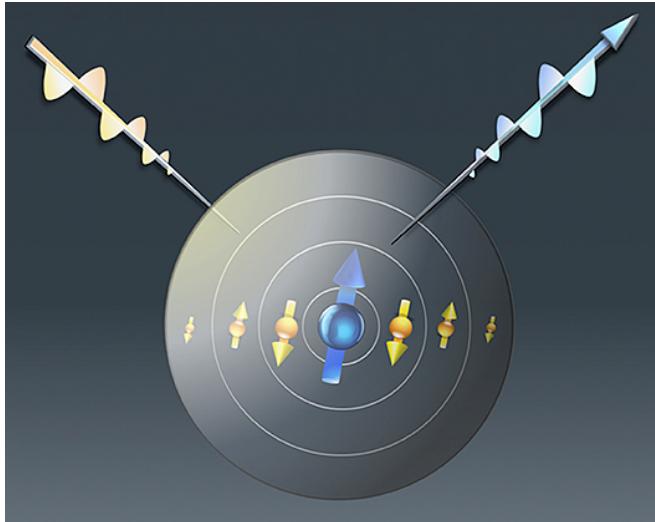
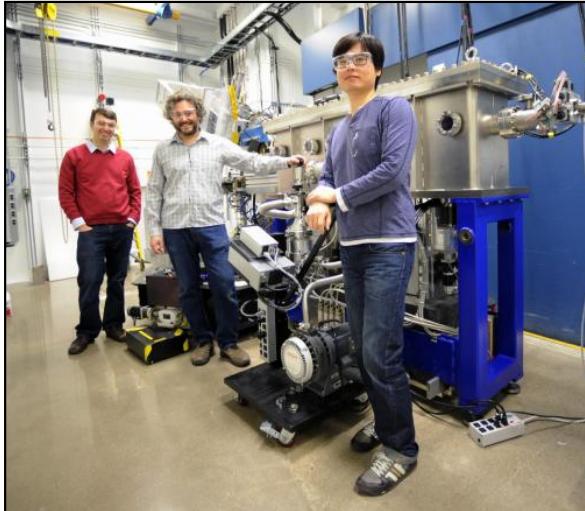


Introduction to Inelastic X-Ray Scattering

Jason N. Hancock

University of Connecticut

jason.hancock@uconn.edu



Neutron and X-ray Scattering School
Argonne National Laboratory, Argonne, IL

Outline

- Introductions
- Dispersion relations in materials physics
- Elastic waves and classification of phonons
- Phonon inelastic X-ray scattering examples
- Resonance phenomena in multi-electron atoms
- Resonant inelastic X-ray scattering examples



UConn Condensed Matter Physics



Alexander Balatsky
Theoretical physics
Superconductivity
Quantum materials



Gayanath Fernando
Theoretical physics
Electronic structure
Quantum materials



Boris Sinkovic
Thin films synthesis
Electron spectroscopy



Lea F. Dos Santos
Theoretical physics
Quantum chaos



Elena Dormidontova
Soft Matter Theory



Jason Hancock
THz, Infrared, X-ray
Applied physics



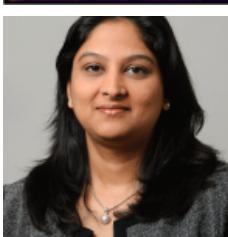
Ilya Sochnikov
Low-T transport
Scanning SQUID



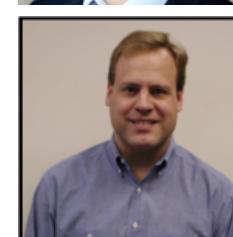
Pavel Volkov
(Harvard until 2023)
Materials theory
Twistrionics



Niloy Dutta
Photonics
Applied physics



Menka Jain
Thin film synthesis



Barrett Wells
PLD films, Muons,
Neutrons, ARPES

Plus strong connections with: AMO group, UConn Institute for Materials Science, New UConn Tech Park

My PhD research does/will likely use...

A

B

C

D

E

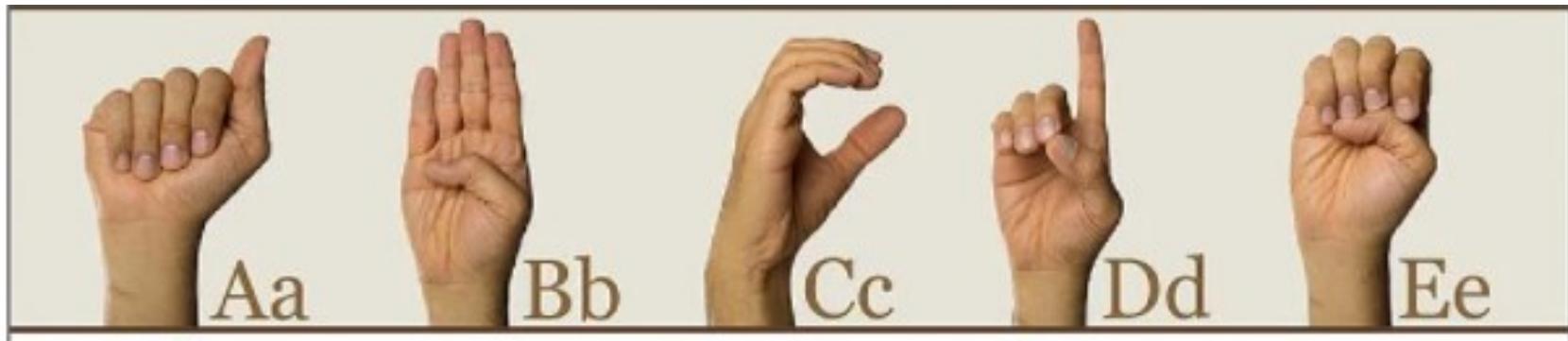
X-rays

Neutrons

Both X-rays and
Neutrons

Neither

I am a theorist



My PhD research does/will likely regard...

