q,t)I(q,t+\delta t) \rangle}{\langle I(q) \rangle^2} = 1 + \beta(q) \left| \frac{S(q,t)}{S(q,0)} \right|^2$$

- Intermediate Scattering Function

$$g^{(1)}(q,t) = \left| \frac{S(q,t)}{S(q,0)} \right| \propto \iint \rho_n(q) \rho_m(q) \exp(iq[r_n(0) - r_m(t)])$$



Correlation Functions

- Signal-to-noise (of $g^{(2)}$) – it's complicated!!

$$R_{sn} = K (T \tau \Omega_x \Omega_z)^{1/2} \Sigma W \exp(-W \Lambda) \tilde{B} (\Delta E / E) r_{snx} r_{snz}$$

K = detector efficiency

T = total experiment duration

τ = accumulation time

Ω = angle subtended by Q of interest

Σ = scattering cross section per unit volume

W = sample thickness

Λ = 1/attenuation length

B = source brilliance

$\Delta E / E$ = normalized energy spread

r = factor depending on source size, pixel size, and slit size

- SNR $\sim B \tau^{1/2} \dots$
- Need an area det
- \sim small pixels
- fast frame rates

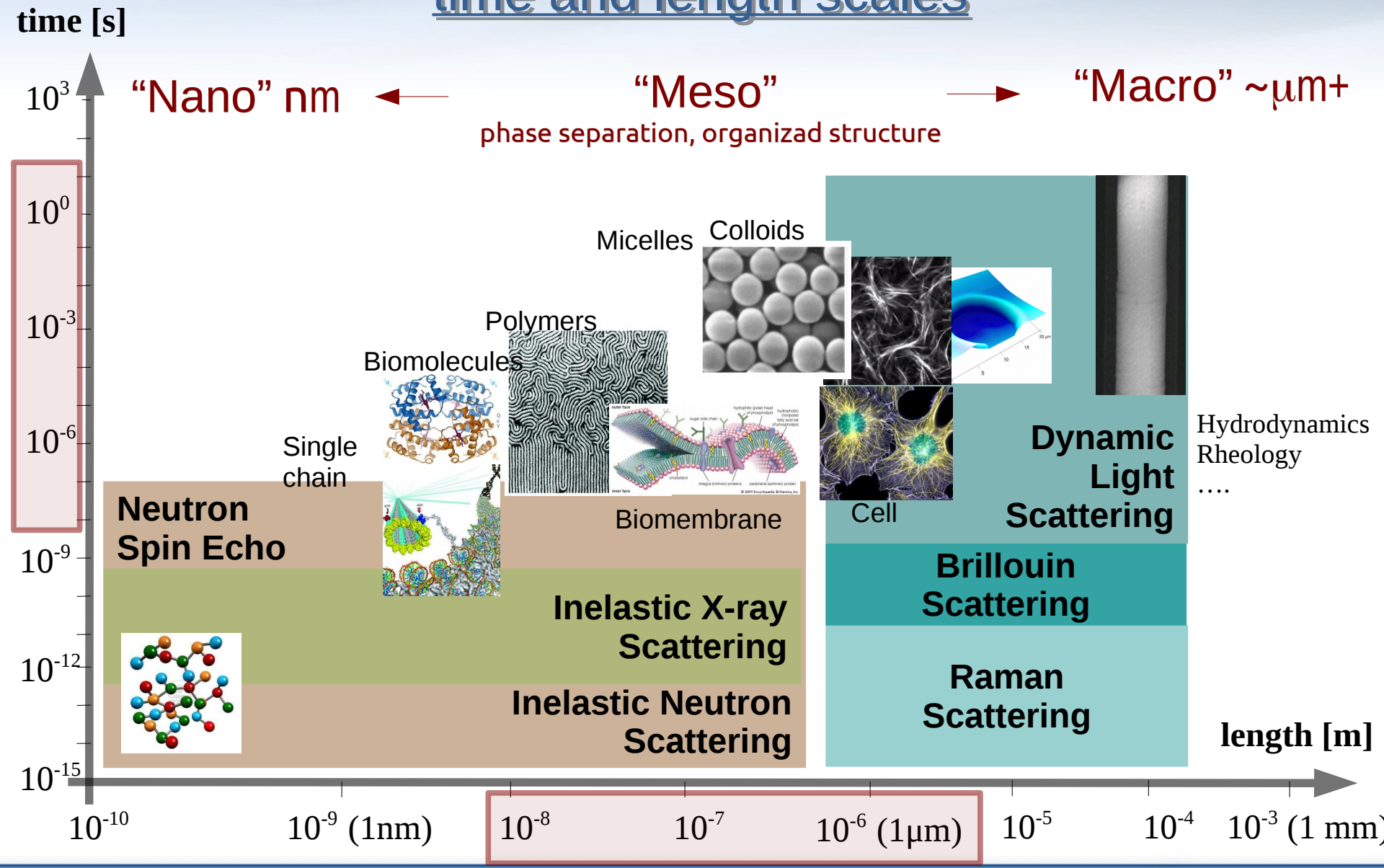


Eiger 1M detector
(Dectris)

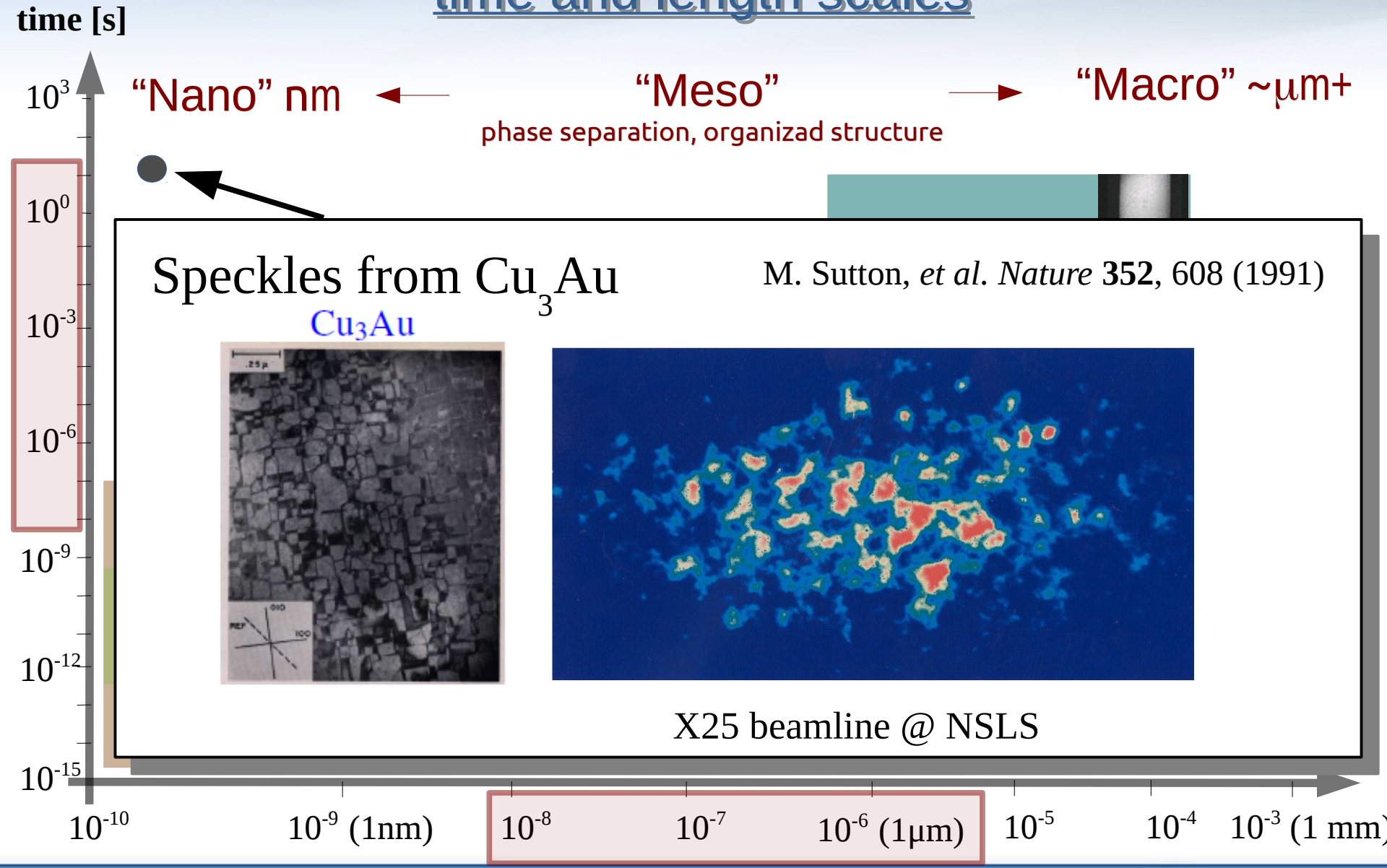
Lumma *et al.* *Rev. Sci. Instrum.* 71, 3274 (2000)

Jackeman *et al.* *J. Phys. A*, 5, 517 (1971)

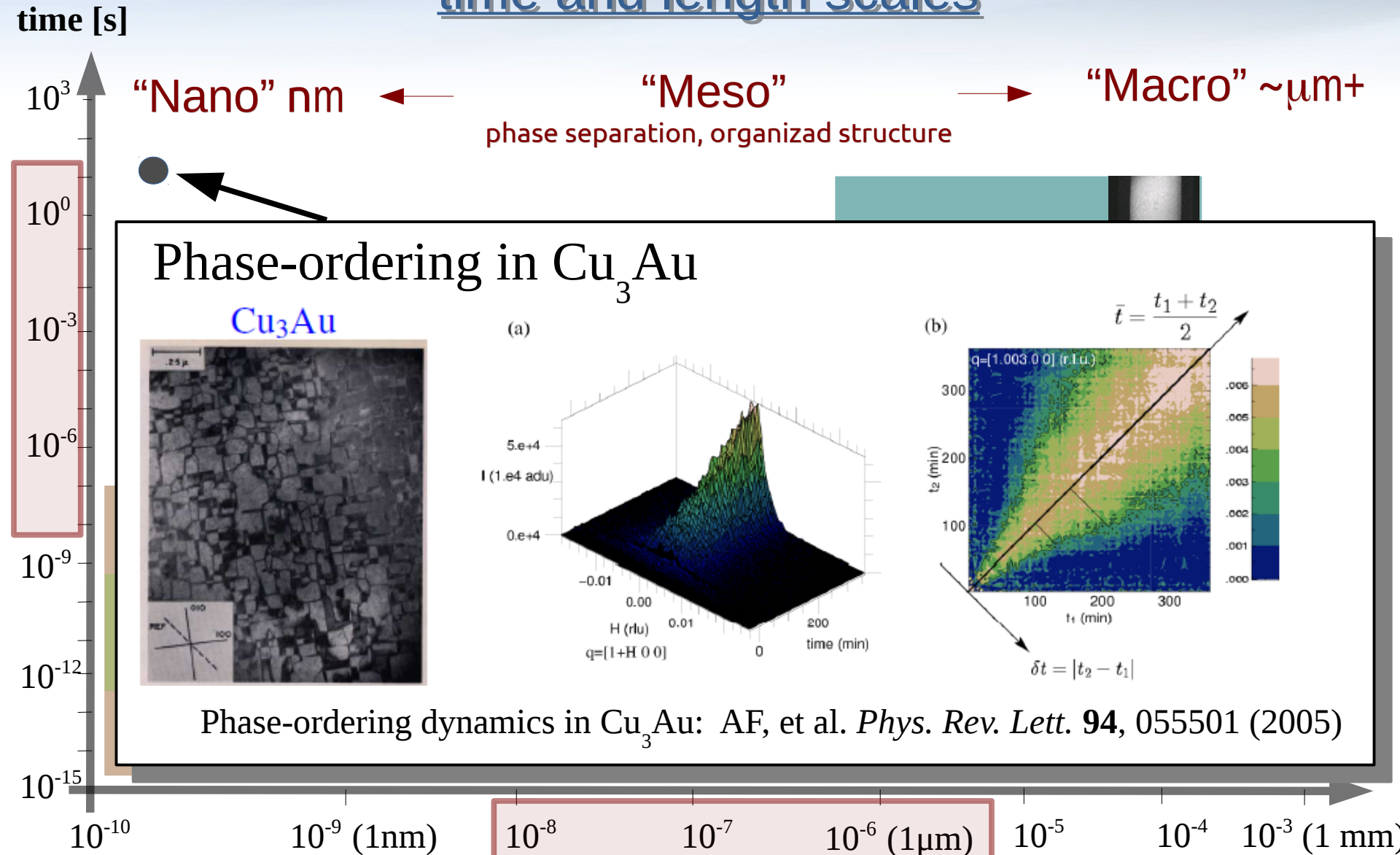
Dynamics of Materials (soft- and bio-): time and length scales



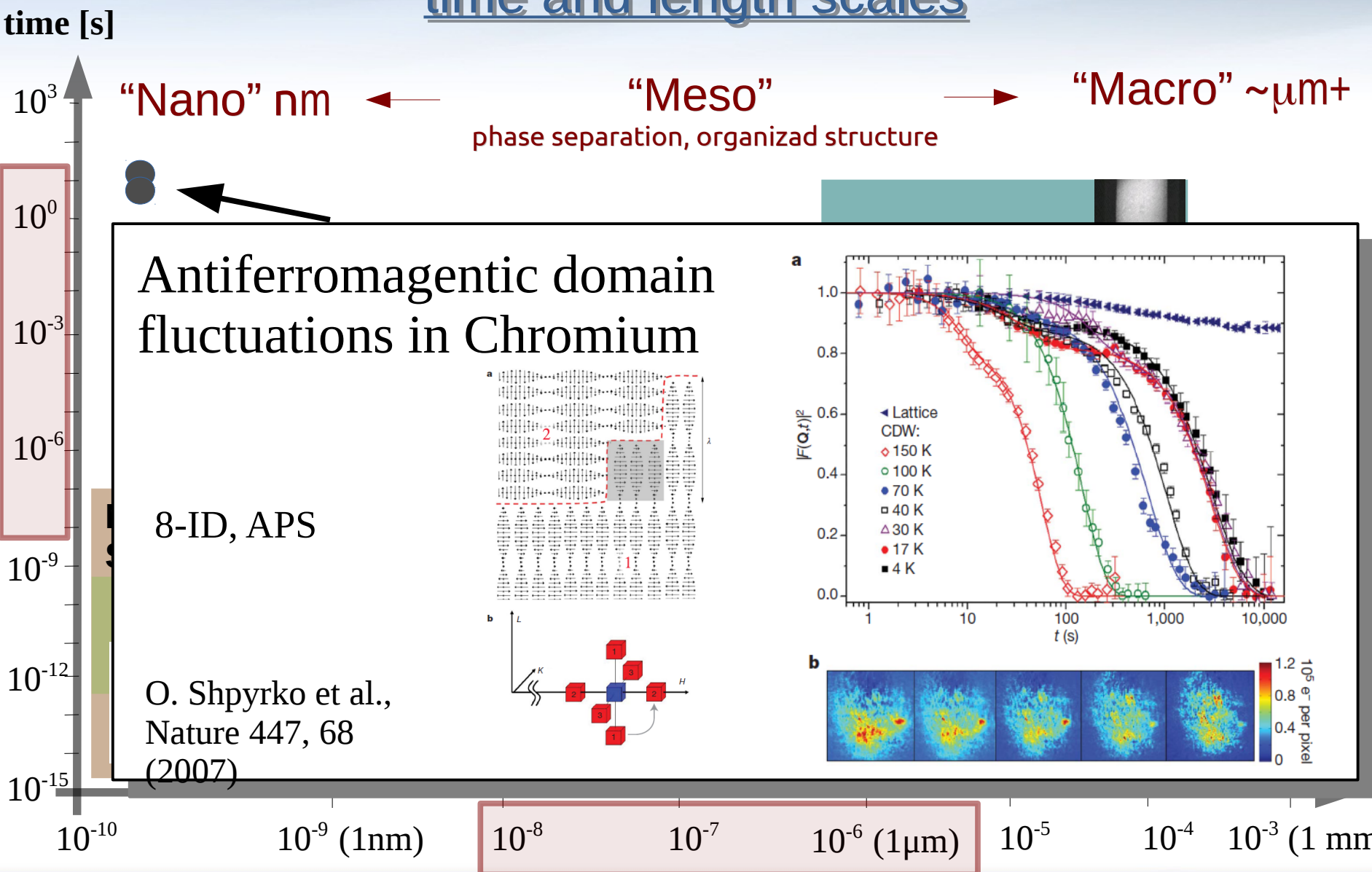
Dynamics of Materials (soft- and bio-): time and length scales