A

B

C

D

E

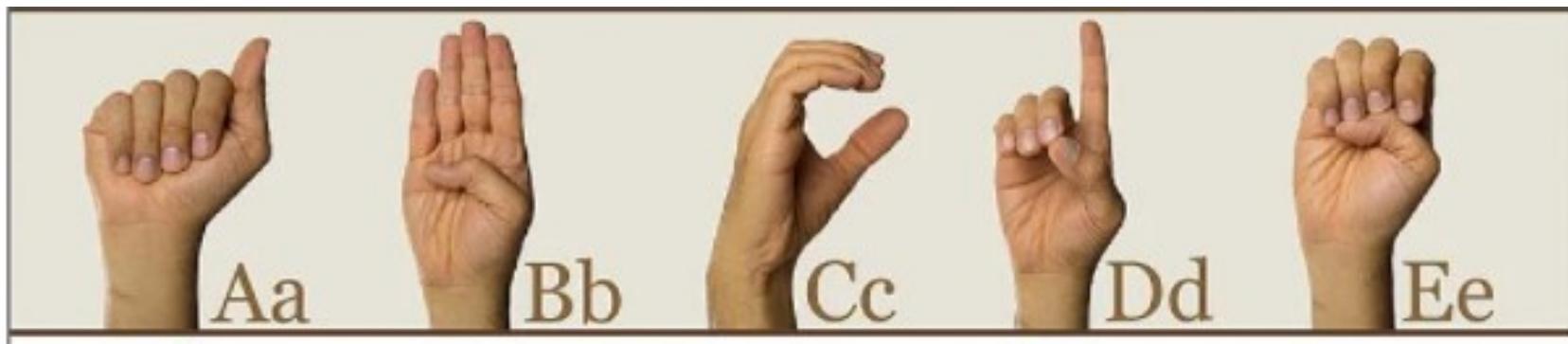
Mostly structural
studies

Spectroscopy and
electronic structure

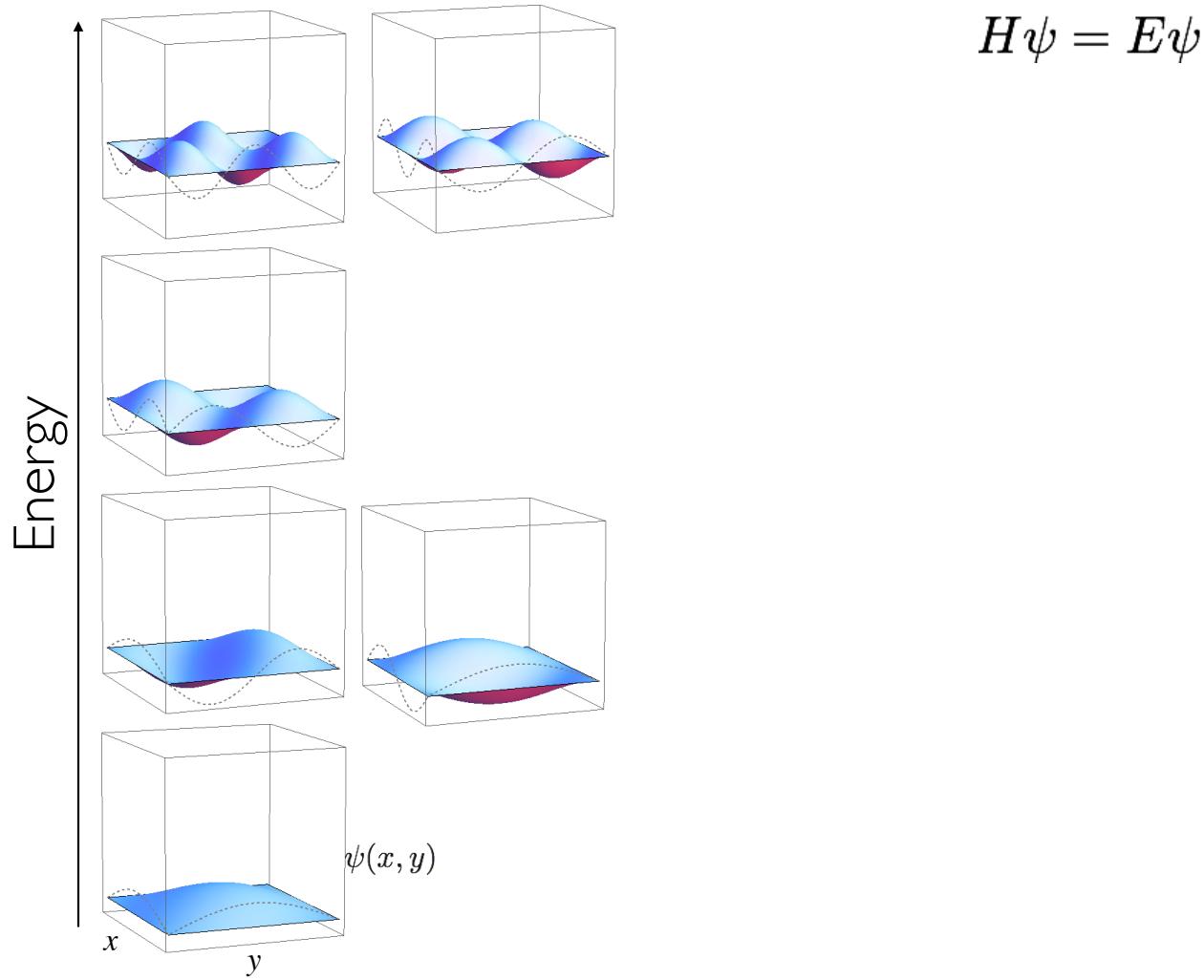
Advanced imaging

Still figuring it out

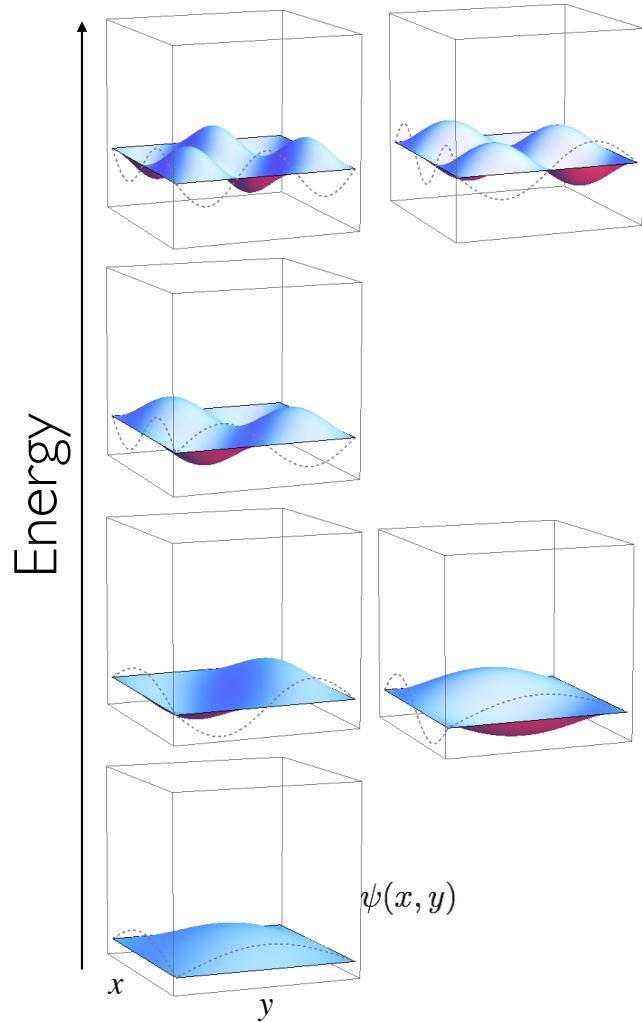
None of these



Particle-in-a-box: the simplest quantum problem



Particle-in-a-box: the simplest quantum problem

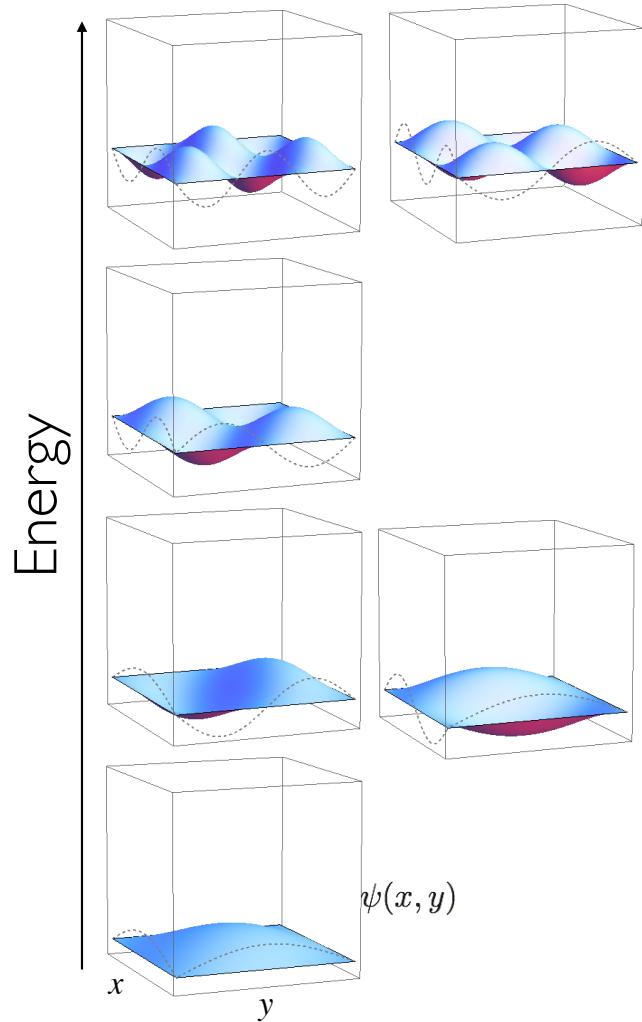


$$H\psi = E\psi$$

$$H = \frac{p^2}{2m} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$$

Kinetic
energy

Particle-in-a-box: the simplest quantum problem

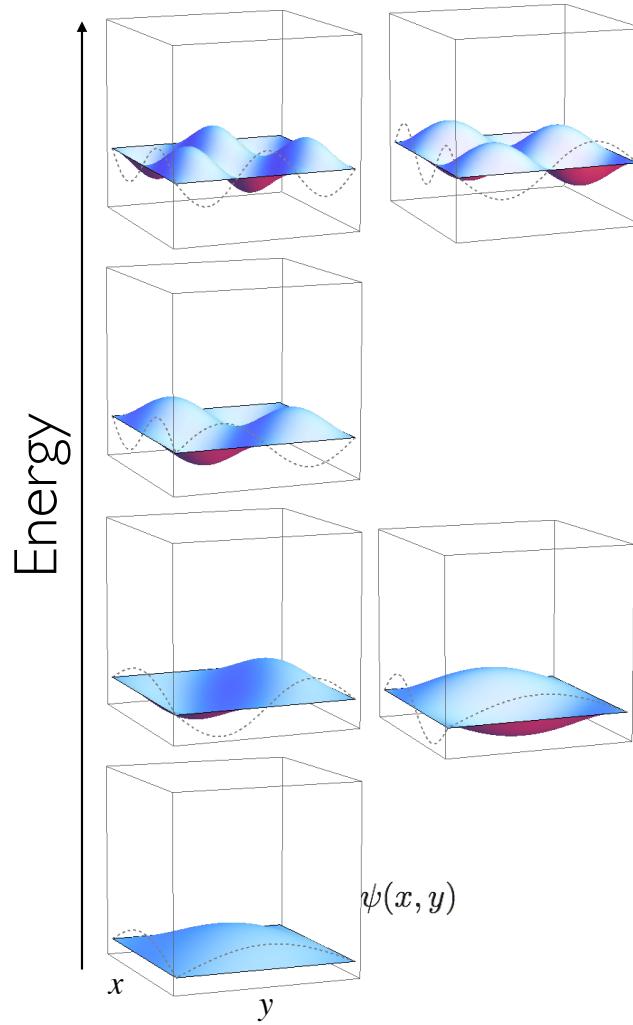


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$$E = \frac{\hbar^2 k^2}{2m}$$

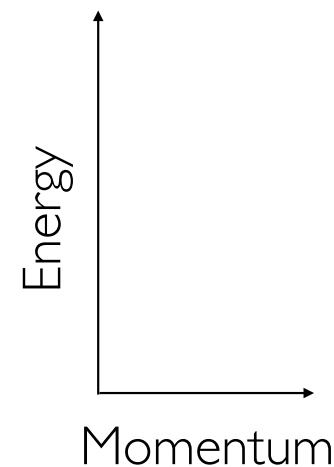
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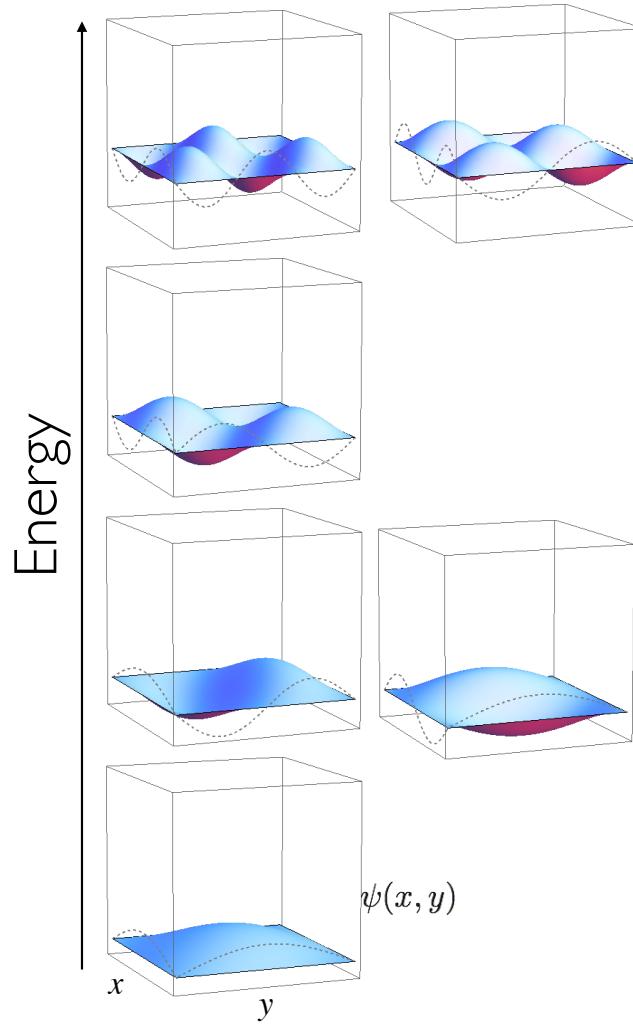
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Kinetic energy

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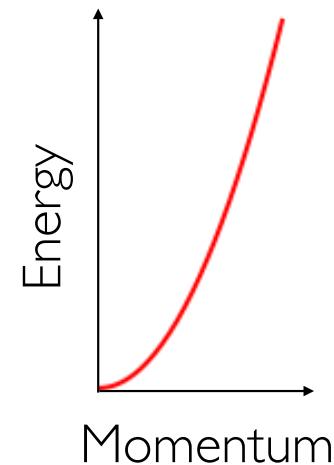
Particle-in-a-box: the simplest quantum problem



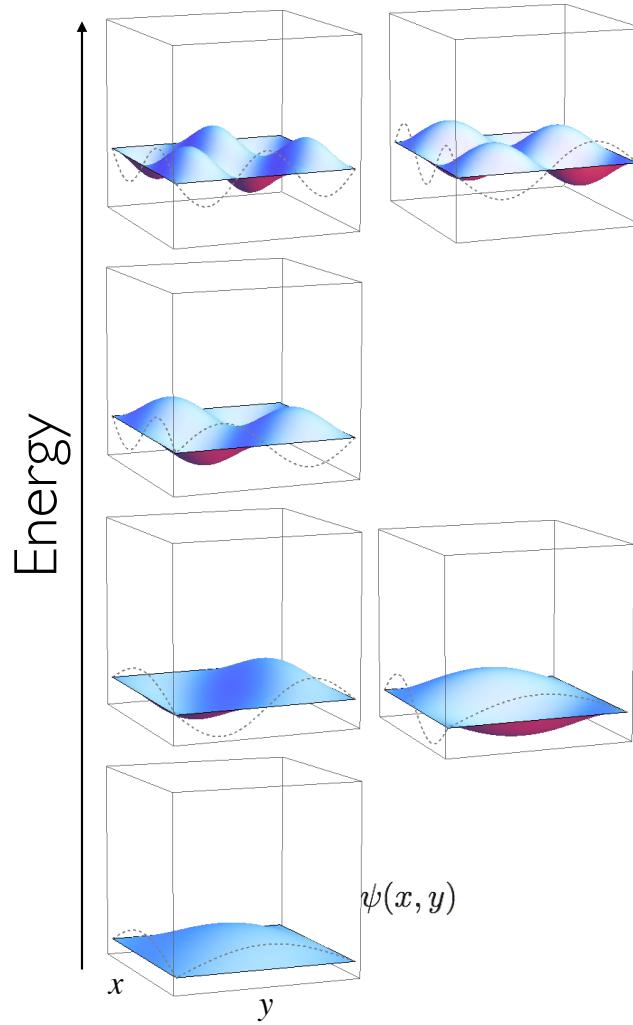
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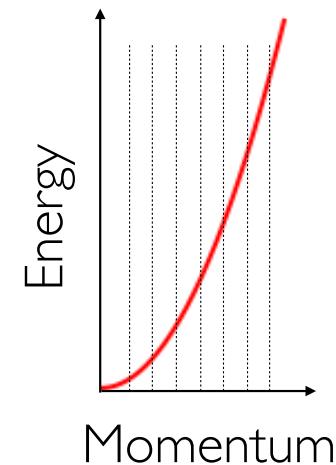
Particle-in-a-box: the simplest quantum problem



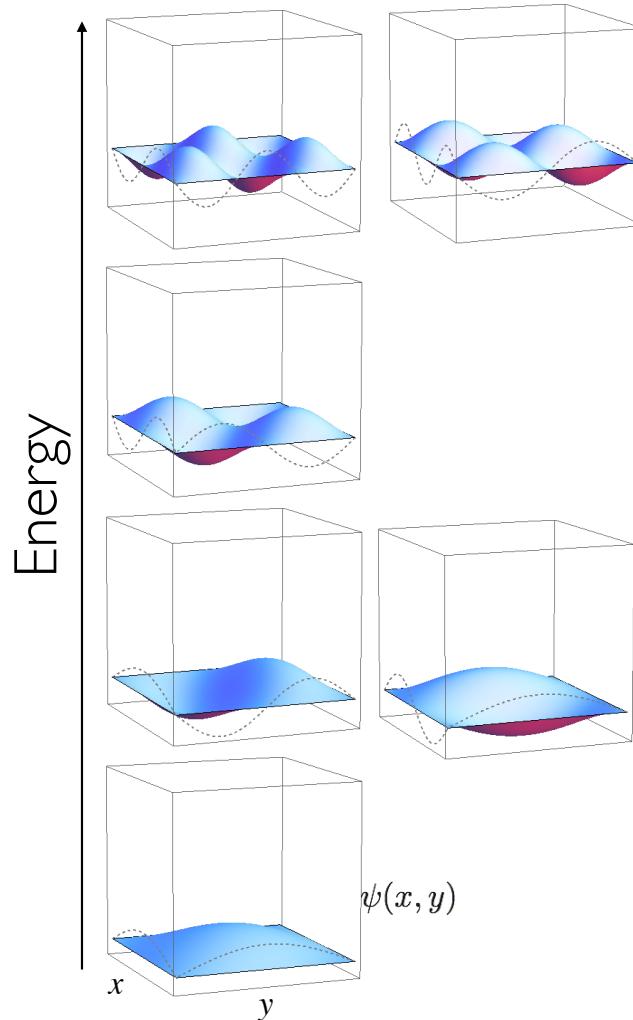
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Kinetic energy

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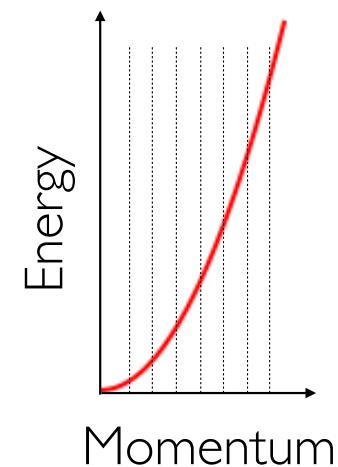
Particle-in-a-box: the simplest quantum problem



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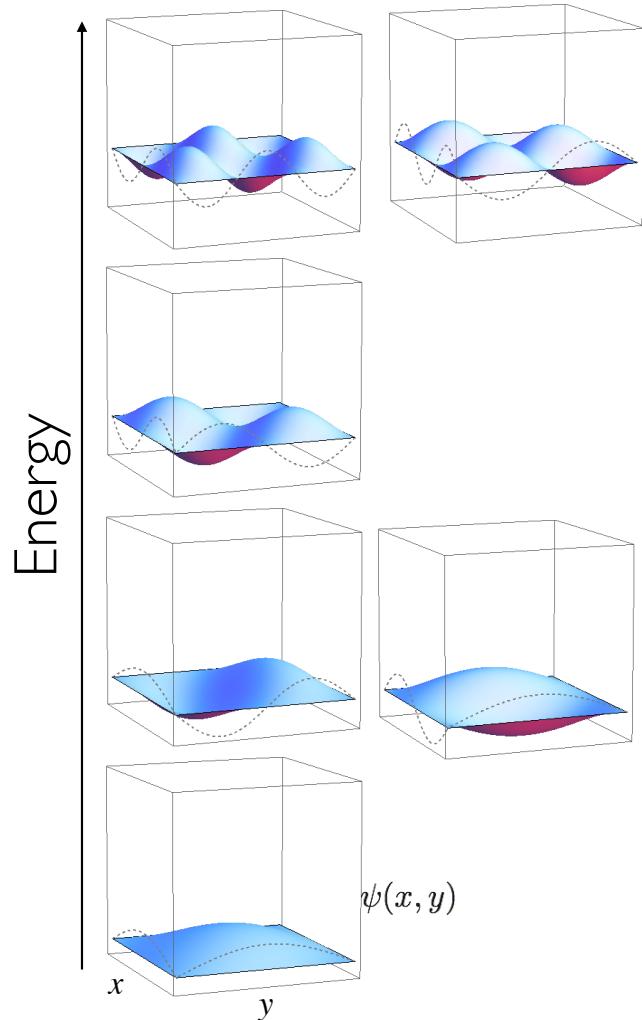
Kinetic energy

$$E = \frac{\hbar^2 k^2}{2m}$$



Dispersion relation

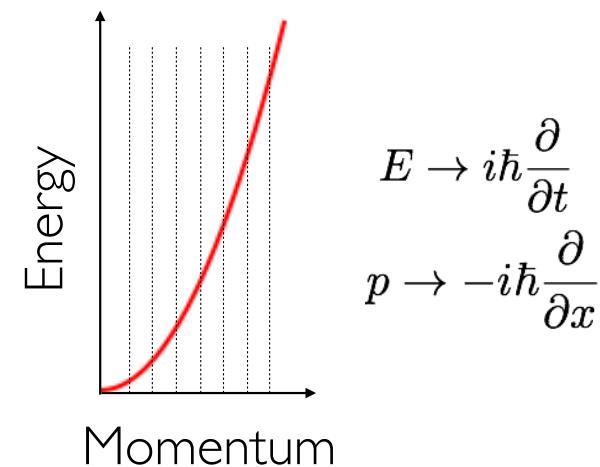
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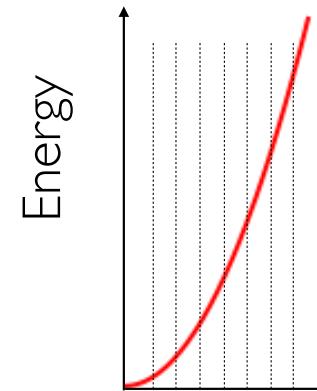
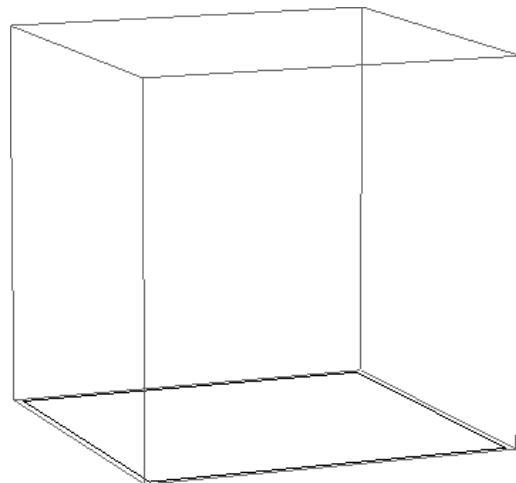


Dispersion relation

Particle in a box to the band theory of solids

$$H\psi = E\psi$$

$$H = \frac{p^2}{2m} + V(\vec{r})$$

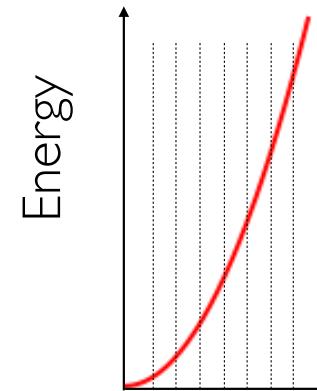
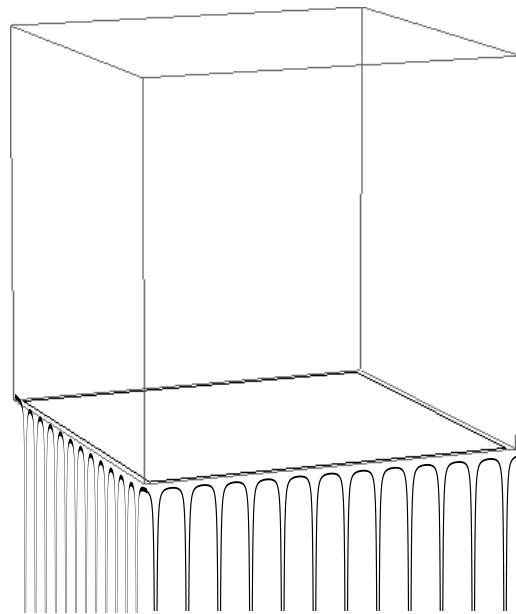