Dynamics of Materials (soft- and bio-): time and length scales



Dynamics of Materials (soft- and bio-): time and length scales



Dynamics of Materials (soft- and bio-): time and length scales

time [s]

“Nano” nm

“Meso”

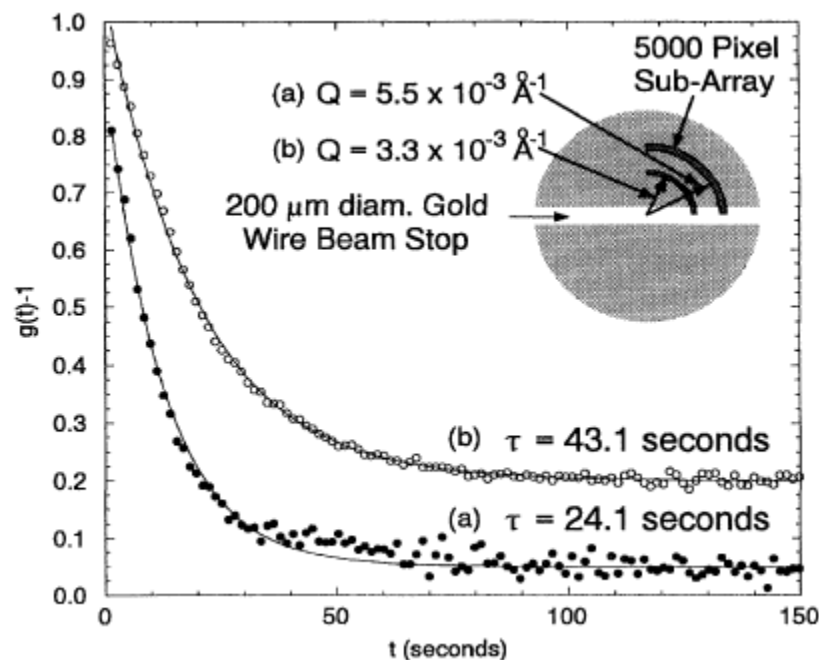
“Macro” $\sim \mu\text{m}+$

phase separation, organized structure

Brownian motion:
Au colloids in glycerol

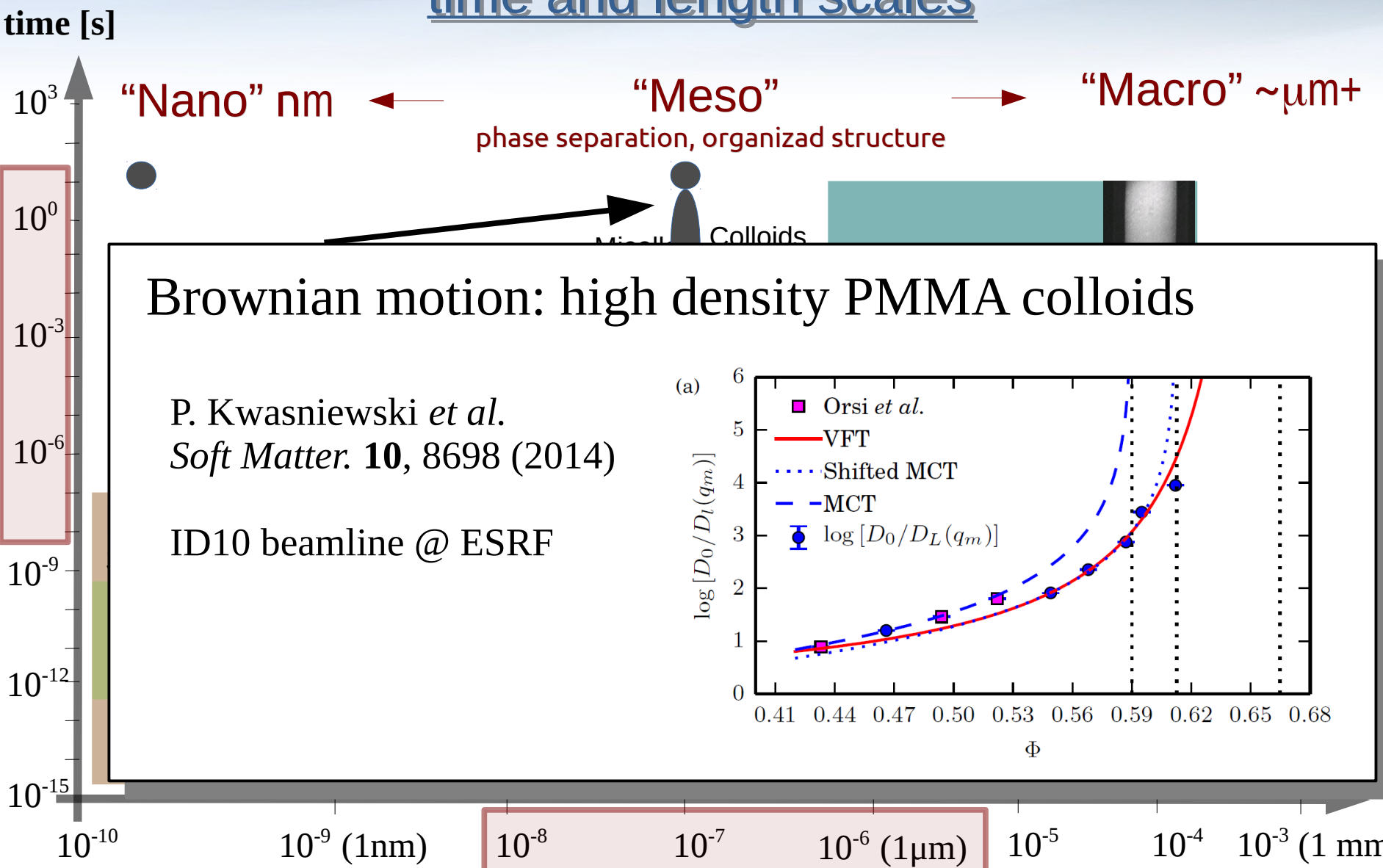
S.B. Dierker, R. Pindak *et al.*
Phys. Rev. Lett. 75, 449 (1995)

X25 beamline @ NSLS

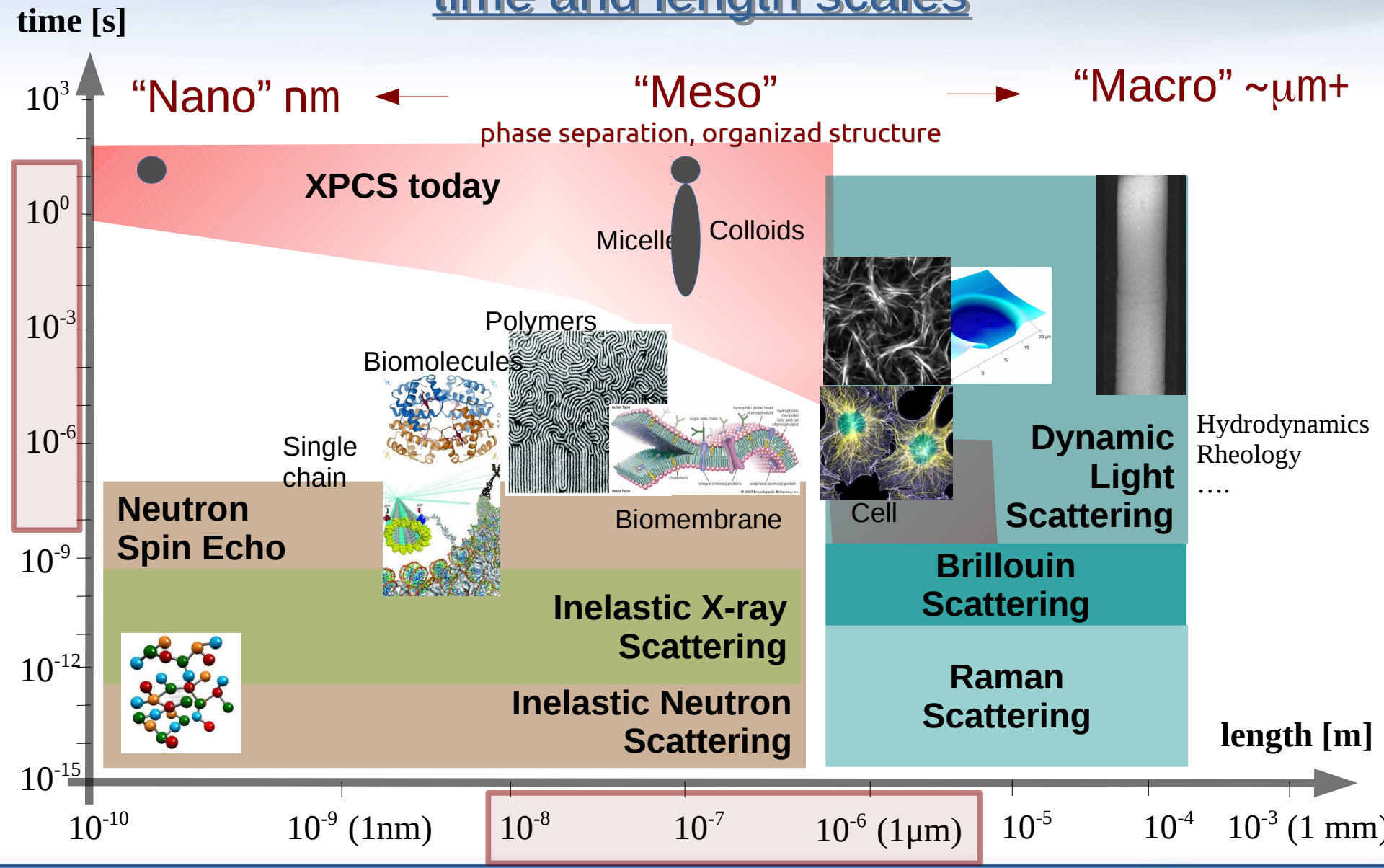


10^{-10} 10^{-9} (1nm) 10^{-8} 10^{-7} 10^{-6} (1 μm) 10^{-5} 10^{-4} 10^{-3} (1 mm)

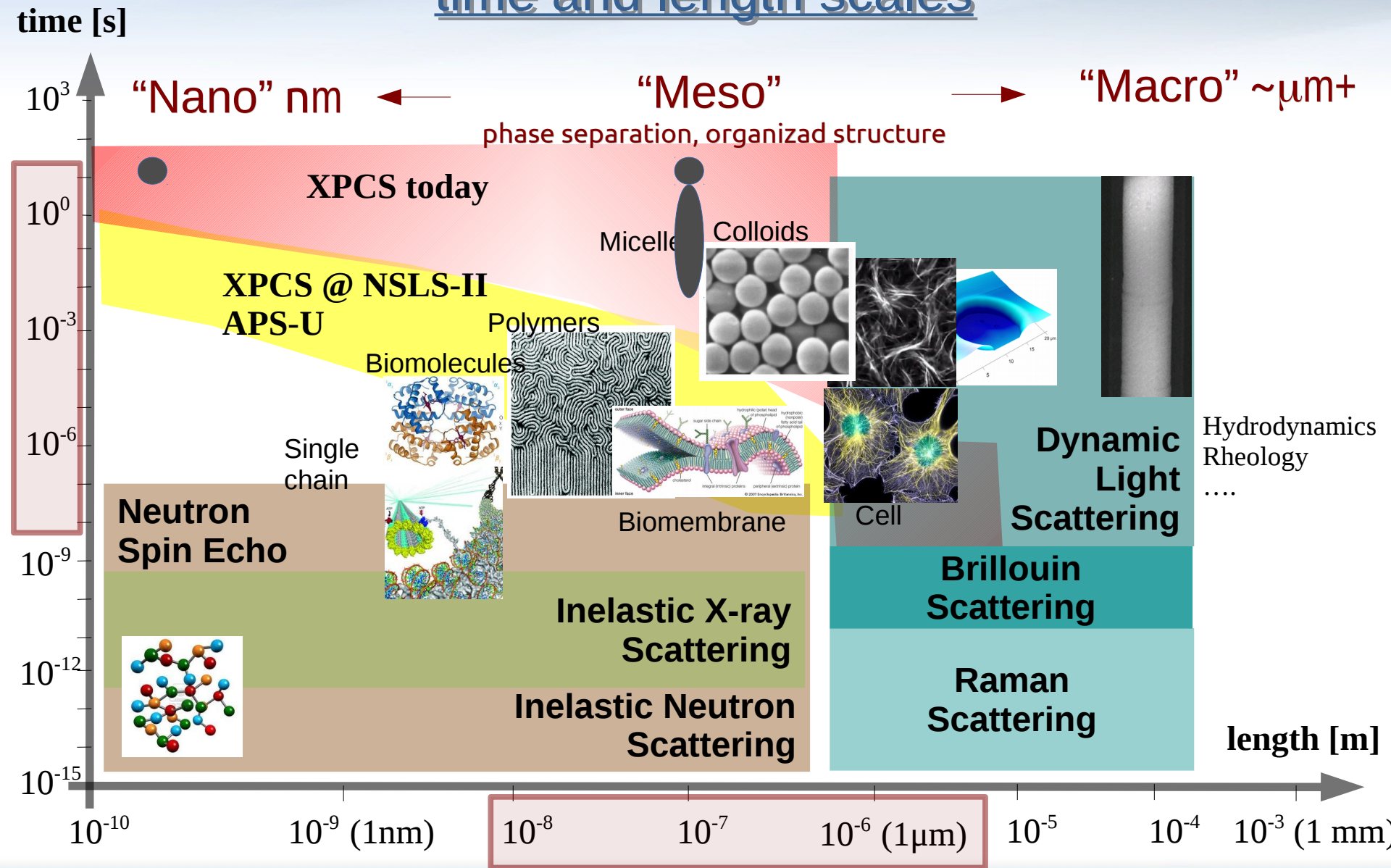
Dynamics of Materials (soft- and bio-): time and length scales



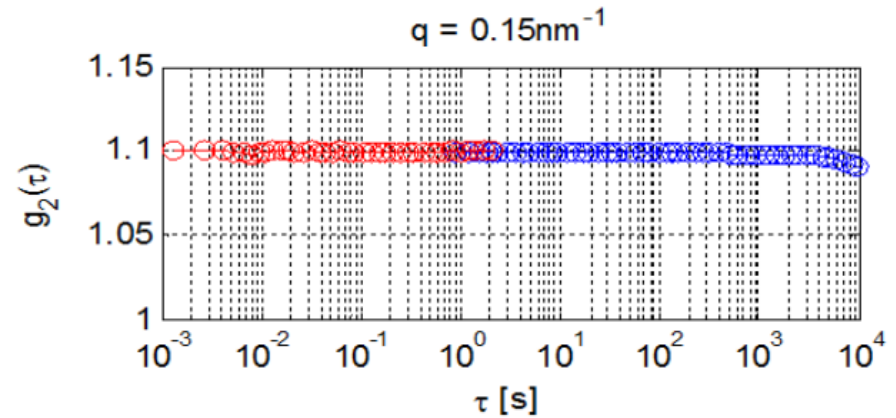
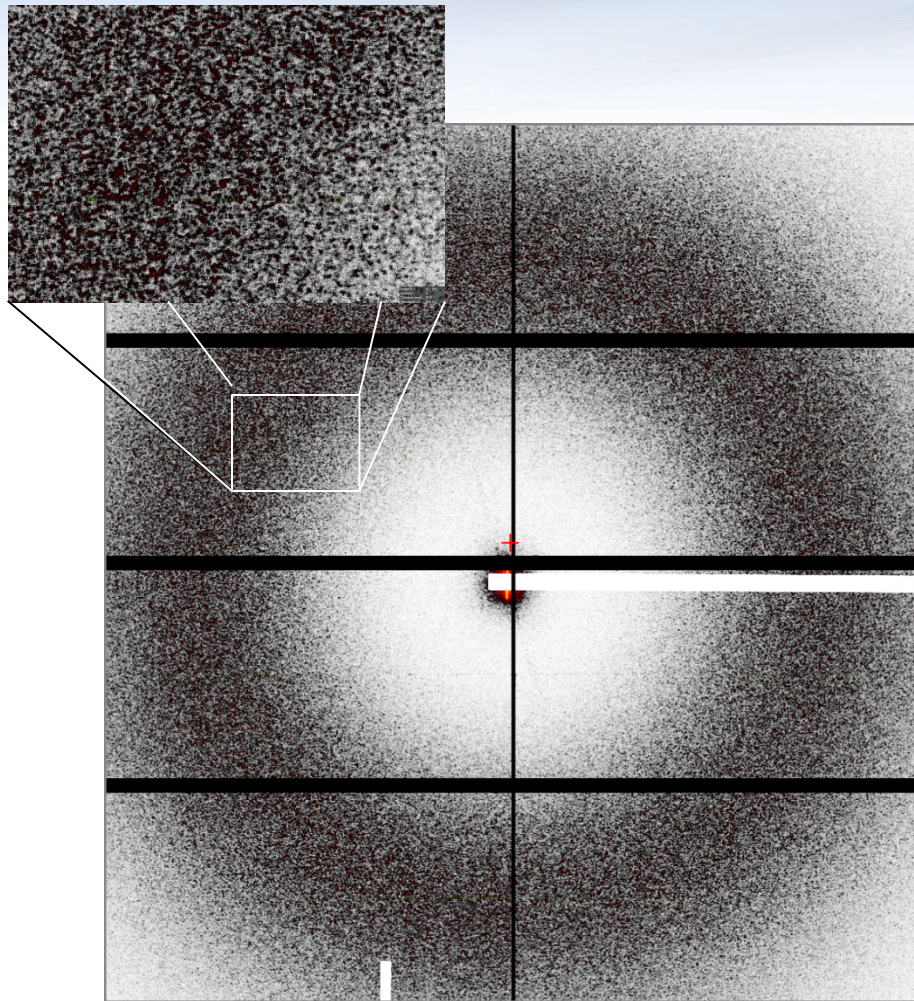
Dynamics of Materials (soft- and bio-): time and length scales



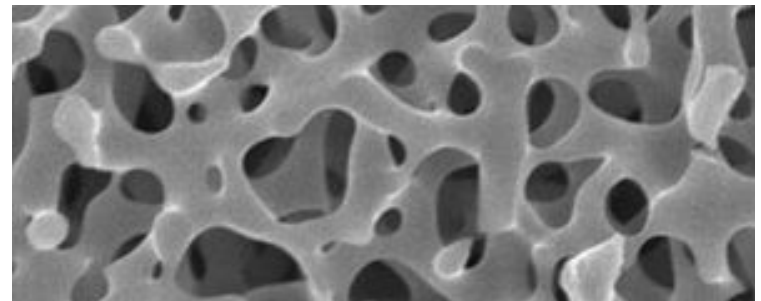
Dynamics of Materials (soft- and bio-): time and length scales



X-ray Speckles (Static!)