Momentum

Particle in a box to the band theory of solids

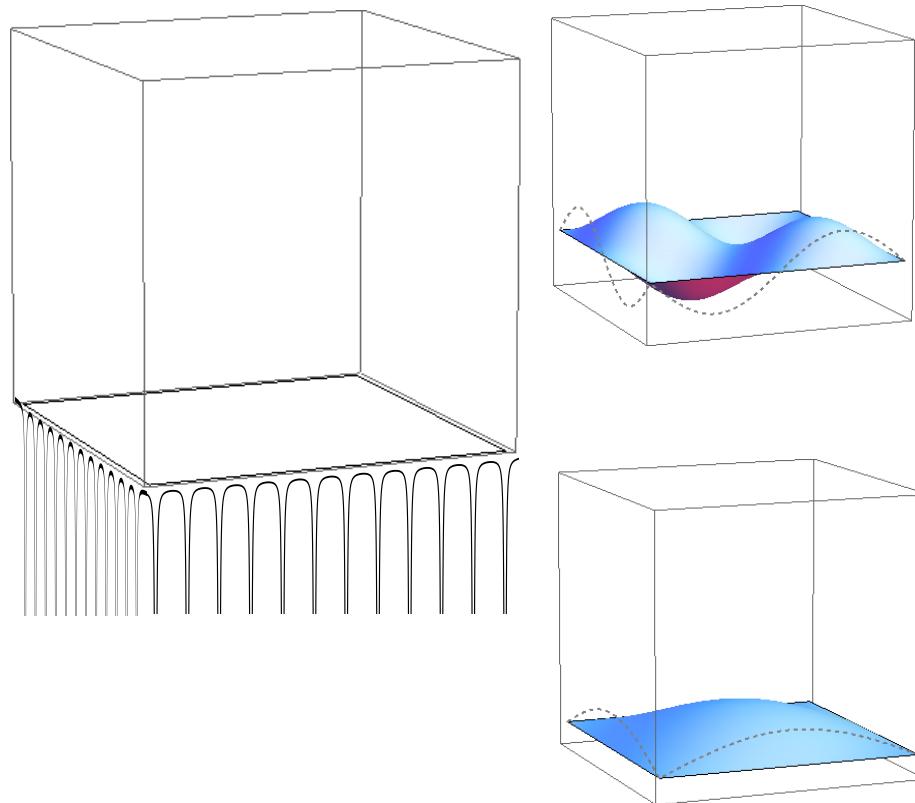
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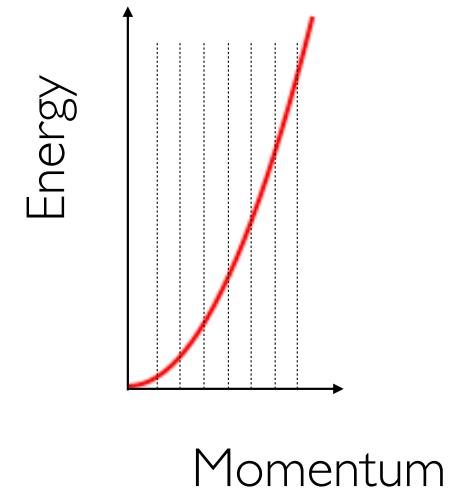
Momentum

Particle in a box to the band theory of solids

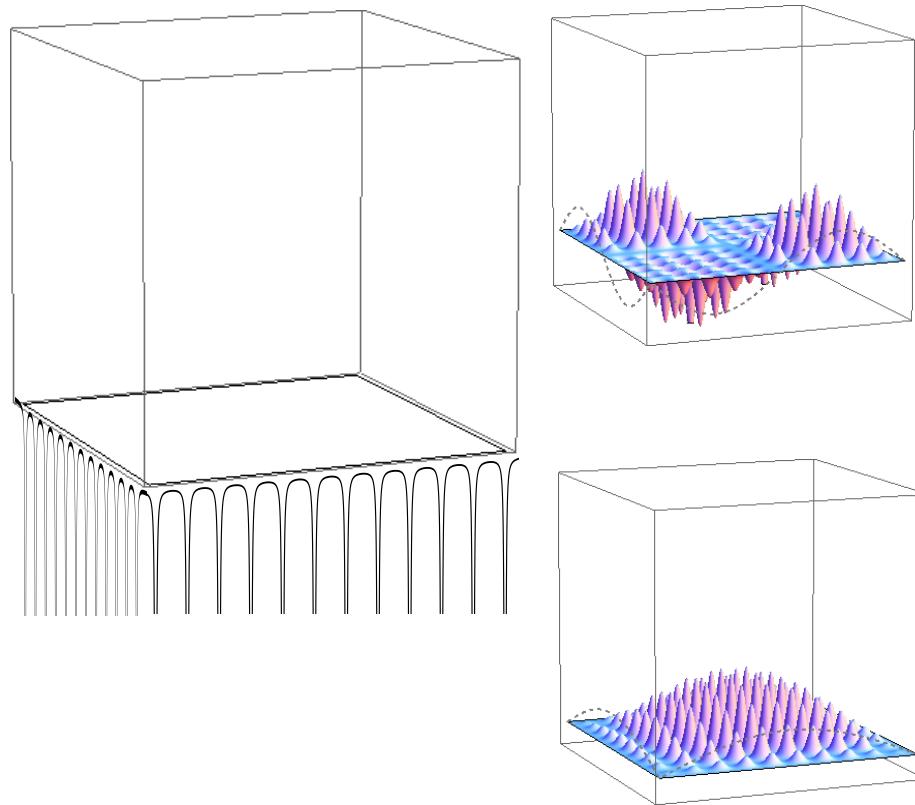


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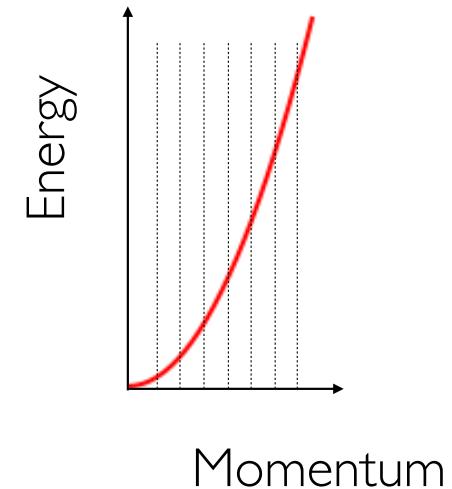


Particle in a box to the band theory of solids

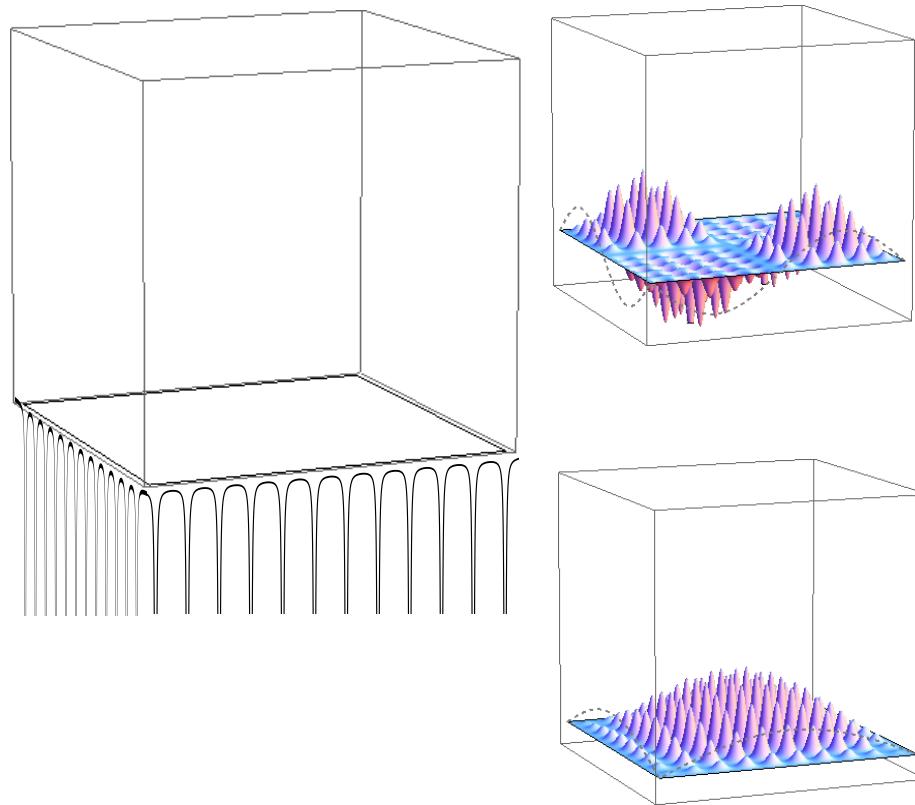


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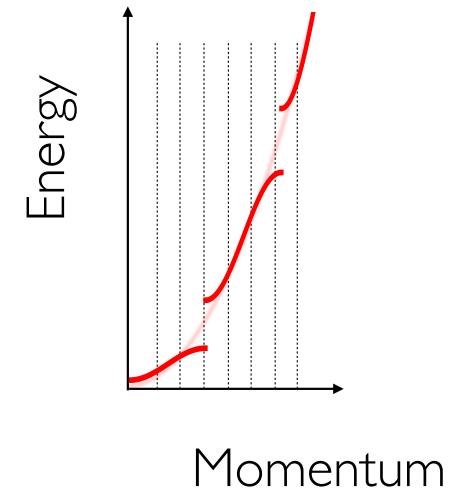


Particle in a box to the band theory of solids

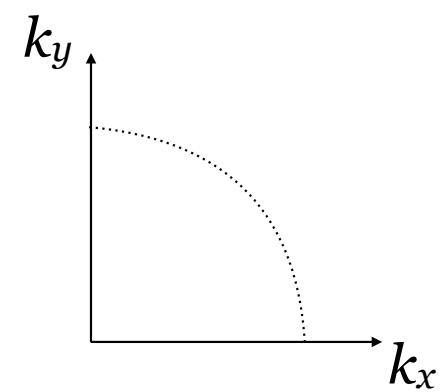
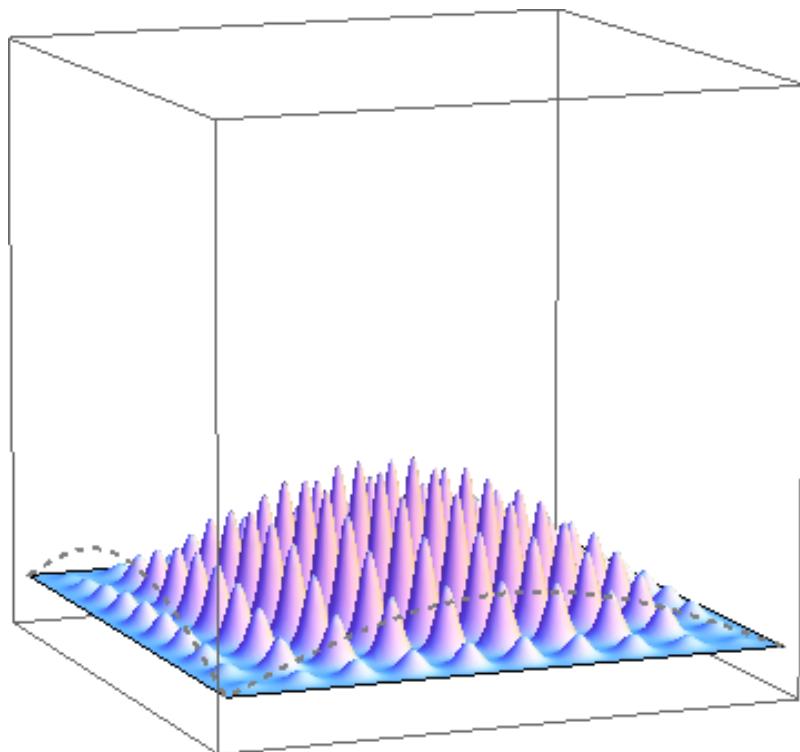


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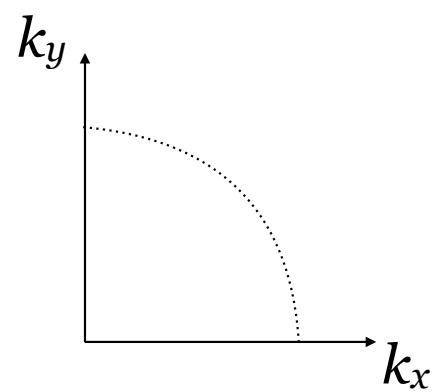
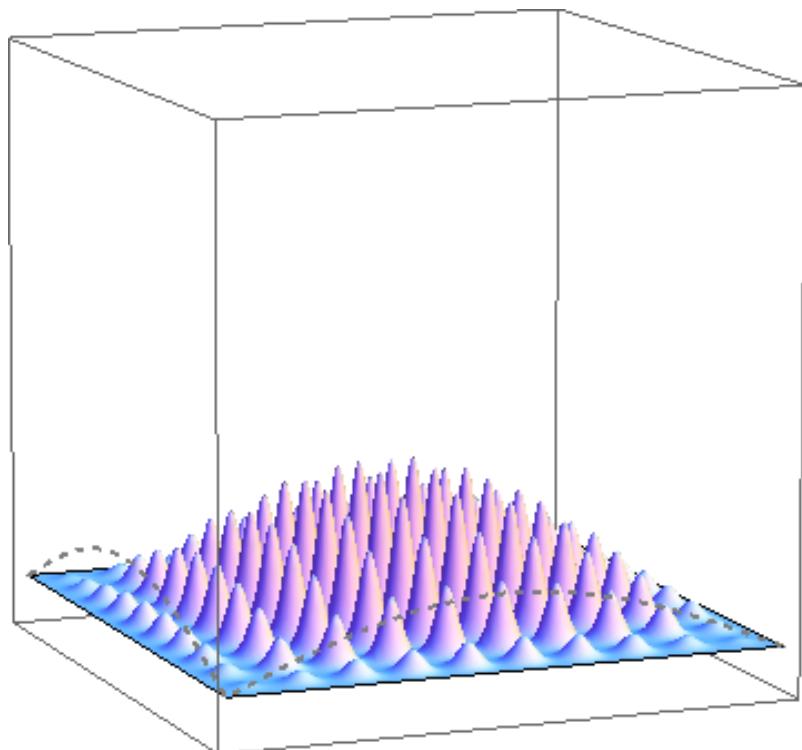
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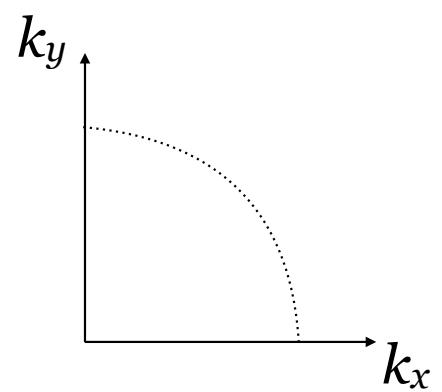
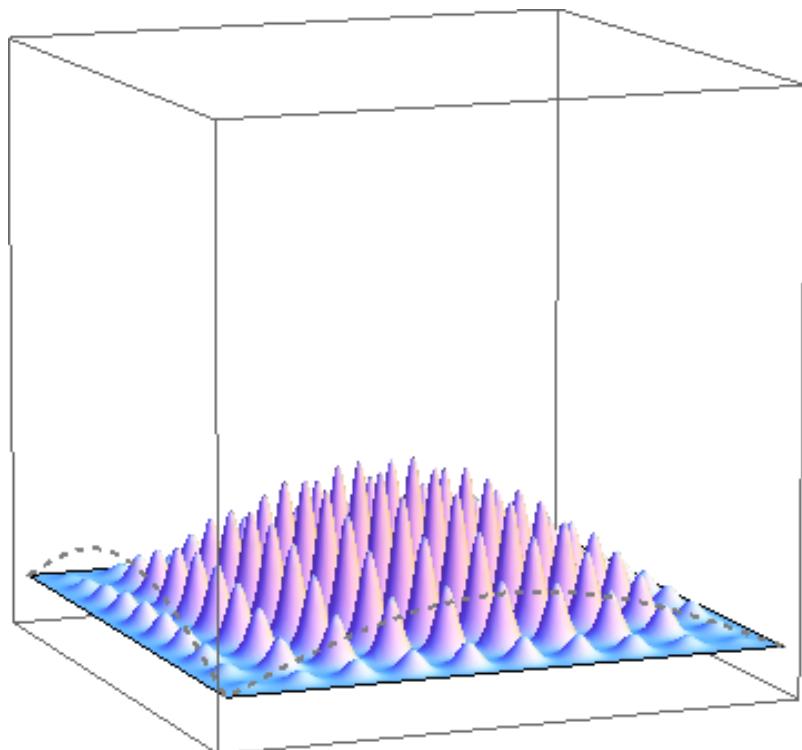
Many electron states



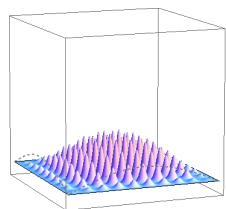
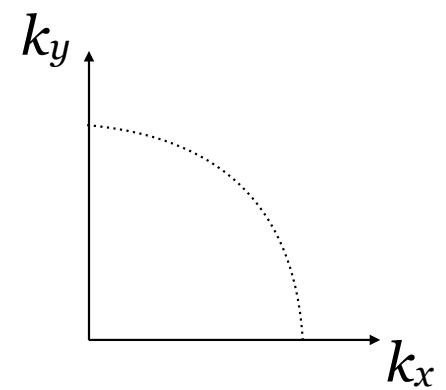
Many electron states



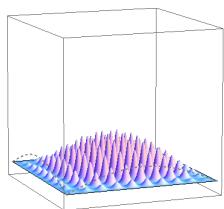
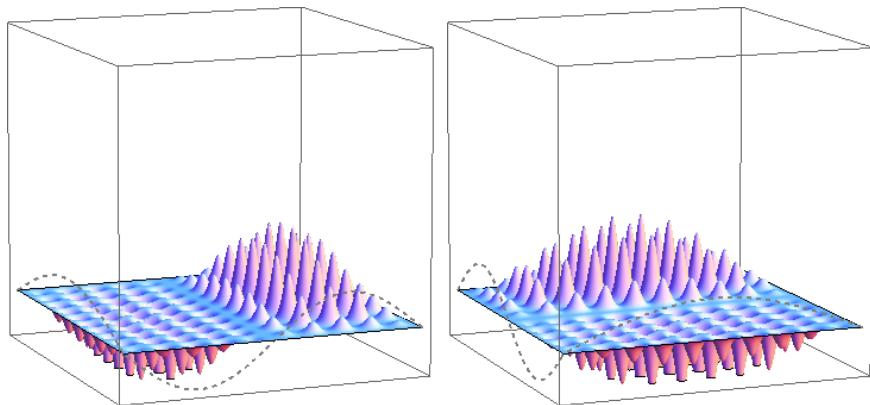
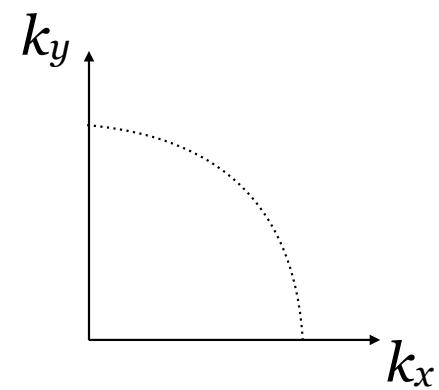
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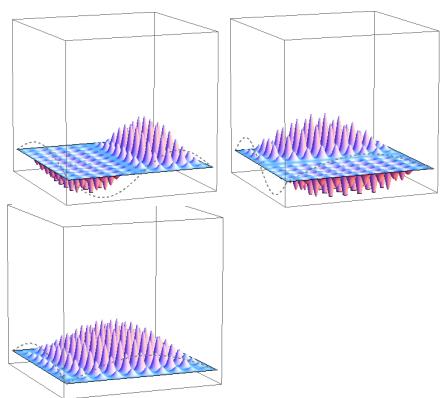
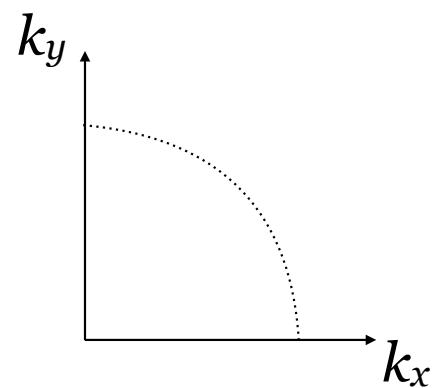
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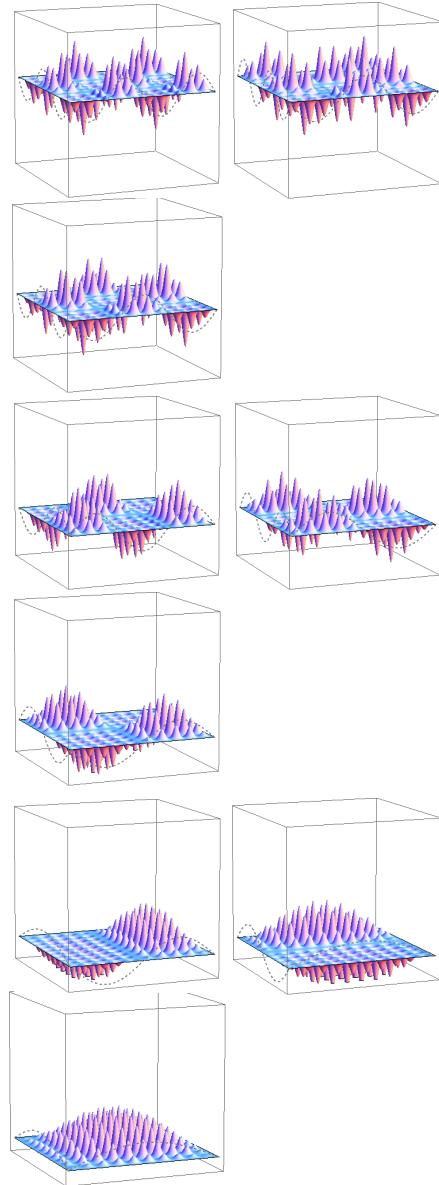


Many electron states

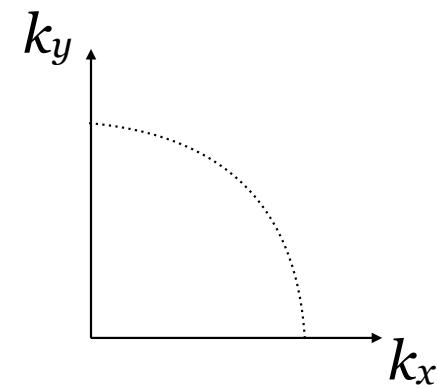


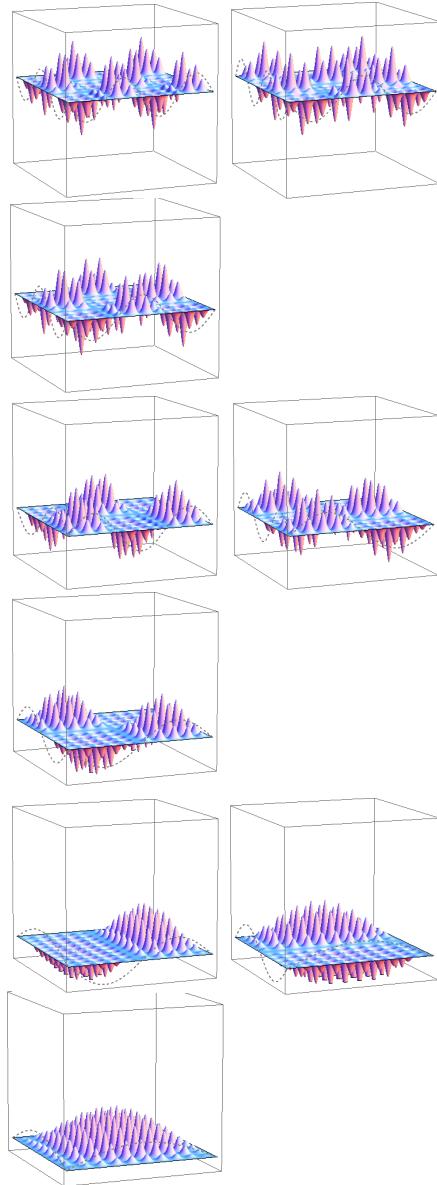
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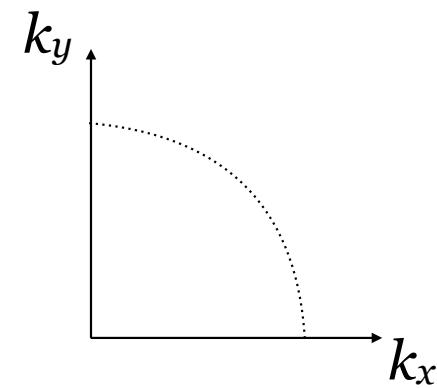