Correlation functions $g^{(2)}(q, \tau)$ measured from a CoralPor® static sample show excellent instrument stability.



www.schott.com

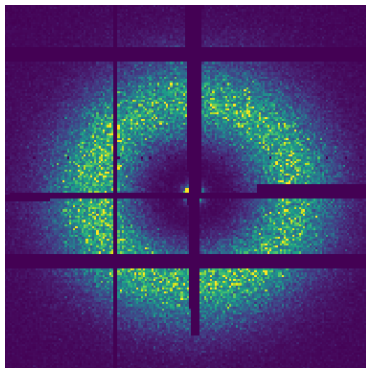
250mA top-off, 1.5×10^{11} ph/s in $10 \times 10 \mu\text{m}^2$; total dose = 101 seconds of “full flux”
Note: decay at $\sim 5 \times 10^3$ seconds due to ‘beam damage’

Speckles

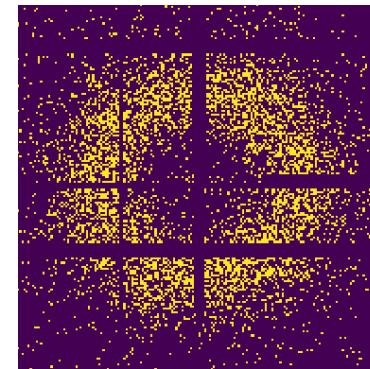
- Speckle statistics is described by the negative binomial distribution with
 - $M=M(q,T)$: # of coherent modes
 - $K=K(q,T)$: avg # of counts at a given q/ring
- Normalized variance becomes:

$$\text{var}_K(q, T) = \frac{1}{M(q, T)} + \frac{1}{K(q, T)}$$

Large $K(q, T)$



Small $K(q, T)$



Mandel, L. (1958). *Proc. Phys. Soc.* **72**, 1037.

Mandel, L. (1959). *Proc. Phys. Soc.* **74**, 233.

Goodman, J. W. (2007). *Speckle Phenomena in Optics: Theory and Applications*. Englewood: Roberts and Company.

Speckles & Speckle Visibility Spectroscopy

- Speckle statistics is described by the negative binomial distribution with
 - $M=M(q,T)$: # of coherent modes
 - $K=K(q,T)$: avg # of counts at a given q /ring

$$P(K) = \frac{\Gamma(K + M)}{\Gamma(K + 1)\Gamma(M)} \left(\frac{M}{\langle K \rangle + M} \right)^M \left(\frac{\langle K \rangle}{M + \langle K \rangle} \right)^K$$

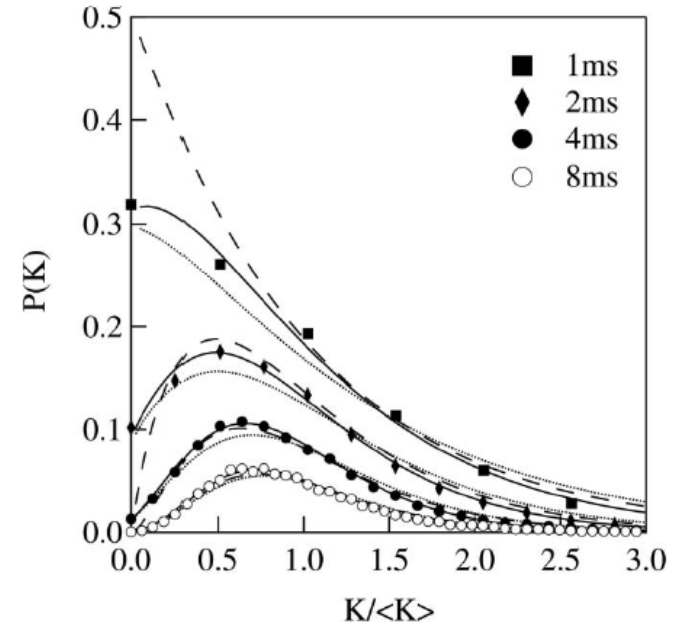
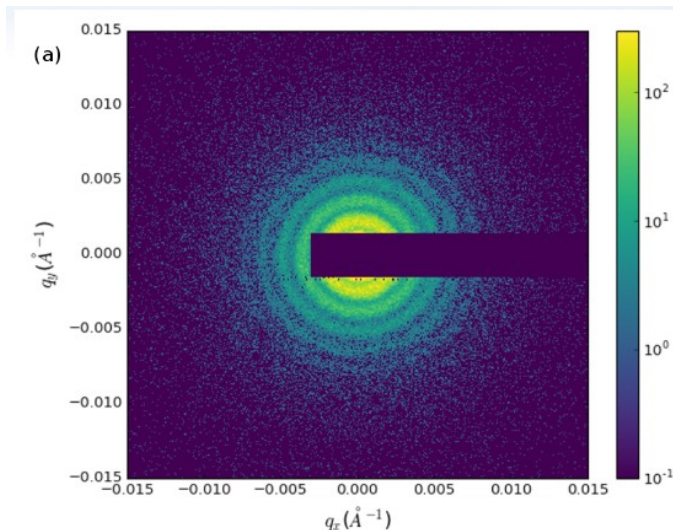
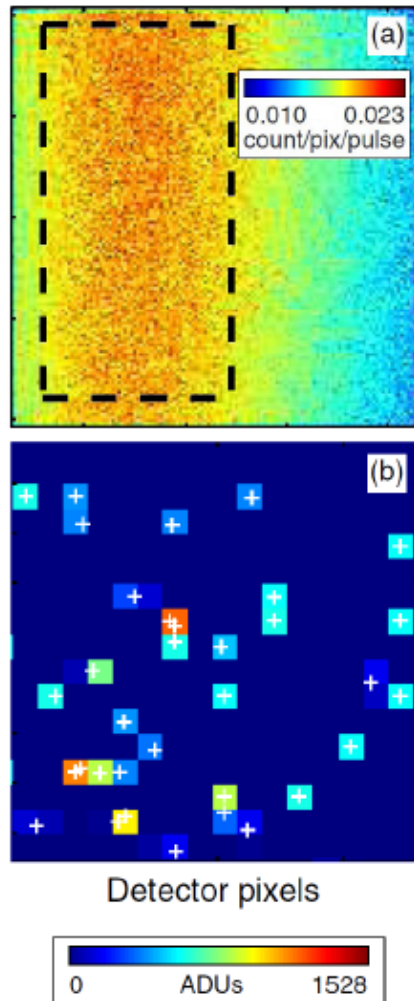


Figure 2

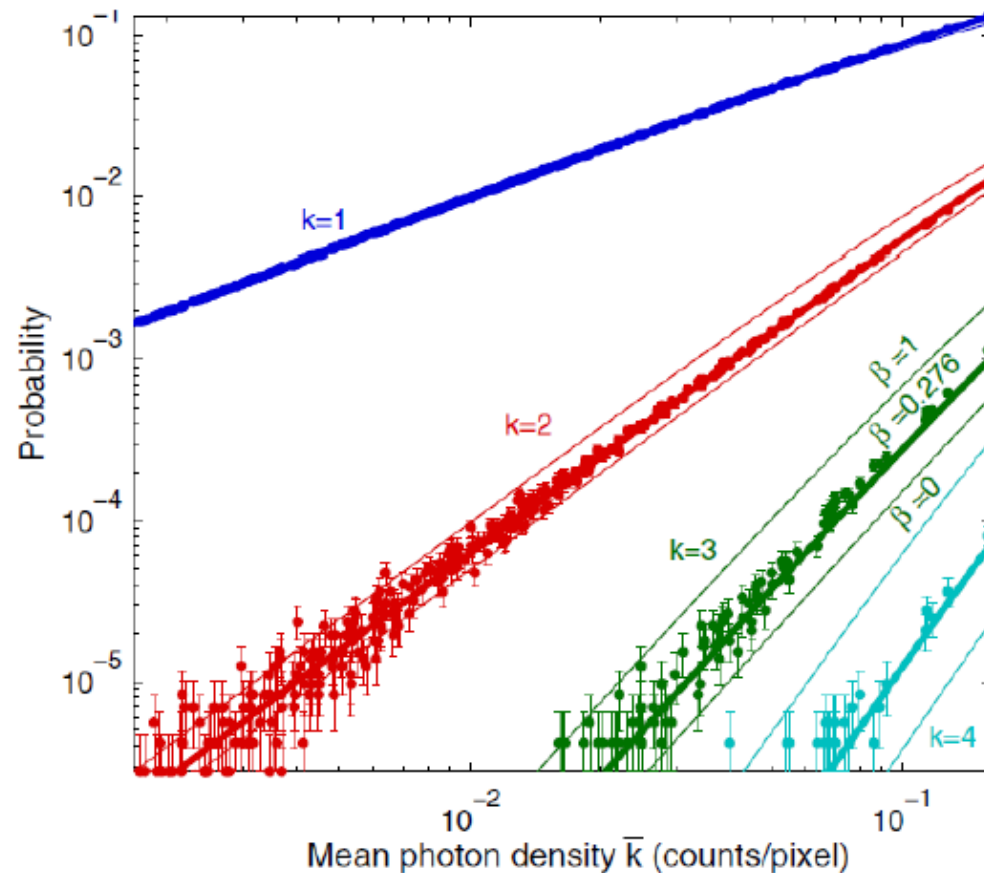
Photon count statistics analysis performed over an ensemble of pixels marked in the circular region in Fig. 1(a) for four integration times. Markers represent the photon count probability density $P(K)$ from the experiments, and solid lines are the fitting curves using the negative-binomial distribution [equation (11)], dashed lines are the fitting curves using the gamma distribution [equation (5)] and dotted lines are the fits using equation (11) with M as the only fitting parameter, while $\langle K \rangle$ is calculated from the measured photon counts. The results are plotted as a function of reduced count $K/\langle K \rangle$, so that $P(K)$ values with different integration times can be stacked in the same figure.

Luxi Li et al. J. Synch. Rad. 2014

Speckles from single shot LCLS pulses



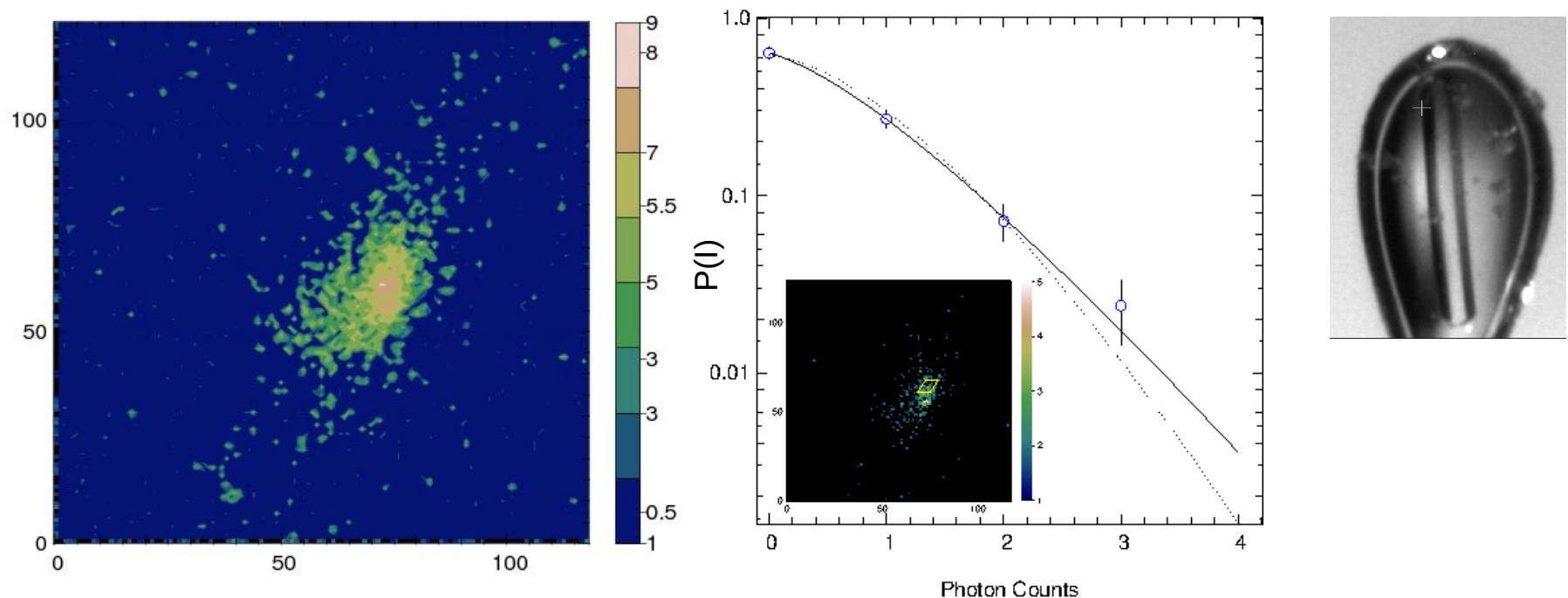
Single-shots at LCLS, Poisson-Gamma statistics



S. O. Hruszkewycz et al., PRL109, 185502 (2012)

X-ray Speckles come to life

- Molecular motion in protein microcrystals coupled over large scales generate diffuse scattering around the main Bragg peaks.



L. Li *et al.*, unpublished

Colloids

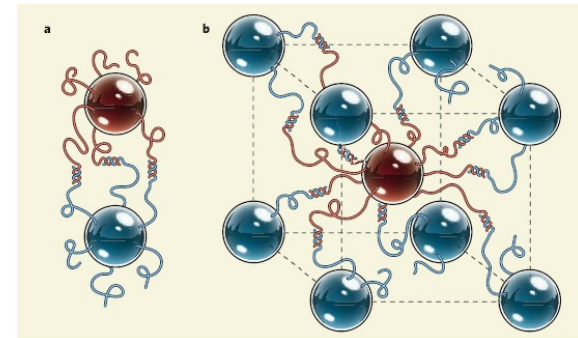
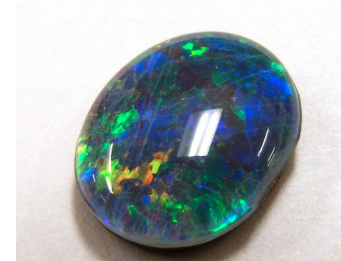
- Colloids are ubiquitous:

- Particles (1-1000 nm) of dispersed phase in dispersion medium



- Phase behavior;
The “magic” of self-assembly ...

- Opals are dried “polycrystalline” colloids” patchy colloids” can be elementary blocks for programmable self-assembly of “colloidal materials”
(O. Gang, BNL & Columbia)
- etc



Colloids: simple diffusive dynamics

- Intermediate Scattering Function

$$g^{(1)}(q, t) \propto \sum_{i=1}^N \sum_{j=1}^N \exp(iq[r_i(0) - r_j(t)])$$

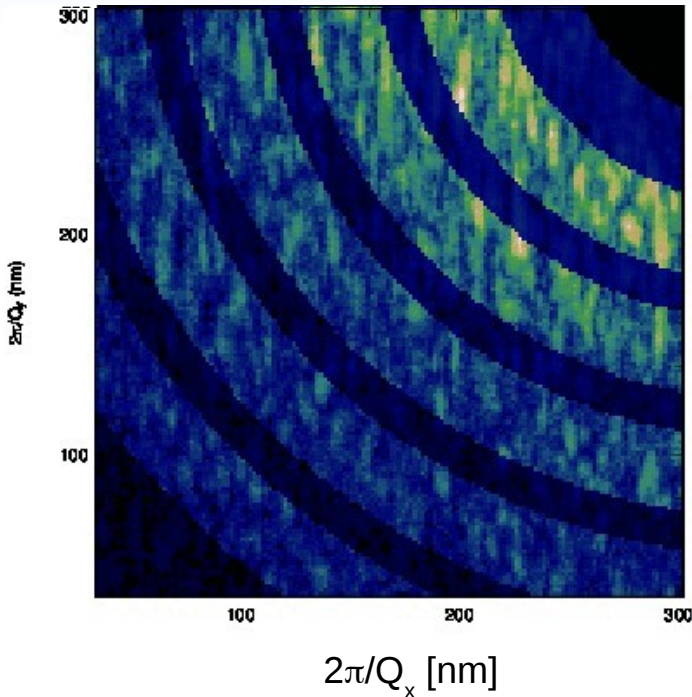
- Mean square displacement

$$\langle [r_i(0) - r_j(t)]^2 \rangle = 6 D_0 t \quad D_0 = \frac{k_B T}{6 \pi \eta a}$$

- Intermediate Scattering Function

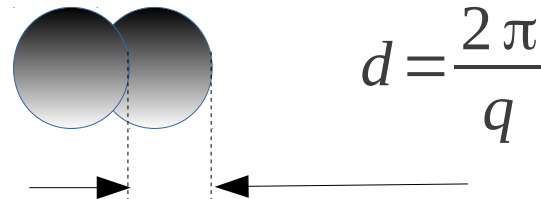
$$g^{(1)}(q, t) = \exp(-D_0 q^2 t)$$

Colloidal Dynamics with XPCS

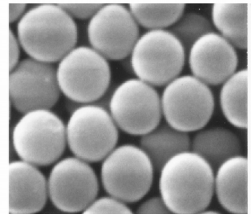


$$q = \frac{4\pi}{\lambda} \sin\left(\frac{2\theta}{2}\right)$$

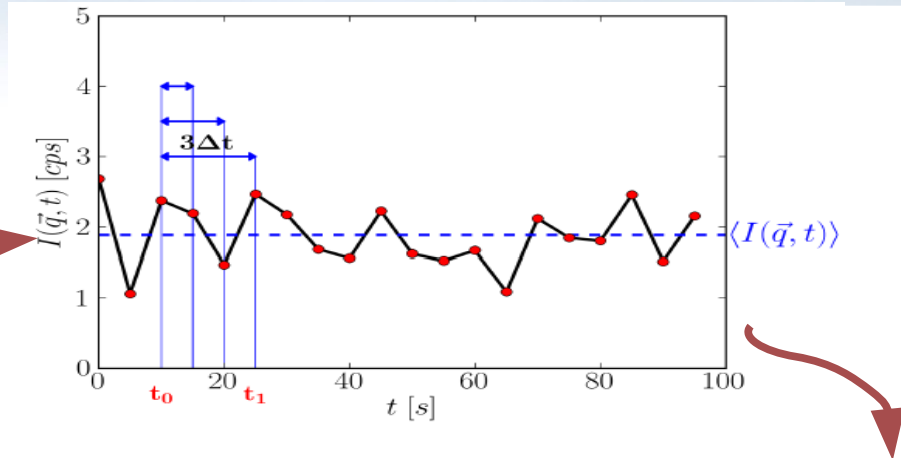
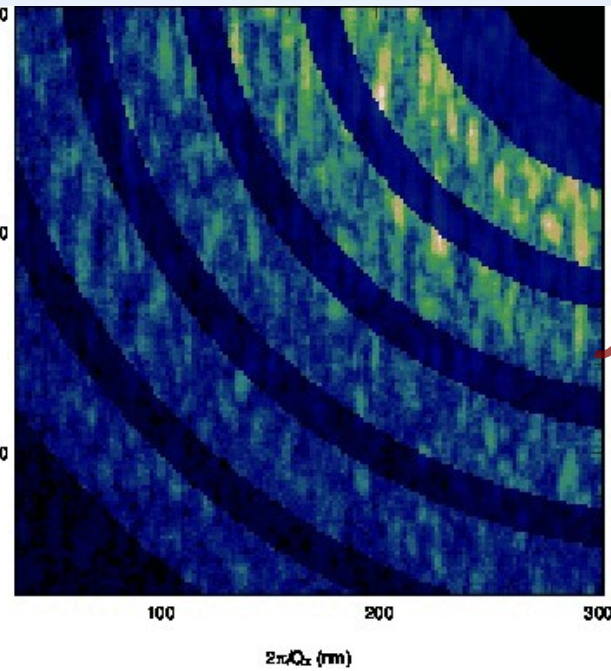
- Measures time scale associated with displacement of colloids



- i.e. measures dynamic structure factor $S(q,t)$
- By averaging over $\sim 10^{11}$ particles
- For different q values

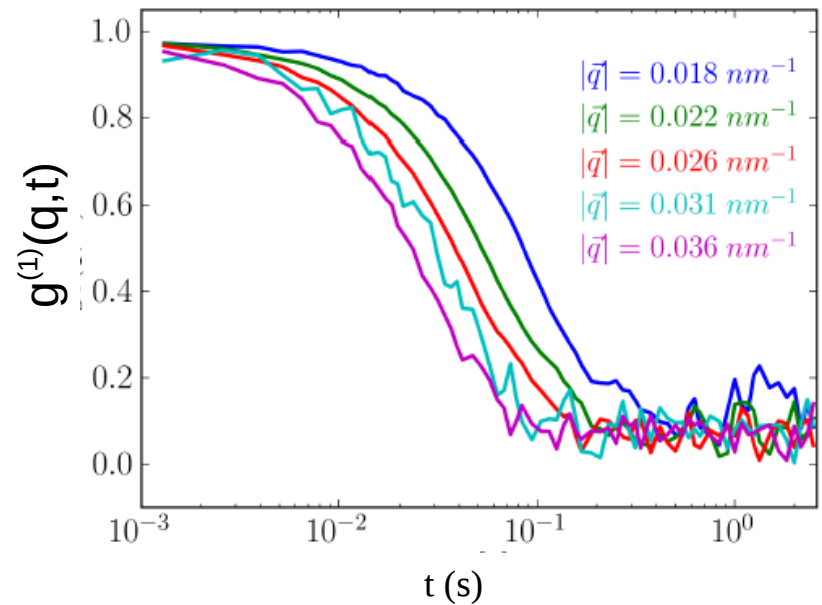


Colloidal Dynamics with XPCS

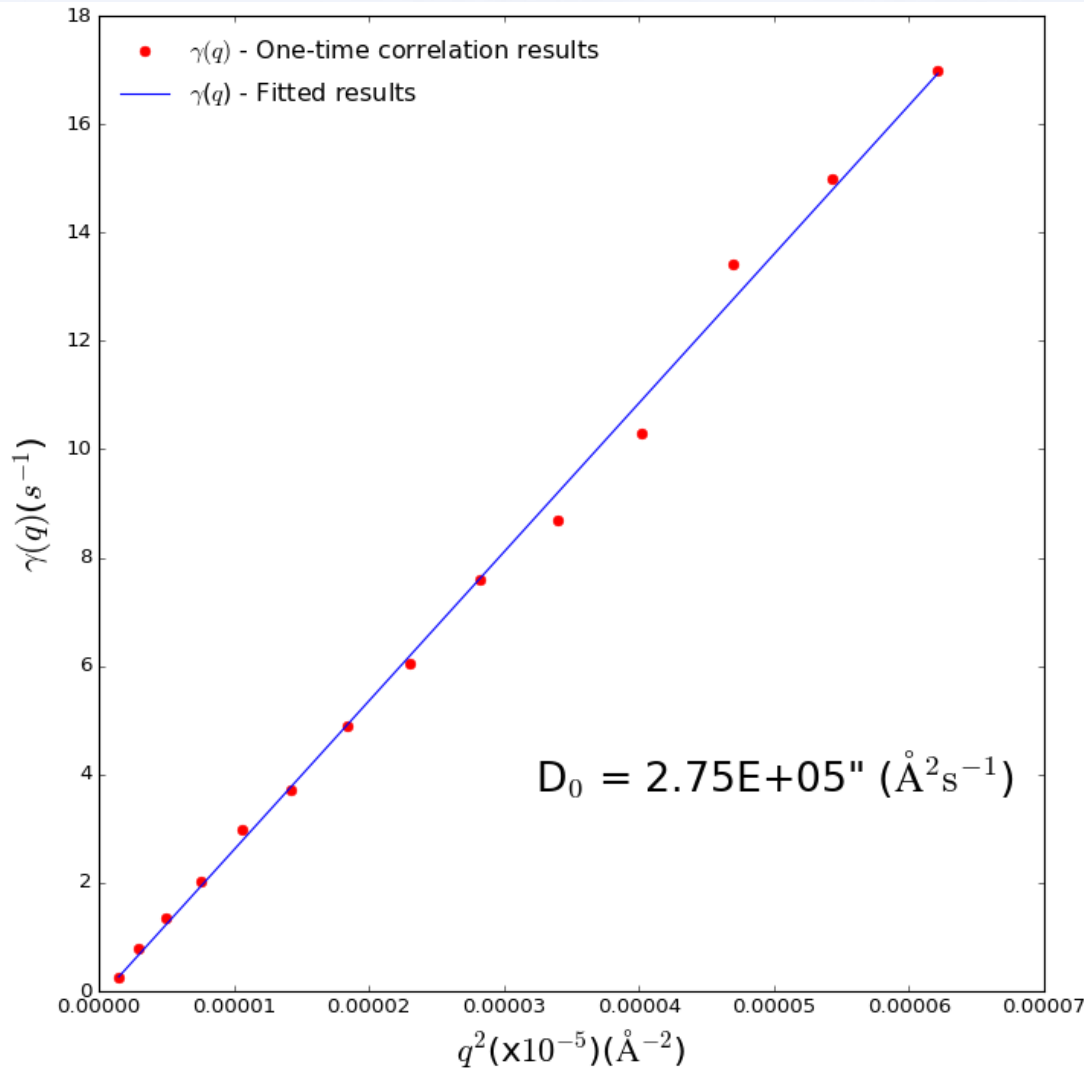


$$g^{(2)}(q, t) = \frac{\langle I(q, t) I(q, t + \delta t) \rangle}{\langle I(q) \rangle^2}$$

$$g^{(2)}(q, t) = 1 + \beta(q) [g^{(1)}(q, t)]^2$$



Colloidal Dynamics with XPCS



Here 500 nm Silica spheres suspended in a water/glycerol mixture

ISF:

$$g^{(1)}(q, t) \propto \exp[-Dq^2 t]$$

Relaxation rate:

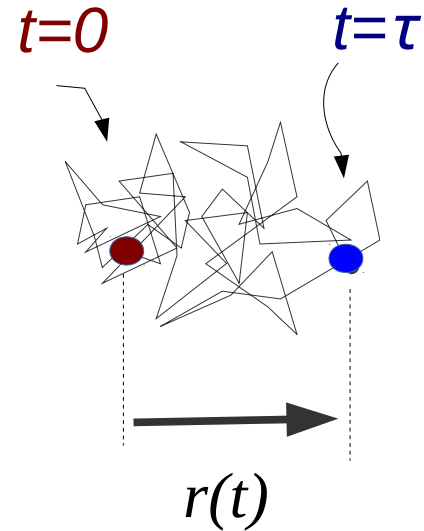
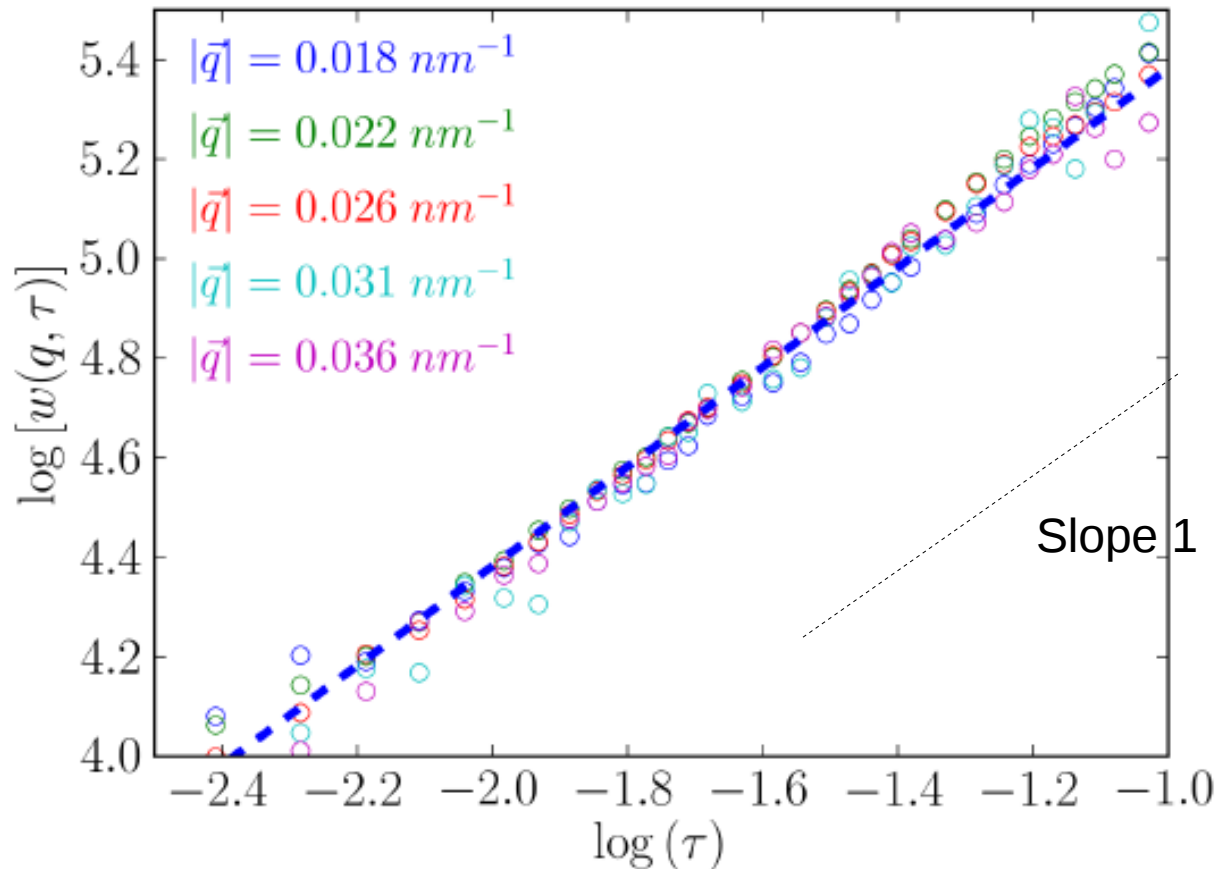
$$\Gamma = Dq^2$$

CHX Analysis Pipeline!