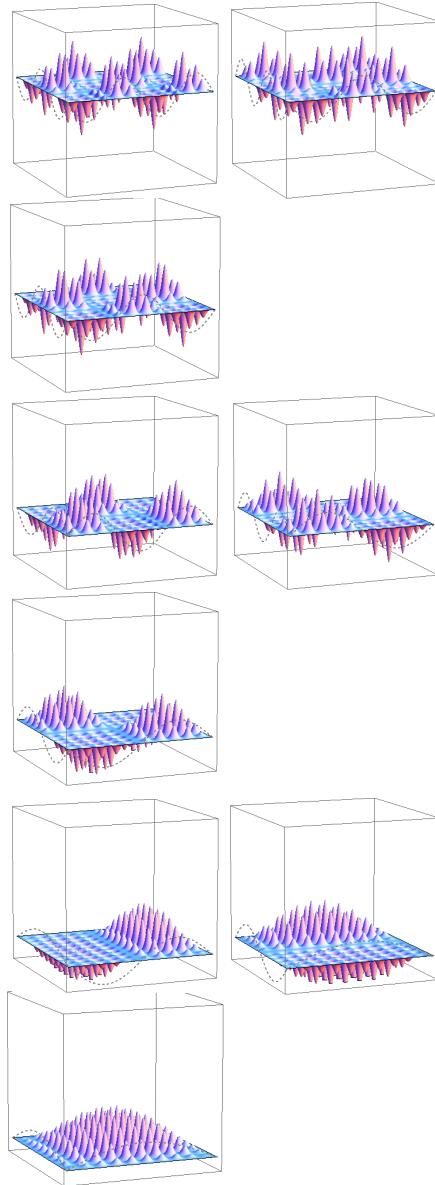
Many electron states



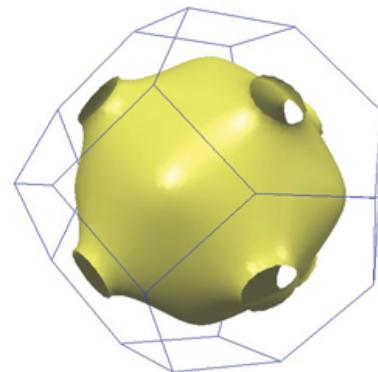


Many electron states

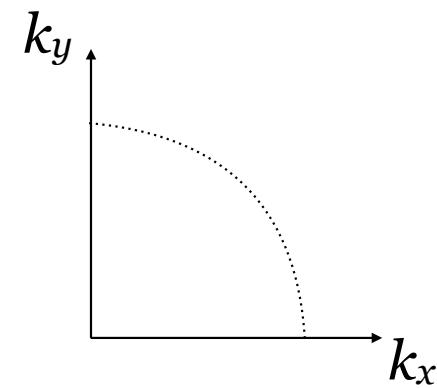


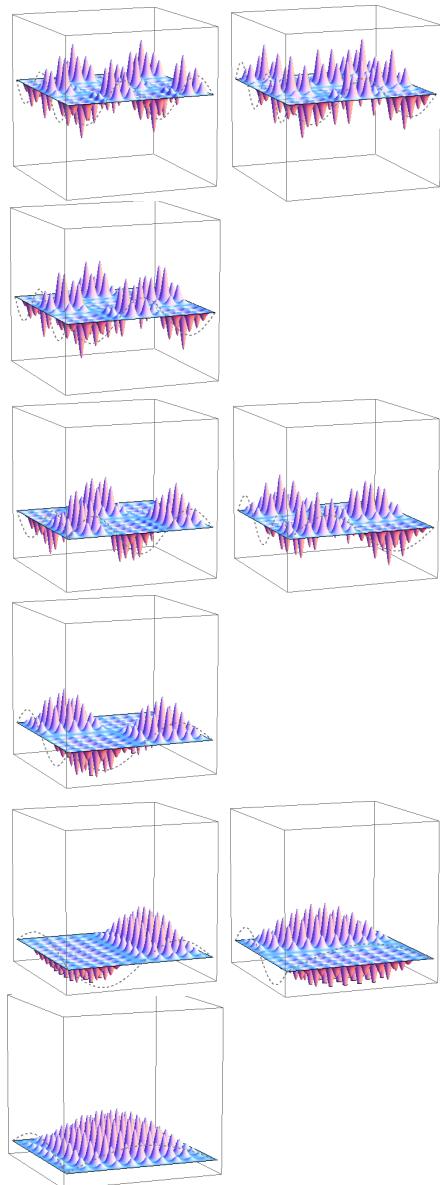


Many electron states

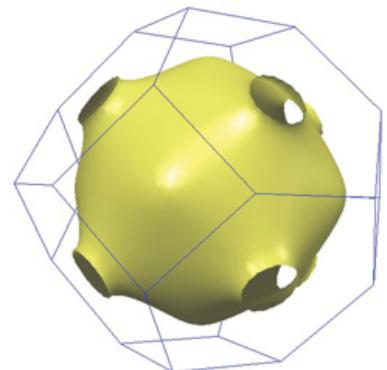


Fermi surface of Cu

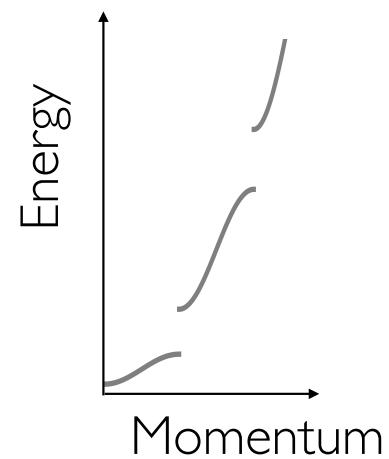
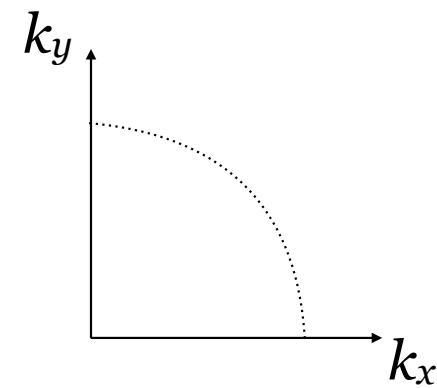


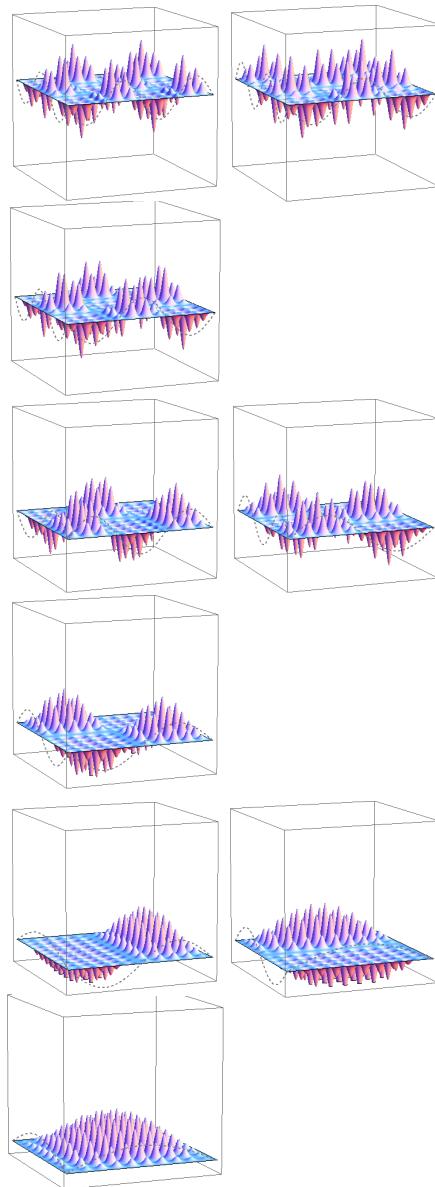


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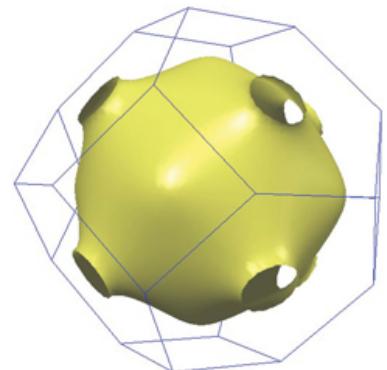


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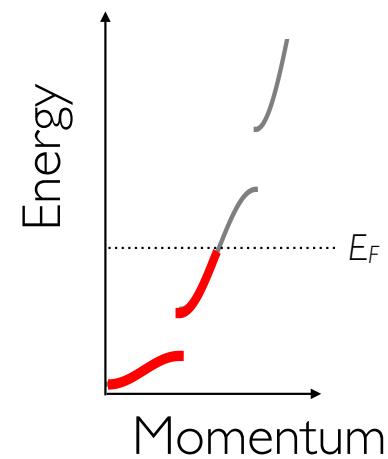
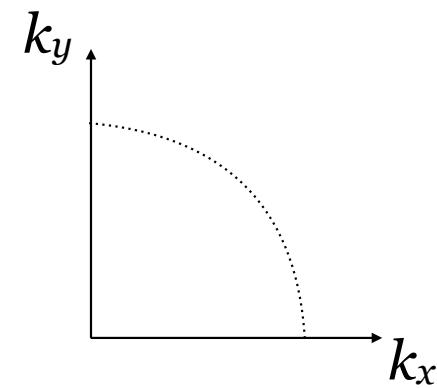