Colloidal Dynamics with XPCS

- Width function analysis
(Martinez, Van Meegen et al. JCP 2011)

$$w(q, t) = -\log[g^{(1)}(q, t)/q^2] \propto Dt \propto \langle r^2(t) \rangle$$



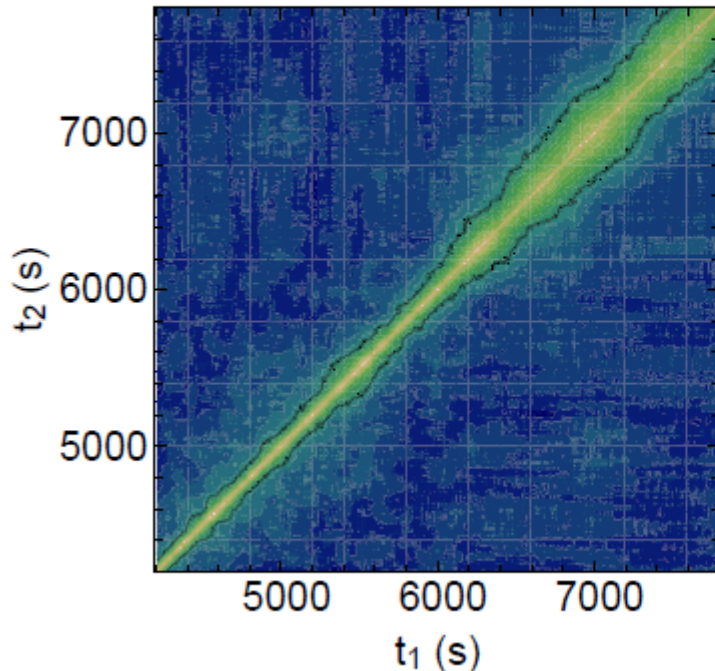
**Free diffusion
at low - Φ**

$$\langle r^2(t) \rangle \sim Dt$$

Two-time analysis

Non-equilibrium dynamics in colloidal depletion gels (colloid/polymer mixtures):

Two-time correlation functions: $C(Q, t_1, t_2) = \frac{\langle I(Q, t_1) I(Q, t_2) \rangle_{pix}}{\langle I(Q, t_1) \rangle_{pix} \langle I(Q, t_2) \rangle_{pix}}$

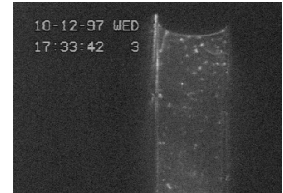


average time (“age”):

$$t_a = \frac{t_1 + t_2}{2}$$

time difference:

$$t = \delta t = |t_1 - t_2|$$

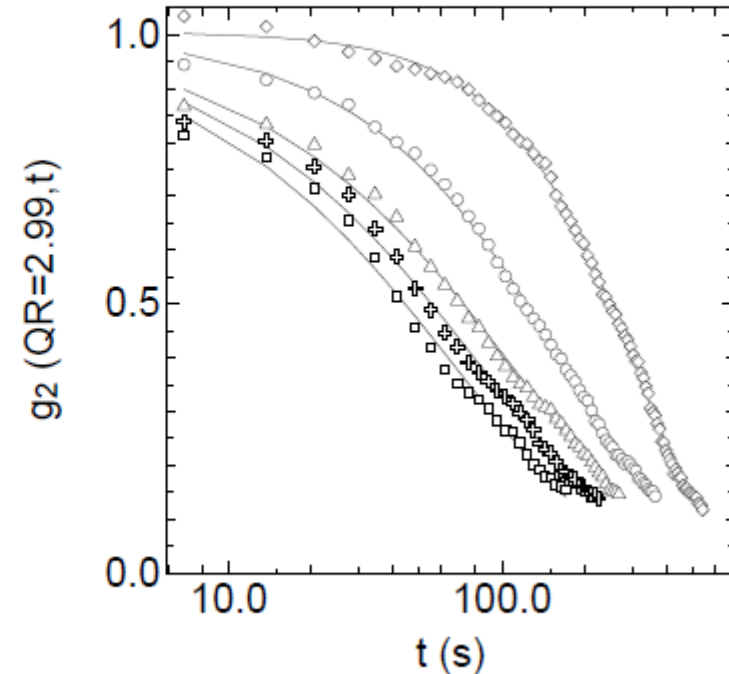
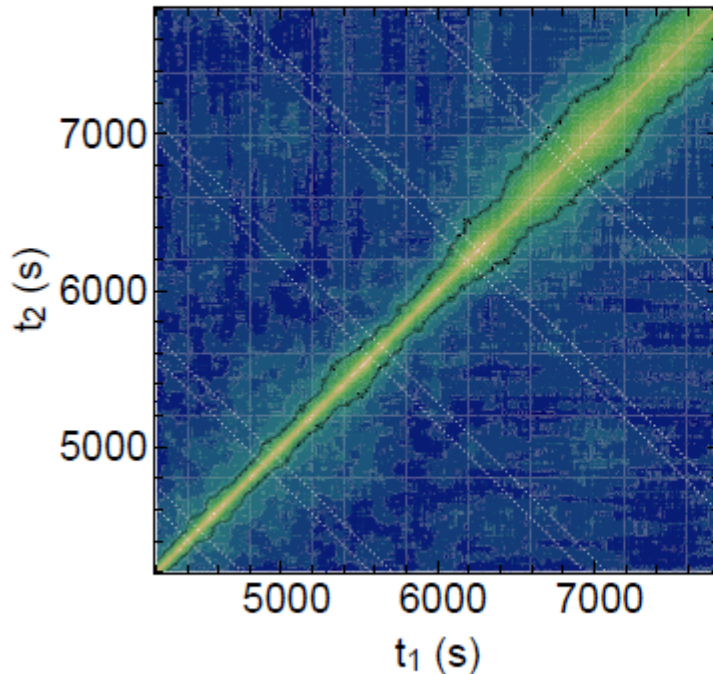


* M.Sutton et al., *Optics Express* **11**, 2268 (2003).

AF et al., *Phys. Rev. E*, **76**, 010401(R) (2007)

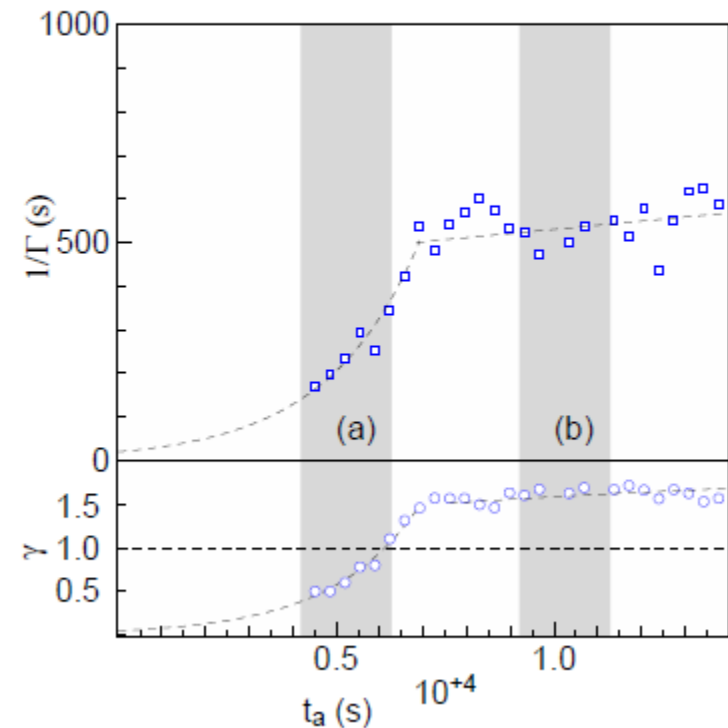
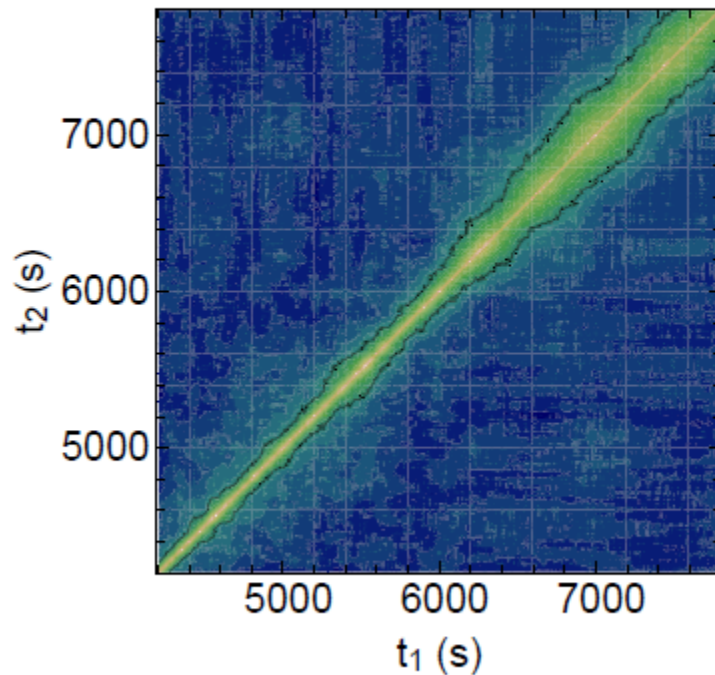
Two-time analysis

Two-time correlation functions: $C(Q, t_1, t_2) = \frac{\langle I(Q, t_1) I(Q, t_2) \rangle_{pix}}{\langle I(Q, t_1) \rangle_{pix} \langle I(Q, t_2) \rangle_{pix}}$



Two-time analysis

Two-time analysis: $g_2(Q, t_a, t) = \beta \exp(-(\Gamma t)^\gamma) + g_\infty$

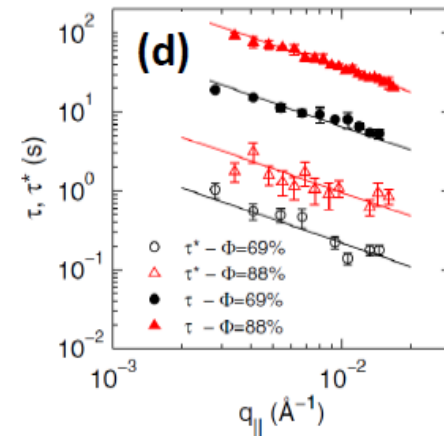
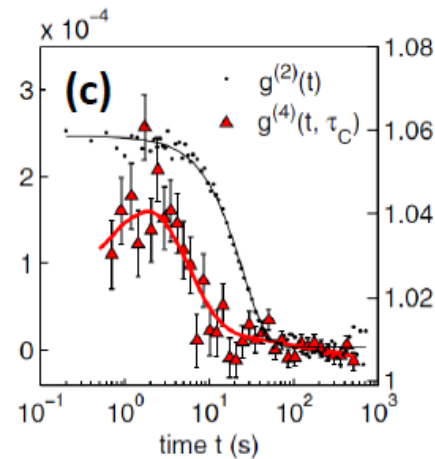
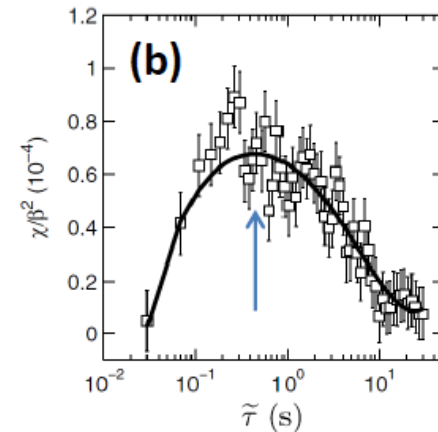
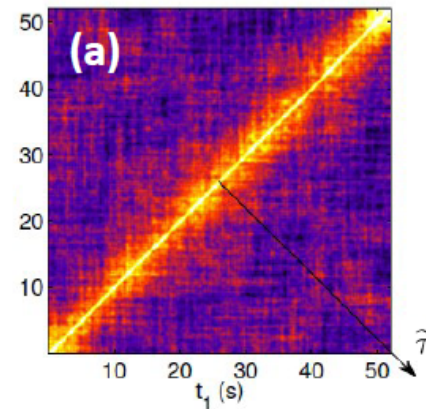


4th order correlations: dynamical heterogeneities

- Orsi et al. - dynamics in langmuir monolayer of nanoparticles using Grazing Incidence (GI)-XPCS
- Heterogeneities (correlations of correlations)

$$g^{(4)}(t, \tilde{\tau}) = \langle C(t_1, t_1 + \tilde{\tau}) C(t_1 + t, t_1 + t + \tilde{\tau}) \rangle_{t_1}$$

$$= \langle I(t_1) I(t_1 + \tilde{\tau}) I(t_1 + t) I(t_1 + t + \tilde{\tau}) \rangle_{t_1}$$



A. Duri et al., *Phys. Rev. E* **72**, 051401 (2005)

D. Orsi et al., *Phys. Rev. Lett.* **108**, 105701 (2012)

A “User Guide” to XPCS

- CHX optimized for Coherent X-ray Diffraction - XPCS, (GI-)SAXS/WAXS, CDI

Unprecedented q-range available in-situ from Angstroms to Microns

- Source: IVU 20 (low β) - highest brightness $E=6-15$ keV

- Beamline Optics: optimized for high stability & wavefront preservation

- COHERENT FLUX:
 $\approx 10^{11}$ ph/sec ($\Delta\lambda/\lambda=10^{-4}$)
 $\approx 10^{12}$ ph/sec ($\Delta\lambda/\lambda=10^{-3}$)

- BEAM SIZE :
 ≈ 10 μm (SAXS)
 ≈ 1 μm (WAXS)

DETECTORS

1. Diagnostics

- Fluorescent Screens; Pin diodes, Monitor counter; beam imaging; BPM

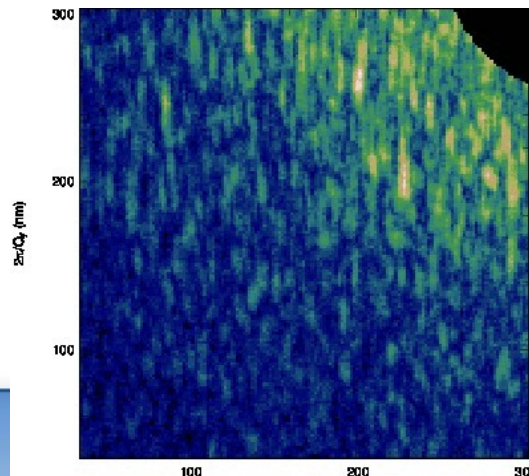
2. EIGER (Dectris)

best in class area detectors 3kHz (up to 15 kHz), 75 μm pixels

- Eiger 1M for c - WAXS
- Eiger 4M for c - SAXS

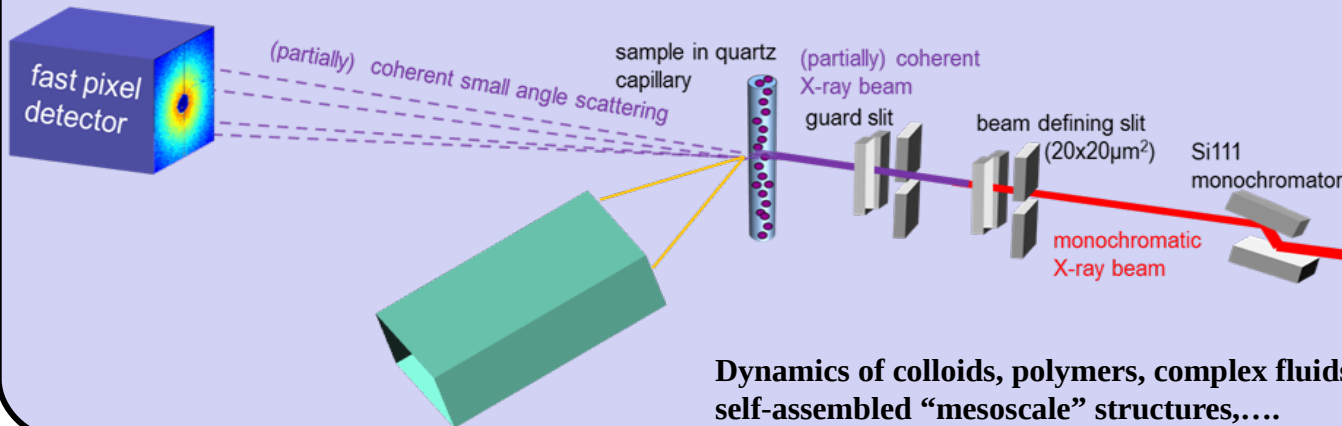
3. Point Detectors (FMB Oxford)

- Scintillator detector systems;
- Avalanche Photodiode (APD)



Example Scattering Geometries

(Transmission) SAXS / WAXS

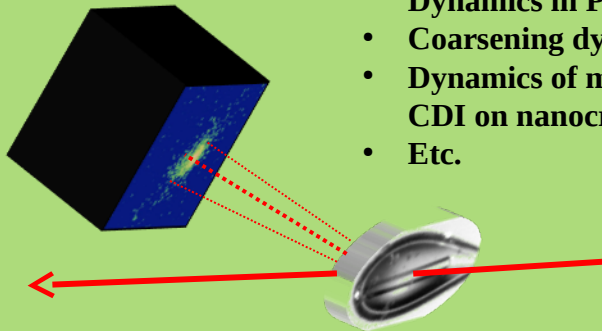


Next deadline for Experiment proposals:
Sept 30, 2016

<https://pass.bnl.gov>

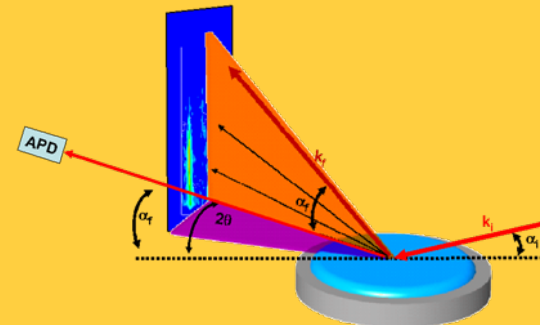
Dynamics in Protein Crystals,

- Coarsening dynamics in alloys
- Dynamics of metallic glasses
- CDI on nanocrystals
- Etc.



Coherent WAXS

GI-SAXS / GI-WAXS, XRR (solid)

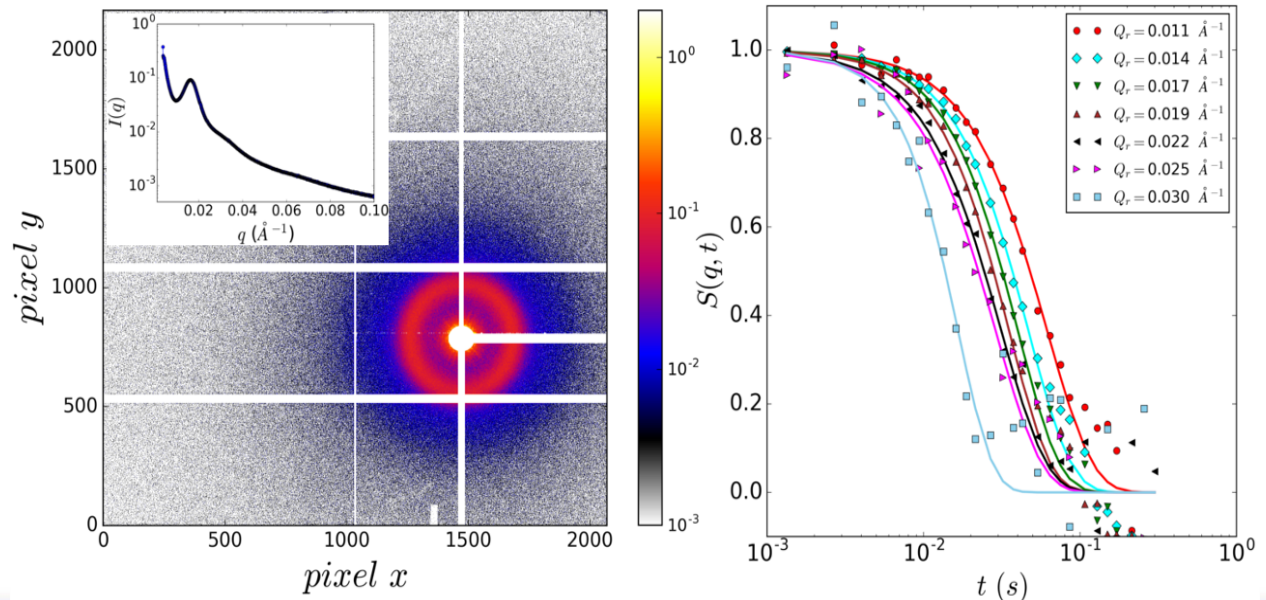


In-plane structure/diffusion,
capillary waves, electron density
profile

A “Mini User Guide” to XPCS

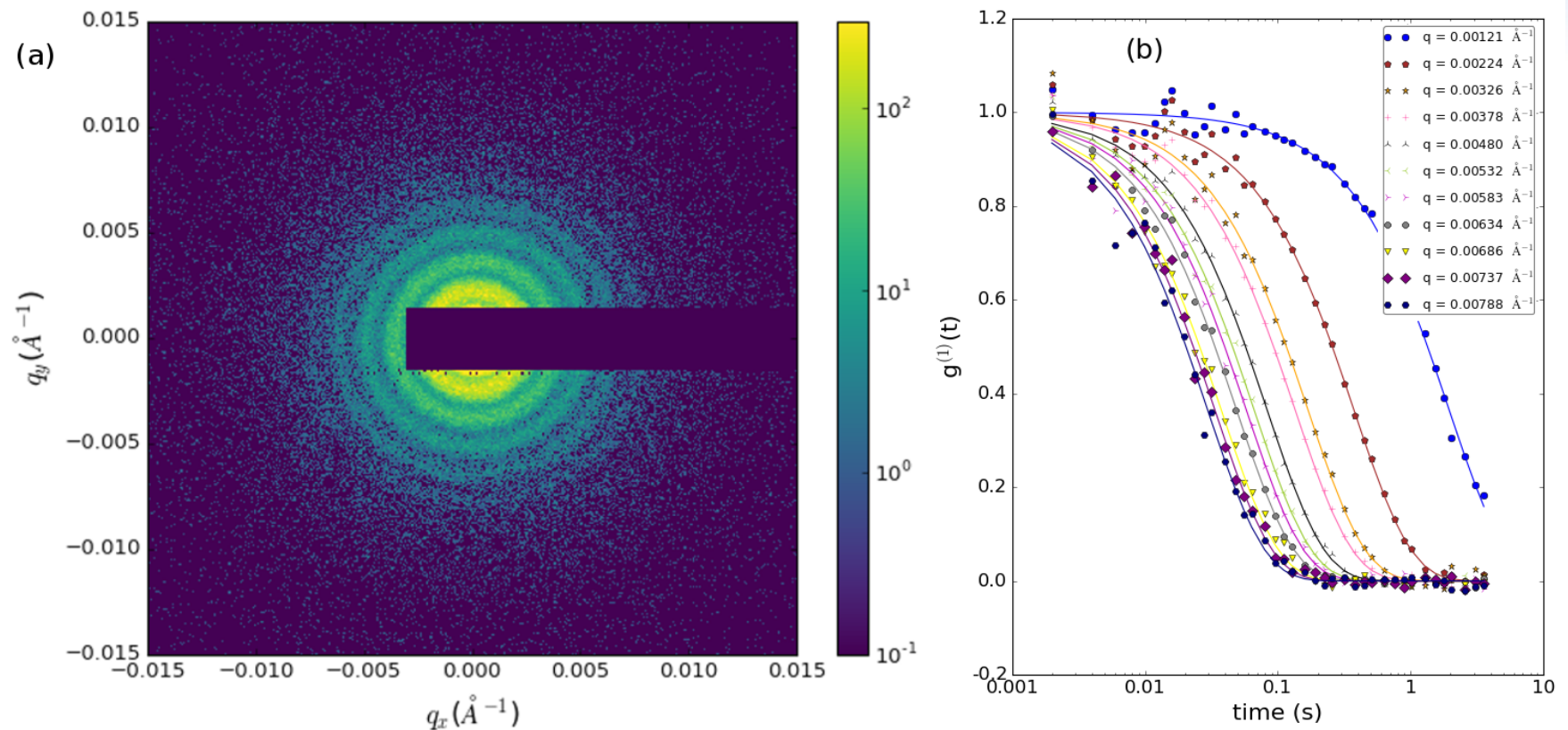
Questions:

- How much does the sample scatter?
 - we need $\sim 10^{-N}$ ph/correlation time/speckle(pixel) - $g^{(2)}$
 - We need $\sim 1/\text{ph/correlation time/speckle(pixel)} - C(t_1, t_2)$
- What time scales are we expecting?
- What is the radiation limit? Is the sample homogeneous? i.e can we build an ensemble by averaging information recorded from different locations?



A “Mini User Guide” to XPCS: Data Analysis

CHX Data Analysis Solutions: <https://github.com/NSLS-II-CHX>



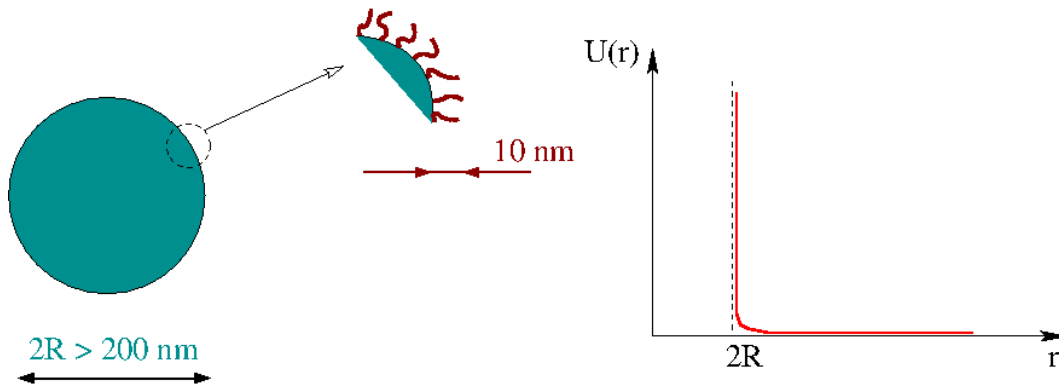
XPCS experiments on the dynamics of silica colloids ($R=250$ nm) suspended in a polymer solution of polypropyleneglycol (PPG) in water.

(a) A single speckle pattern recorded in 2 ms from the colloidal suspension.

(b) Intermediate scattering function (dynamic structure factor)

S. K. Abeykoon *et al.*, 2016 New York Scientific Data Summit (NYSDS), New York, NY, 2016, pp. 1-10. doi:10.1109/NYSDS.2016.7747815

A more detailed science example: high density hard-sphere (colloidal) suspensions

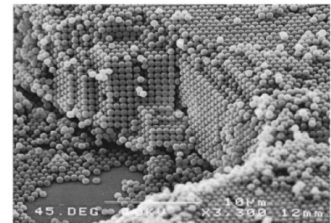
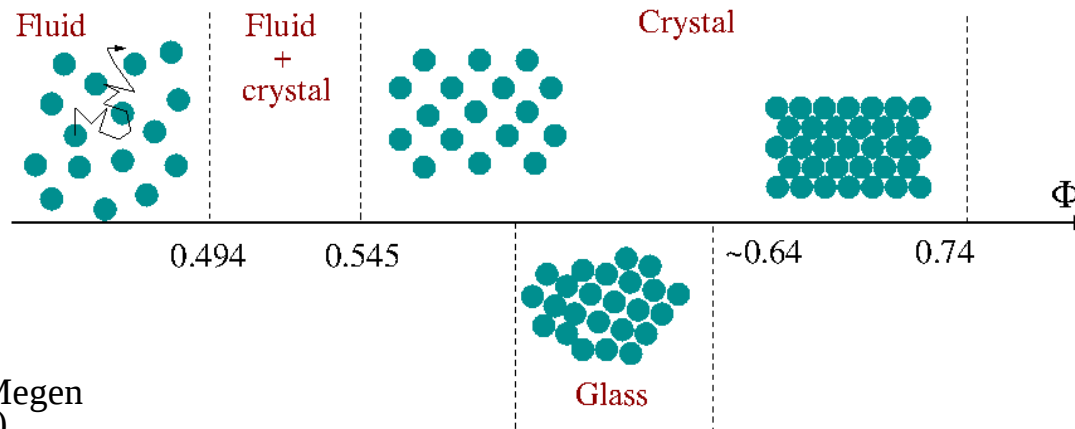


Hard-sphere colloids:

- Spherical PolyMethylMethacrylate (PMMA) particles coated with 12 hydroxystearic acid in cis-decalin (A. Schofield, Edinburgh)
- Entropic forces between polymer coating layers → infinite “hard-sphere-like” repulsions

- The phase behavior depends on the *particle volume fraction* Φ

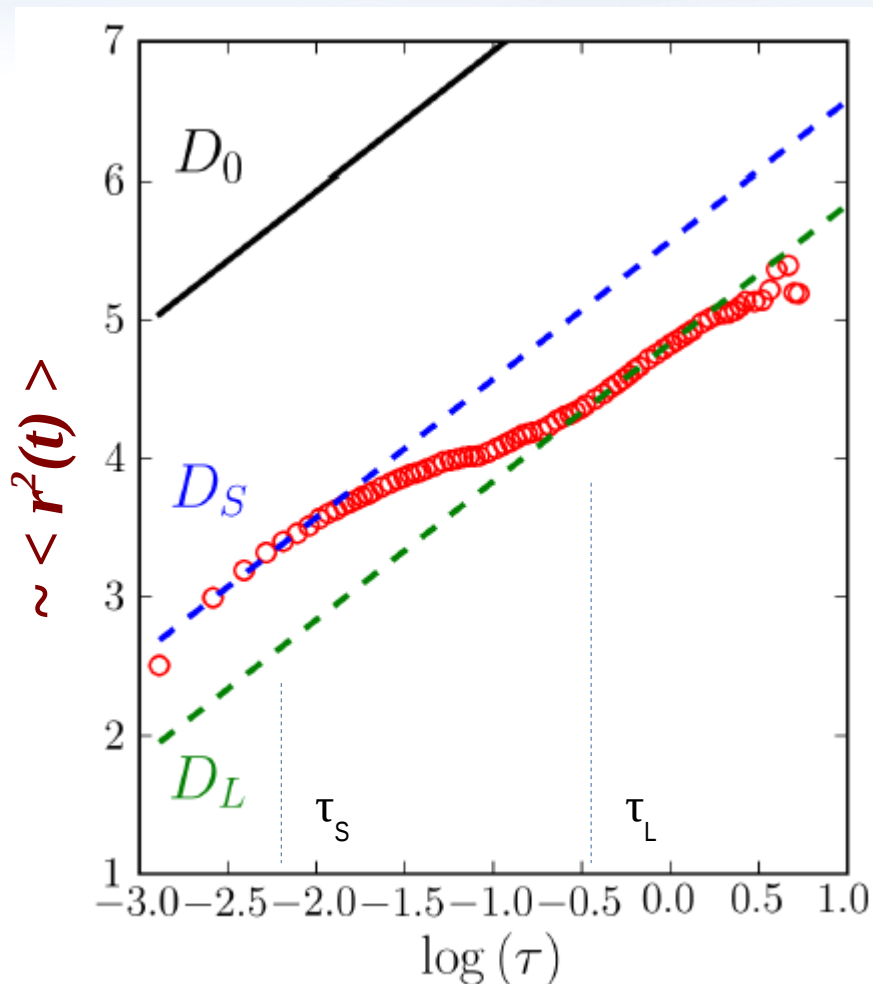
$$\Phi = \frac{N V_{colloid}}{V_{total}}$$



(b) R.M. Amos et al., PRE 61, 2929 (2000)

P.N. Pusey & W. Van Meegen
Nature **320**, 340 (1986)

Dynamics in high density hard-sphere suspensions



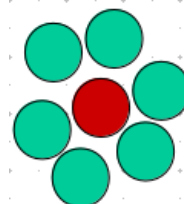
P. Kwasniewski, PhD Thesis 2012

Short-time diffusion D_s ($t < \tau_s$)

Motion of particles inside of “cages”
created by other particles

Slowed down (compared to D_0) by
hydrodynamic interactions

D. Orsi, AF et al. *Phys. Rev. E* 2012



Long-time diffusion D_L ($t > \tau_L$)

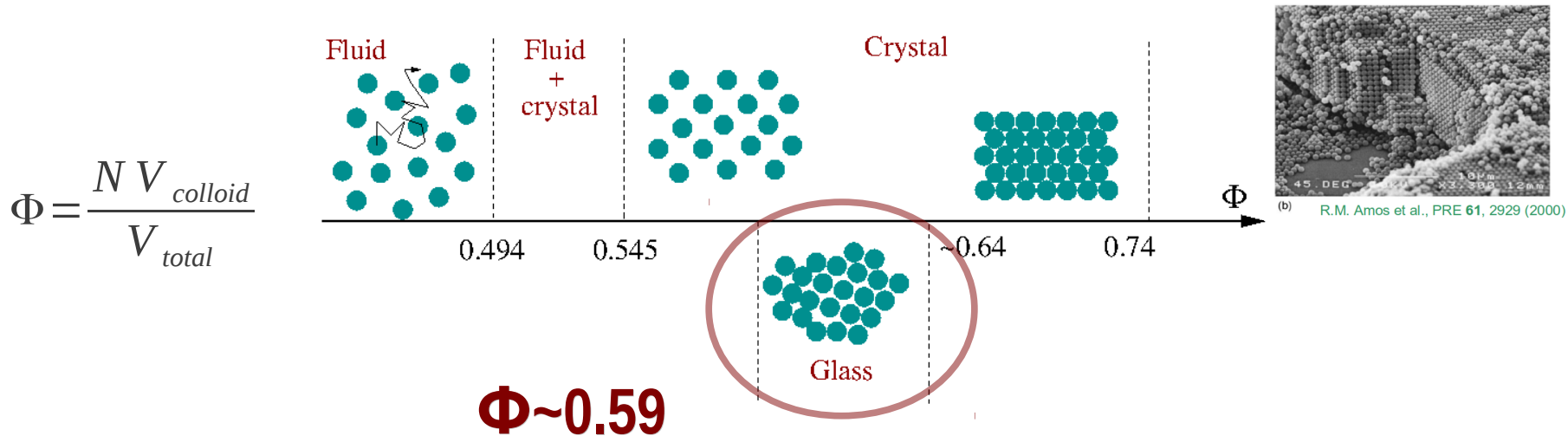
Structural rearrangements i.e.
“Rearrangements of cages”

Slowed down (compared to D_s) by
direct interactions

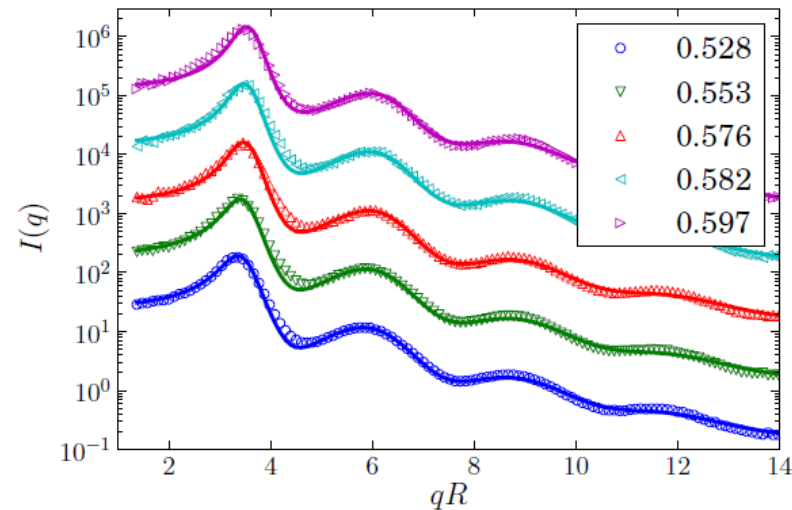
P. Kwasniewski, AF, A. Madsen, *Soft Matter*, 2014, **10**, 8698-8704

The Colloidal Glass Transition

- What happens here?