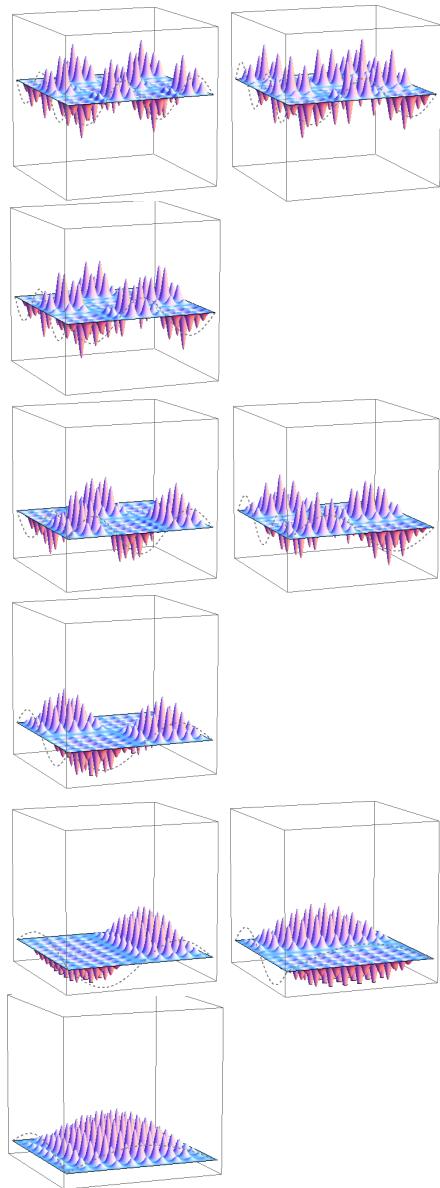


Many electron states

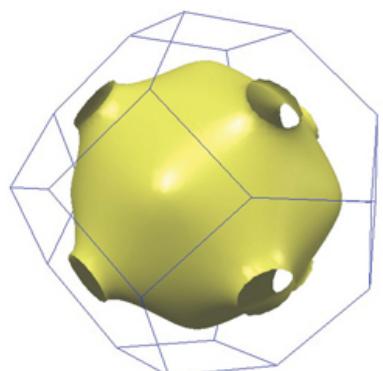


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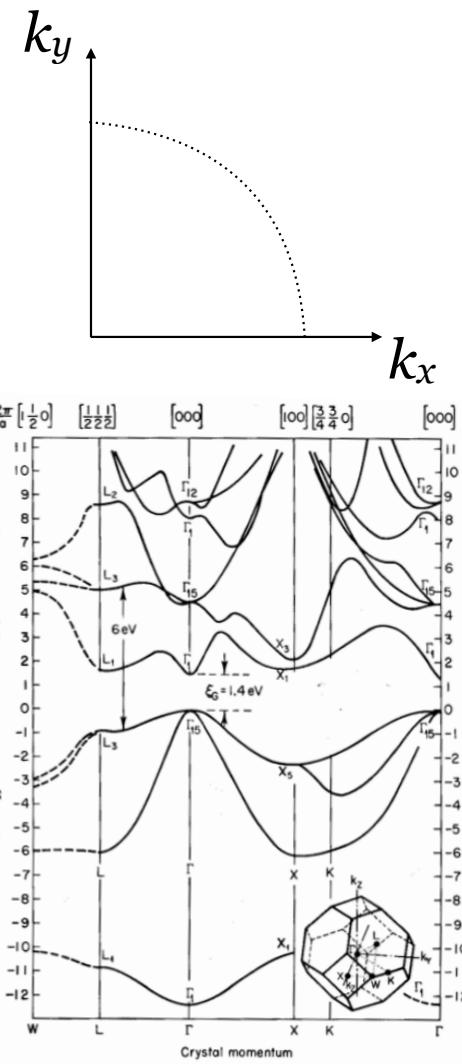
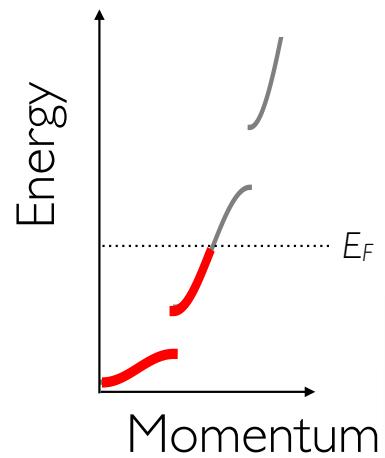




Many electron states



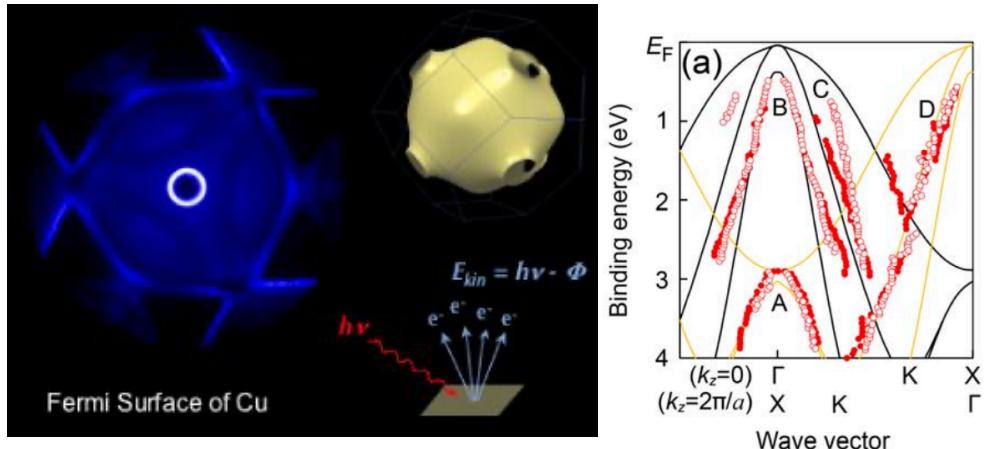
Fermi surface of Cu



Uncorrelated electron systems

Electron-electron interactions do not dominate material behavior

Well-established theoretical framework describes most properties of semiconductors (Si,Ge,GaAs), good metals (Cu,Ag,Al,Au)



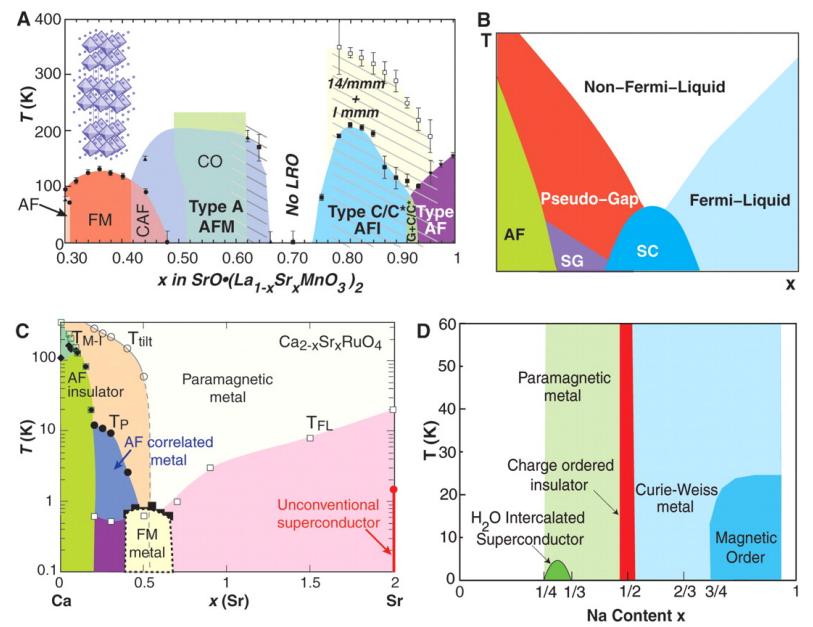
Correlated electron systems

Electron-electron interactions dominate material behavior

Often includes magnetism, more exotic behavior

Theoretical treatment is very limited

High potential for applications and interesting basic science

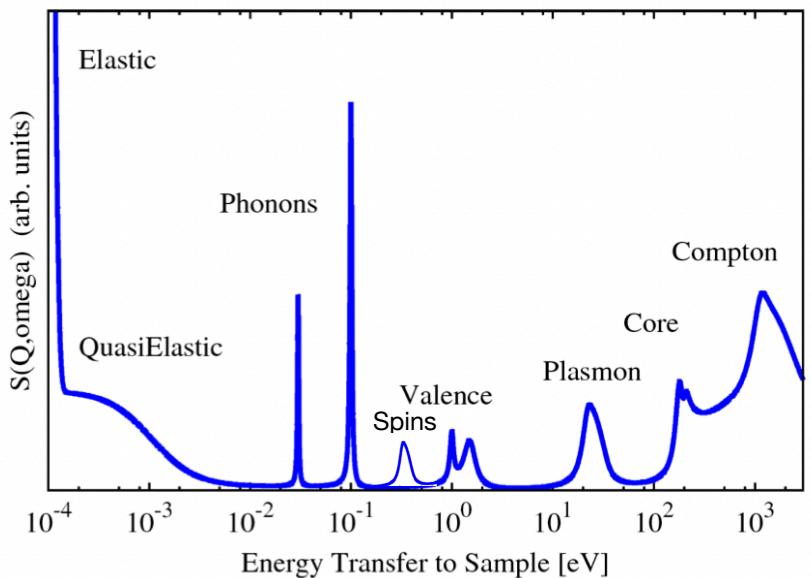


Collective excitations in materials



- “Fields” of spin, charge, nuclear displacement self organize according to interactions and confining potential
- Disturbances in the pattern form excited states relevant to material behavior

Inelastic X-ray scattering can probe many different fundamental excitations



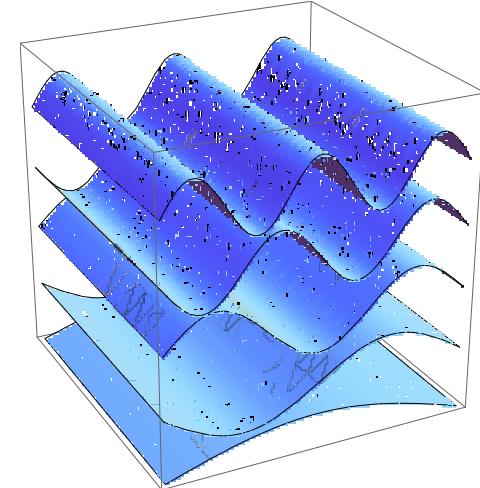
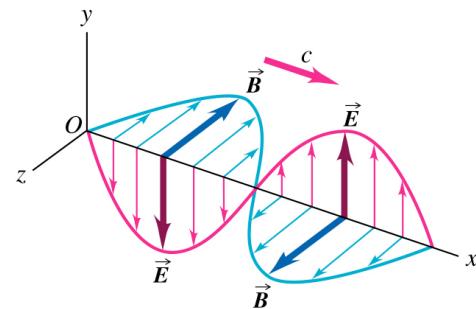
The Electromagnetic Field

- Photons are bosons, independent, luminal
- Described by Maxwell's equations at low energy, QED at high energy
- Fully understood, quantized, modes are described as

$$|v\rangle = |n_0, \dots n_{\mathbf{k}} \dots\rangle$$

Vector potential:

$$\mathbf{A}(\mathbf{r}, t) = \sum_{\mathbf{k}} \sum_{\mu=-1,1} \left(\mathbf{e}^{(\mu)}(\mathbf{k}) a_{\mathbf{k}}^{(\mu)}(t) e^{i\mathbf{k}\cdot\mathbf{r}} + \bar{\mathbf{e}}^{(\mu)}(\mathbf{k}) \bar{a}_{\mathbf{k}}^{(\mu)}(t) e^{-i\mathbf{k}\cdot\mathbf{r}} \right)$$