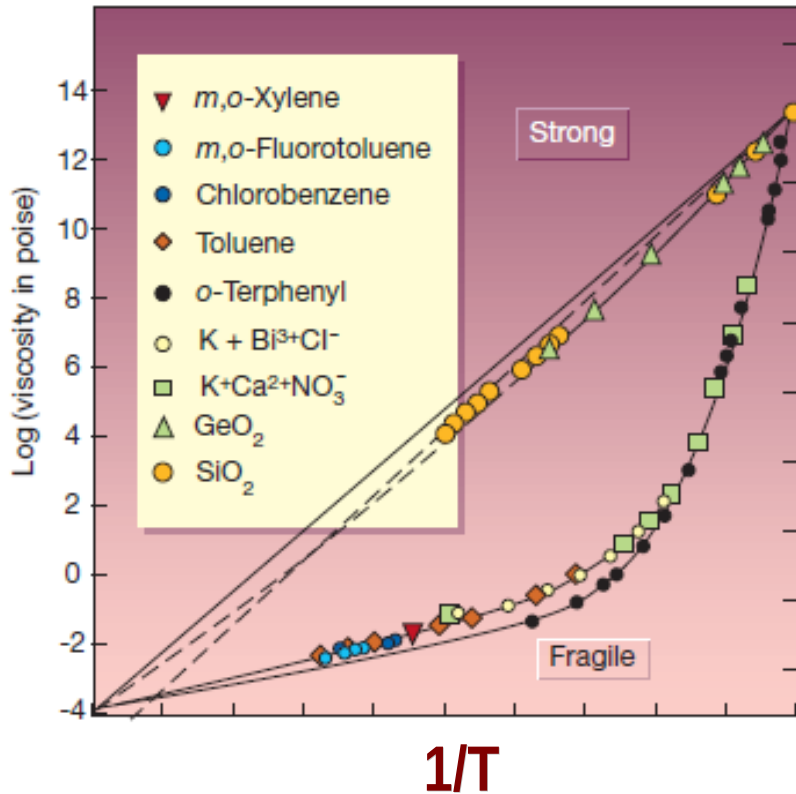


- From SAXS / static scattering: pretty much *nothing* ...

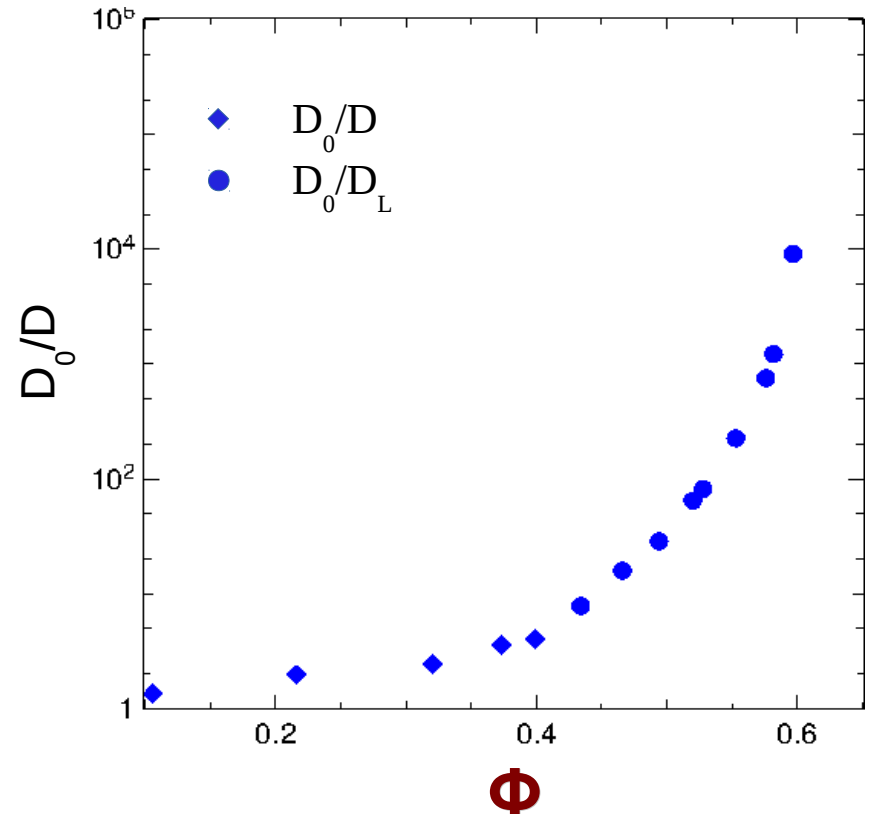


Supercooled Liquids vs. Hard-Sphere Colloids

- In addition to being interesting/useful in their own right, colloids are an excellent model system for supercooled liquids and molecular glassformers

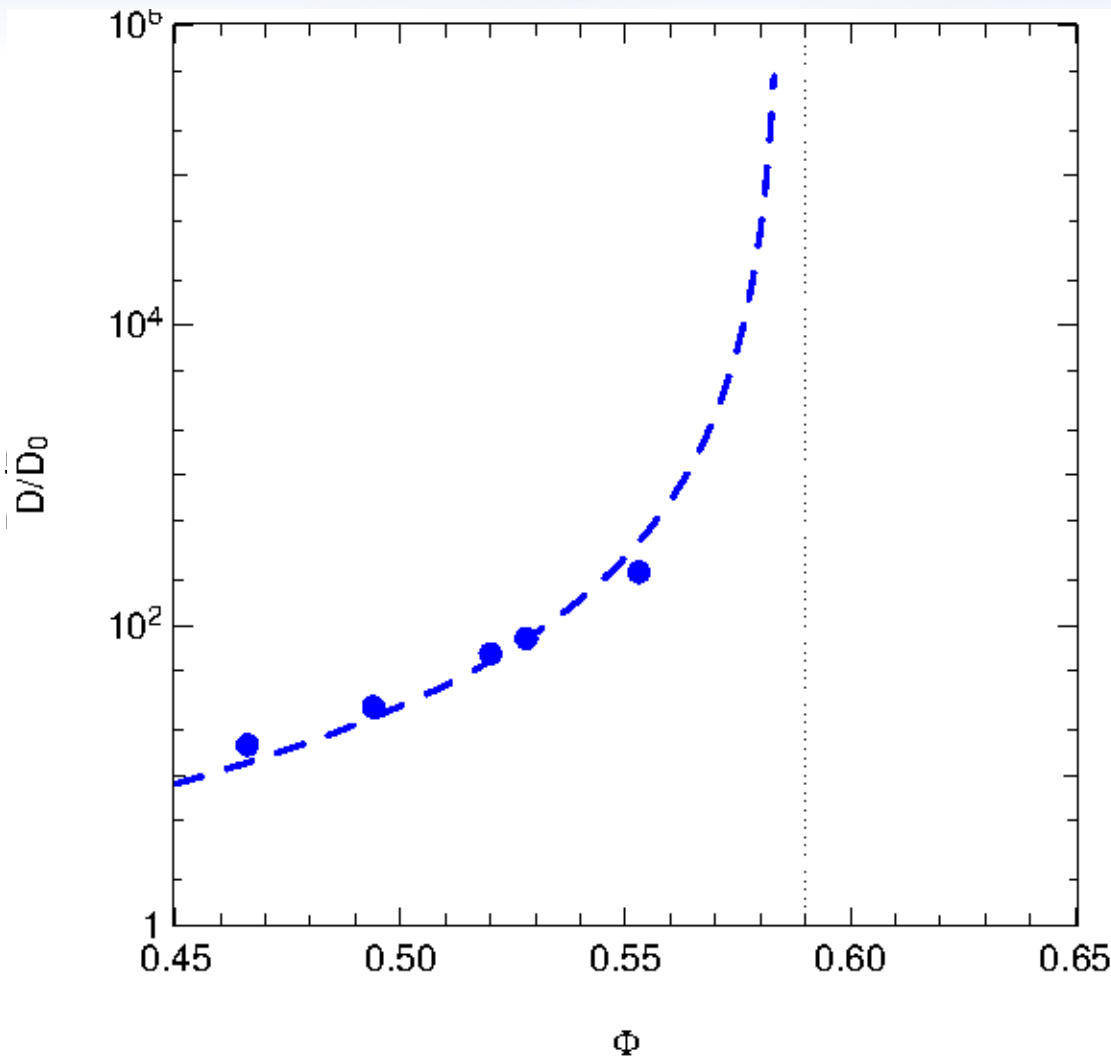


Denenedetti, Stillinger, *Nature* 2001



D. Orsi, AF et al. *Phys. Rev. E* 2012
 P. Kwasniewski, AF, A. Madsen, *Soft Matter* 2014
 $\eta/\eta_0 \rightarrow D_0/D_L$ (Segre et al., *Phys. Rev. Lett* 2001)

Structural Relaxations near the Hard-Sphere Glass Transition

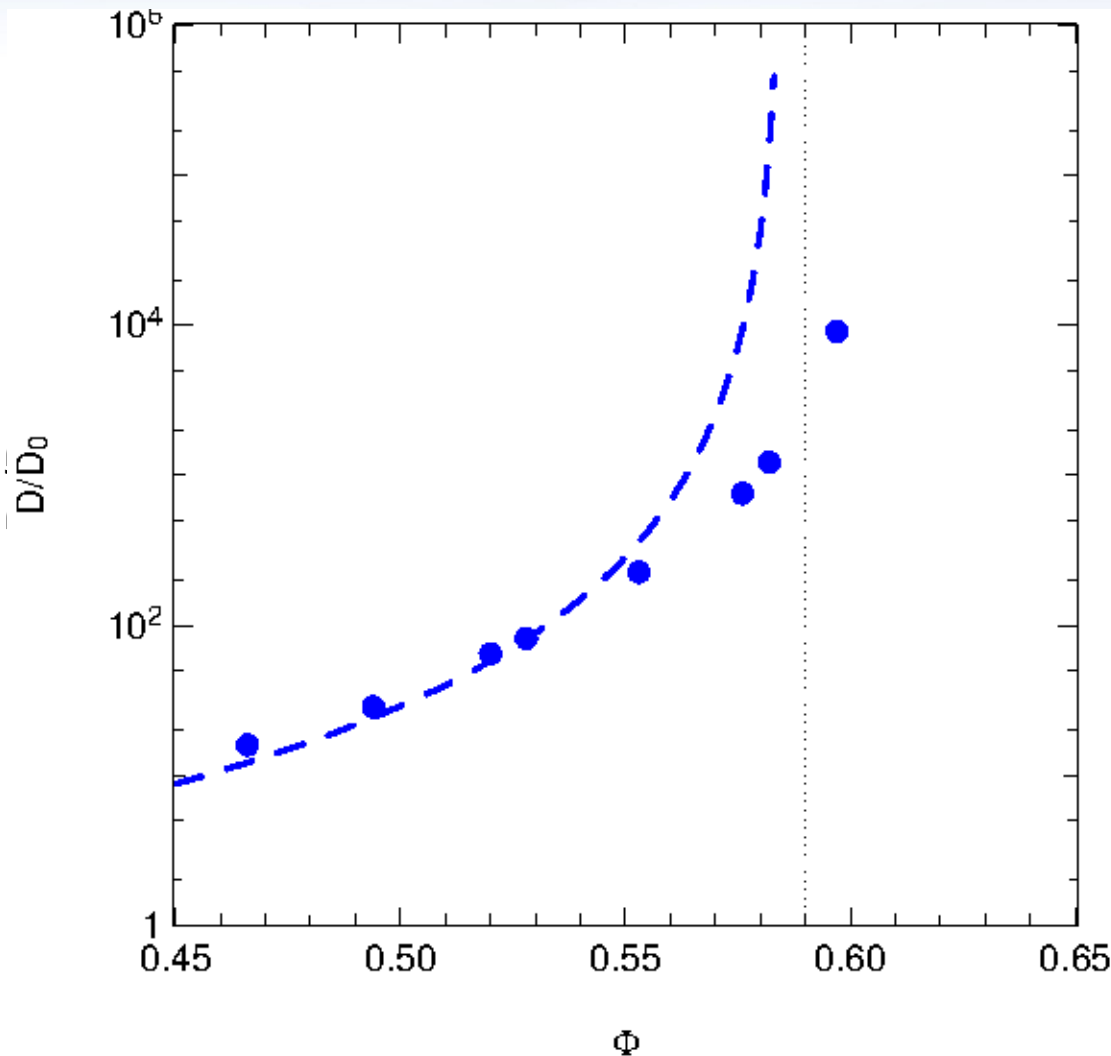


Structural relaxations:

- Structural relaxations slow-down with increasing Φ
- And are expected to *diverge* at the colloidal glass transition concentration Φ_g - "Mode Coupling Theory"-(MCT)
- $D_0/D_L \rightarrow \infty$ at $\Phi_g \sim 0.59$

$$\frac{D_0}{D_L(q_m)} \propto \left| \frac{\Phi_g - \Phi}{\Phi_g} \right|^{-\gamma}$$

Structural Relaxations near the Hard-Sphere Glass Transition



Structural relaxations:

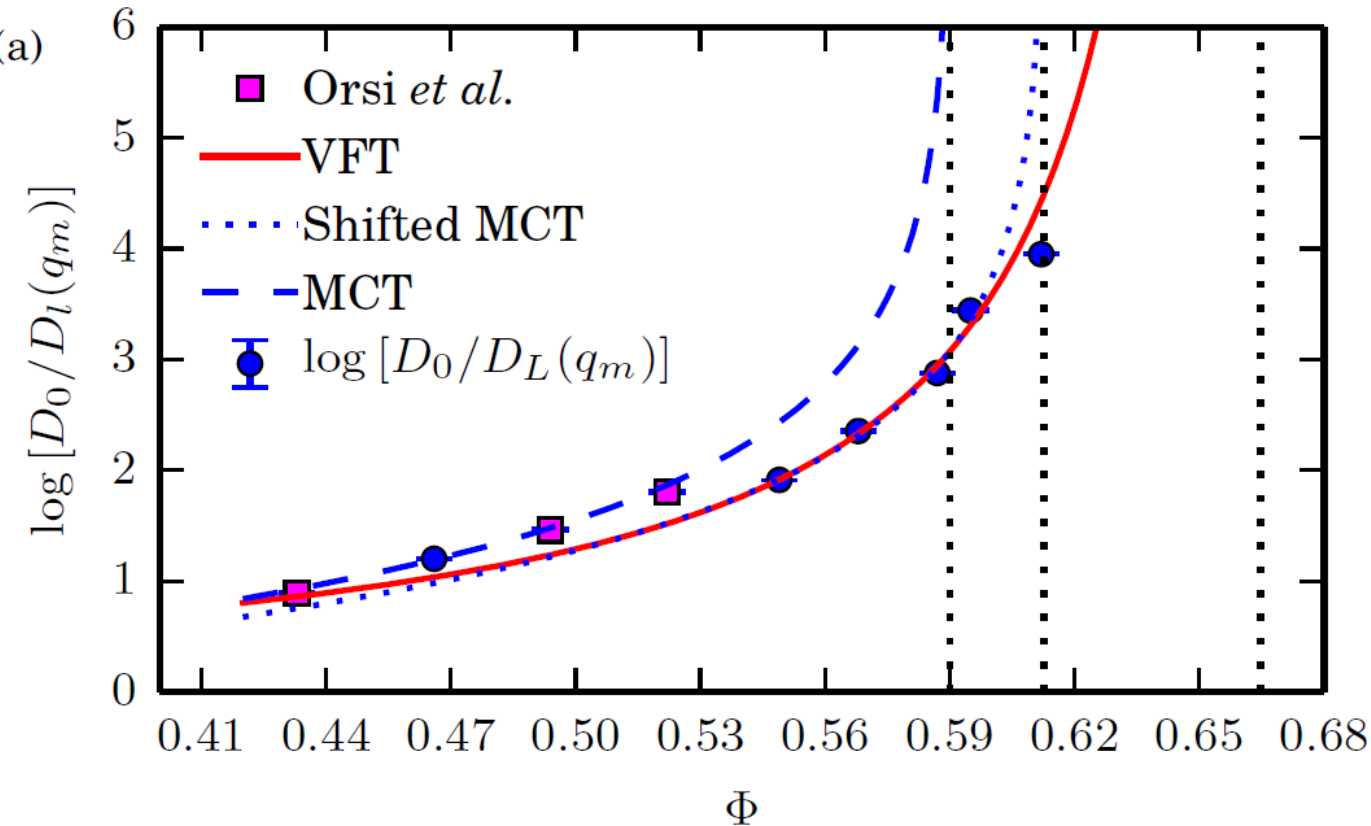
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$$\frac{D_0}{D_L(q_m)} \propto \left| \frac{\Phi_g - \Phi}{\Phi_g} \right|^{-\gamma}$$

Not so simple:

- Instead of diverging the relaxations remain finite (but slow!) above Φ_g

Structural Relaxations near the Hard-Sphere Glass Transition



MCT:

$$\frac{D_0}{D_l(q_m)} \propto \left| \frac{\Phi_g - \Phi}{\Phi_g} \right|^{-\gamma}$$

$g \sim 2.58$

VFT:

$$\frac{D_0}{D_l(q_m)} = \tau_\infty \exp \left[\frac{F}{(\Phi_0 - \Phi)^\delta} \right]$$

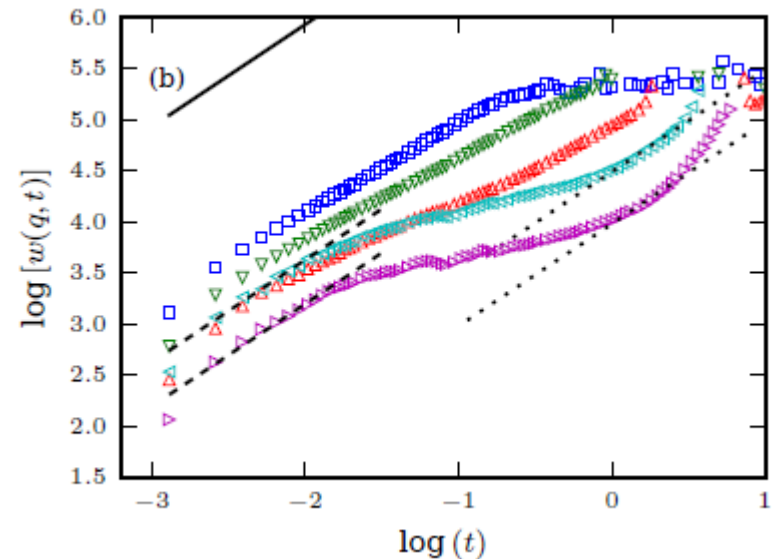
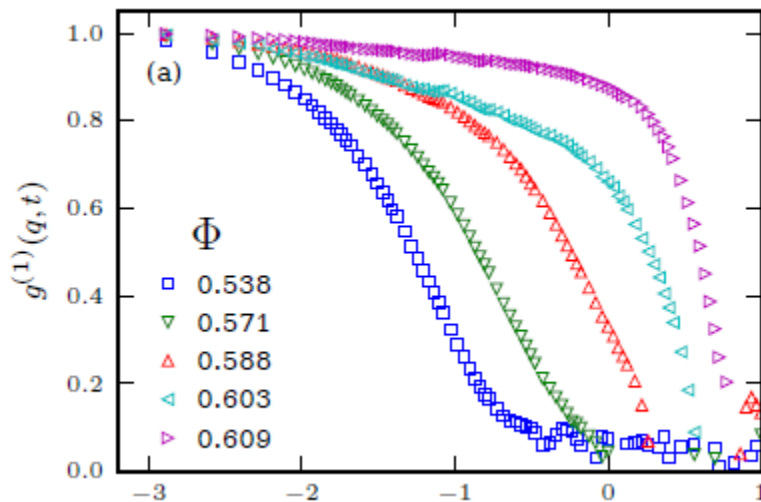
- relaxations follow an unexpected functional (VFT) form suggesting a kinetic arrest near the “random close packing concentration $\Phi_{RCP} \sim 0.67$ (~10% polydispersity)
- Suggests connection with **Jamming**

P. Kwasniewski, AF, A. Madsen, *Soft Matter*, 2014, 10, 8698-8704

See also; Brambilla, Cipelletti *et al.*, *Phys. Rev. Lett.* 104, 169602 (2010)

Anomalous Dynamics near the Hard-Sphere Glass Transition

- Near the colloidal Glass Transition the dynamics becomes anomalous
 - Compressed exponential relaxations
 - Hyperdiffusive dynamics: $\langle r^2(t) \rangle$ “faster than” $\sim t$



- Is this behavior a signature of *jamming*?

Universal non-diffusive slow dynamics in aging soft matter

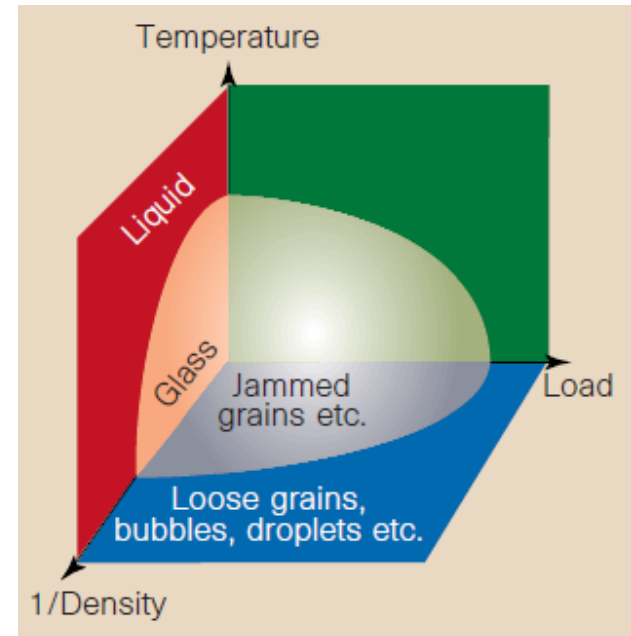
L.Cipelletti *et al.*, *Faraday Discuss.*, 2003, **123**, 237

Jamming?

- Is this behavior a “universal” ?
- Common behavior in seemingly different systems: hyperdiffusive & faster-than-exponential relaxations associated with *Jamming*

L.Cipelletti *et al.*, *Faraday Discuss.*, 2003, **123**, 237

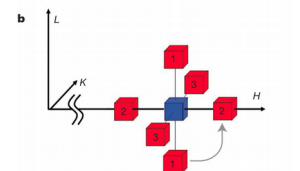
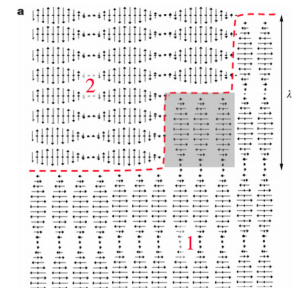
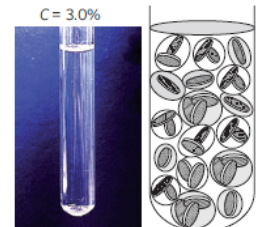
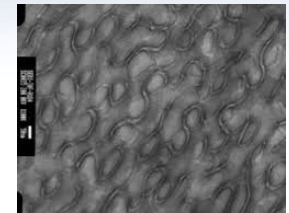
- Jamming – heterogeneities & response to flow/shear



A. Liu *et al.* Nature 1998

Anomalous Dynamics near the Hard-Sphere Glass Transition

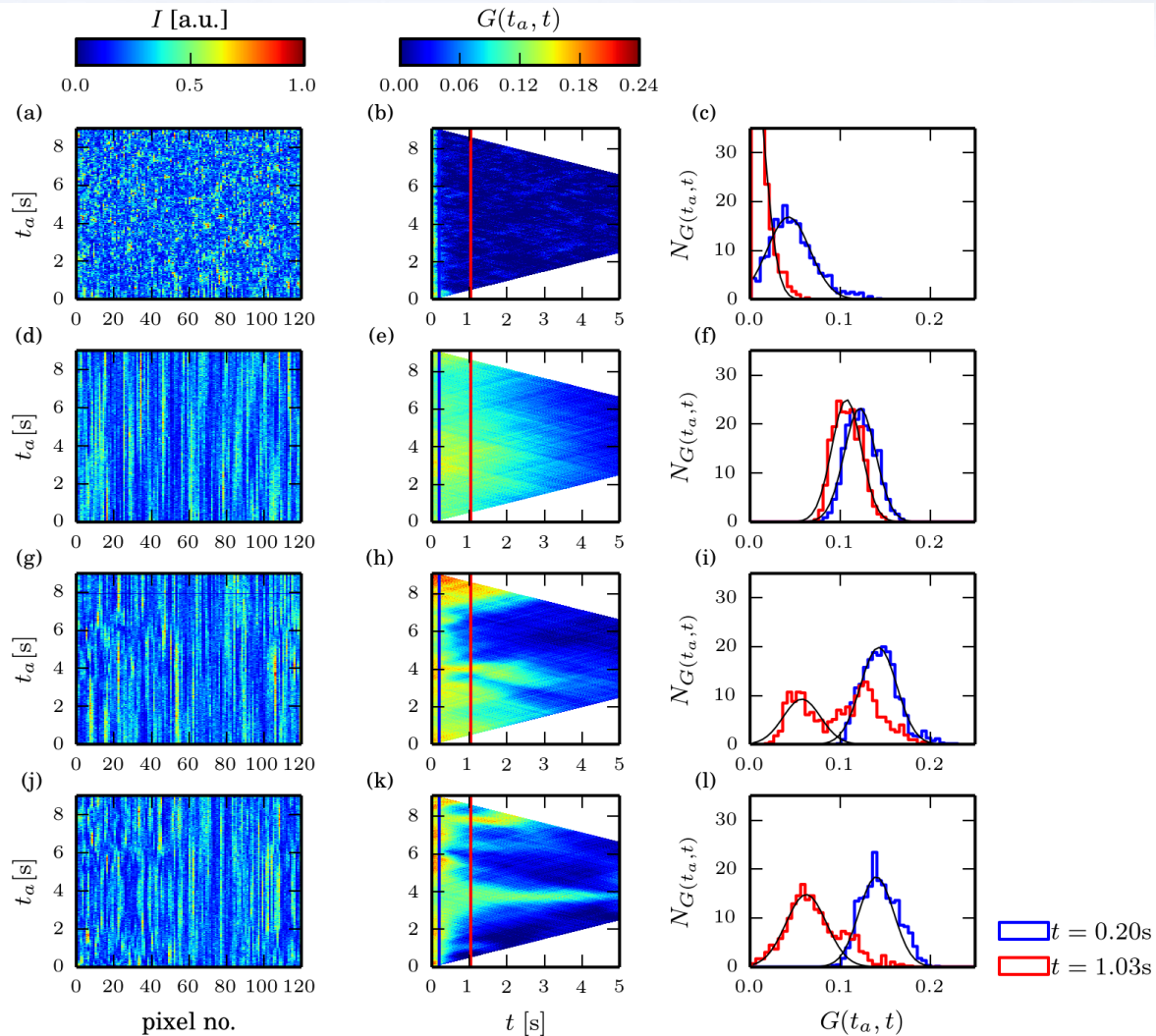
- Polymer-based sponge phases
P. Falus *et al.* *Phys. Rev. Lett* 2006
- Aging Clay (Laponite) Gels
B. Bandyopadhyay *et al.*, *Phys. Rev. Lett.* 2004;
R. Angelini *et al.*, *Soft Matter* 2013
- Antiferromagnetic domain fluctuations (Cr)
O. Shpyrko *et al.*, *Nature* 2007
- Aging Ferrofluids
A. Robert *et al.* *Europhys. Lett.* 2007
- Aging colloidal gels (“transient gels”)
A. Fluerasu *et al.*, *Phys. Rev. E* 2007
- Cross-linked Polymer Gels
R. Hernandez *et al.*, *J. Chem Phys* 2014
O. Czakkel, *Europhys. Lett.* 2011, K. Laszlo *et al.*, *Soft Matter* 2010
- Atomic-scale dynamics & aging in metallic glasses
B. Rutta *et al.*, *Phys. Rev. Lett.* 2012
- Etc. etc. etc. ...



Dynamical Heterogeneities

$\Phi \sim 0.57$

$\Phi \sim 0.61$



Age ~ 30 min

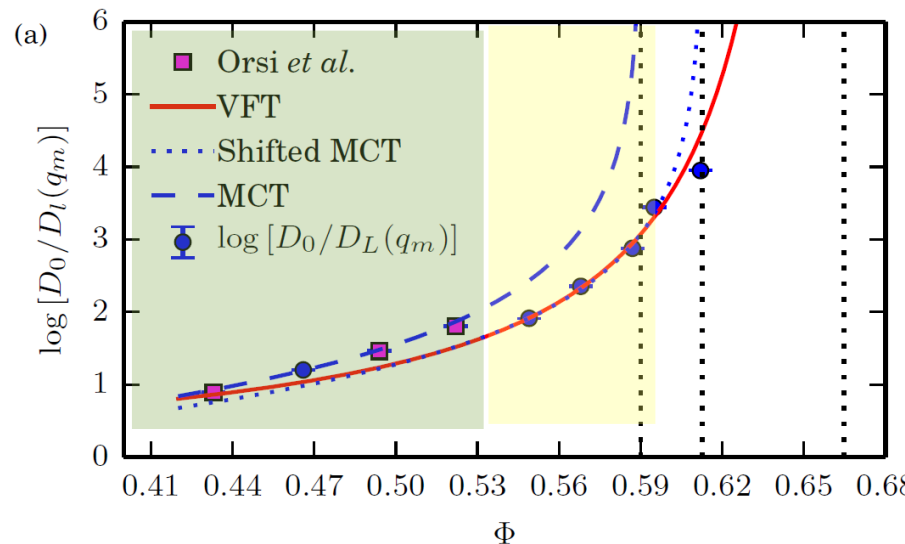
Age ~ 2 h30

Age ~ 9 h

Pawel Kwasniewski *et al.*

Colloidal Glasses: Conclusions

- Low- Φ : Dynamics of colloids well explained by existing many-body theories (MCT)
- $\Phi \geq 0.57$ -0.59 Stress in the network and stress-induced (nonthermal) fluctuations become dominant and hinder the expected glass transition
- Non-equilibrium, complex dynamics determined by “rough” energy landscape (heterogeneities)
Hyperdiffusive relaxations
→ *jamming*
(common also in other systems)
- Response to perturbations?
→ *flow, shear*



Acknowledgements

- Colloids** *Pawel Kwasniewski* (ESRF), *Davide Orsi* (U. Parma)
A. Madsen (XFEL)
- Proteins** *Luxi Li*, V. Stojanoff, L. Wiegart (BNL), S. Mochrie (Yale)
- CHX** *Lutz Wiegart, Yugang Zhang*,
M. Carlucci-Dayton, S. Antonelli, R. Greene,
D. Chabot, W. Lewis,
- Beamlines** ID 10 ESRF - Y. Chushkin, 34-ID APS - R. Harder
8-ID APS - A. Sandy, S. Narayanan
- NSLS-II** Ron Pindack, Qun Shen, P. Zschack, J. Hill, A. Broadbent
O. Chubar, K. Evans-Lutterodt, P. Siddons ...
- Funding** NSLS-II project: DOE# E-AC02-98CH10886
BNL SC0012704
BNL LDRD 11-025