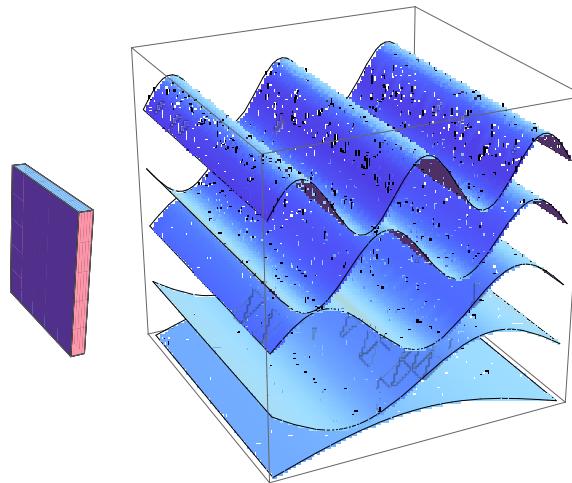


Spectroscopy, general remarks

- Sample can exchange energy, momentum with EM field

$$|\psi_g\rangle \otimes |v\rangle \rightarrow |\psi_e\rangle \otimes |v'\rangle$$

(sample) (vacuum)

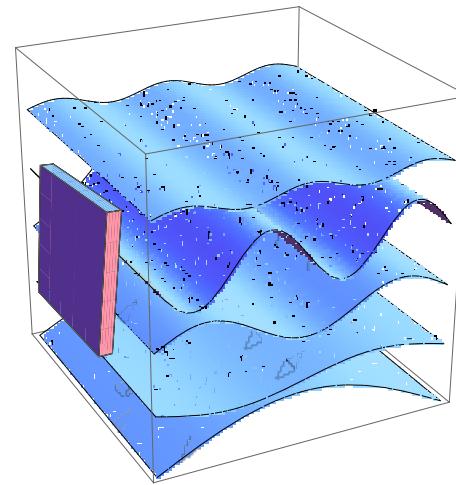


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Spectroscopy, general remarks

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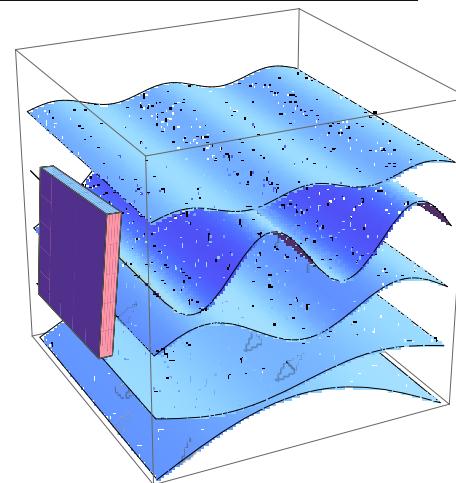
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(sample) (vacuum)



One photon
Optical spectroscopy
X-ray absorption

$$|n_0, \dots n_{\mathbf{k}} \dots \rangle \rightarrow |n_0, \dots n_{\mathbf{k}} - 1 \dots \rangle$$



Spectroscopy, general remarks

- Sample can exchange energy, momentum with EM field

$$|\psi_g\rangle \otimes |v\rangle \rightarrow |\psi_e\rangle \otimes |v'\rangle$$

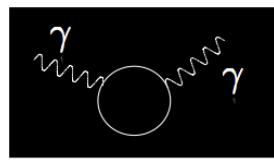
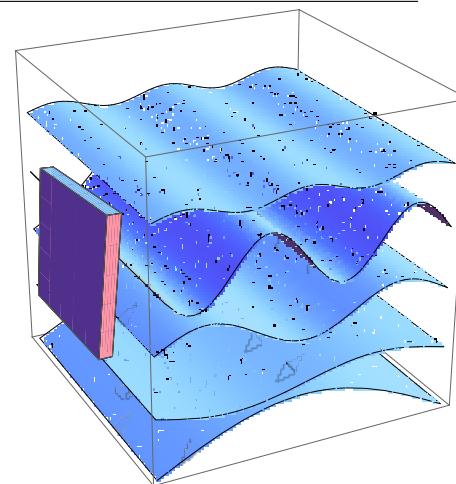
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One photon

Optical spectroscopy
X-ray absorption

$$|n_0, \dots n_{\mathbf{k}} \dots \rangle \rightarrow |n_0, \dots n_{\mathbf{k}} - 1 \dots \rangle$$



Two photon Raman Inelastic X-ray Scattering

$$|n_0, \dots n_{\mathbf{k}} \dots \rangle \rightarrow |n_0, \dots n_{\mathbf{k}} - 1 \dots n_{\mathbf{k}}, +1 \dots \rangle$$

Spectroscopy, general remarks

- Sample can exchange energy, momentum with EM field

$$|\psi_g\rangle \otimes |v\rangle \rightarrow |\psi_e\rangle \otimes |v'\rangle$$

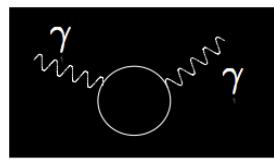
(sample) (vacuum)



One photon

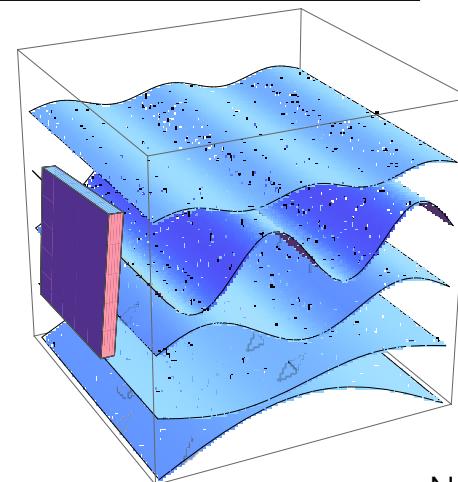
Optical spectroscopy
X-ray absorption

$$|n_0, \dots n_{\mathbf{k}} \dots \rangle \rightarrow |n_0, \dots n_{\mathbf{k}} - 1 \dots \rangle$$



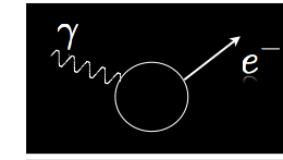
Two photon Raman Inelastic X-ray Scattering

$$|n_0, \dots n_{\mathbf{k}} \dots \rangle \rightarrow |n_0, \dots n_{\mathbf{k}} - 1 \dots n_{\mathbf{k}}, +1 \dots \rangle$$

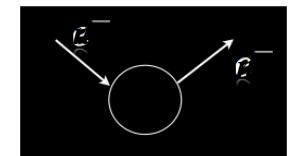


Related spectroscopies:

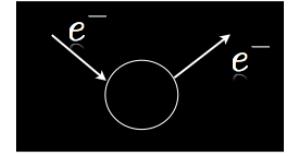
ARPES



Neutron



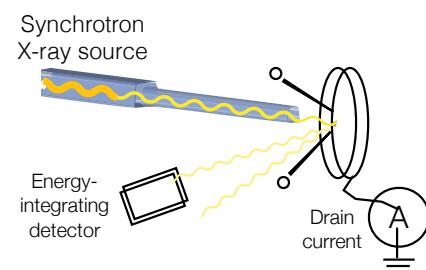
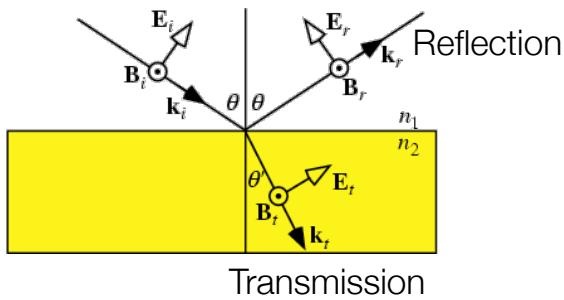
EELS



Photon spectroscopies

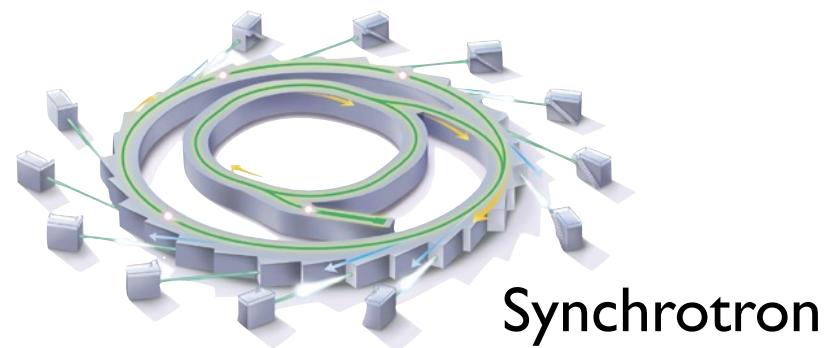
One-photon

- Can measure from reflection or transmission in time or frequency domain
- Use Fresnel's equations to determine $n + ik$, $\sigma_1 - i\sigma_2, \epsilon_1 + i\epsilon_2$ and relate to fundamental behavior
- X-ray absorption collected differently, through total electron or total fluorescence yield



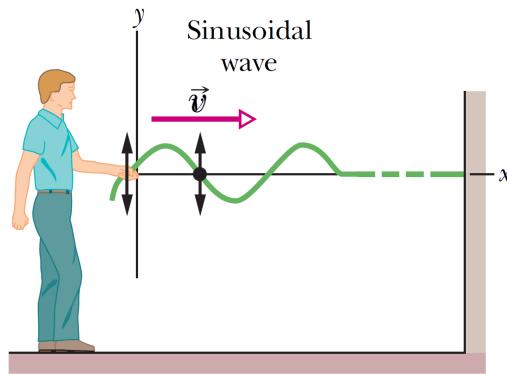
Two-photon

- IXS and RIXS (also REXS, Raman)
- Probes electronic excitations in a momentum-resolved way
- Technical advances making rapid progress



Elastic waves

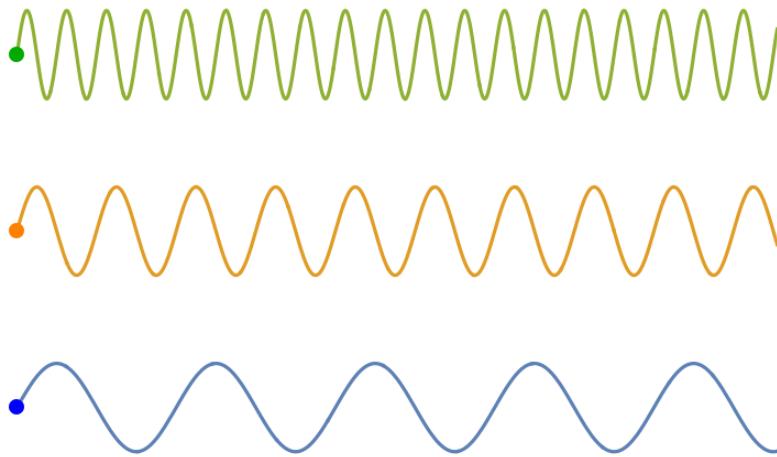
- Materials can distort in wavelike patterns, typical of transverse waves on a string



A common wave equation: $\frac{\partial^2 f}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 f}{\partial t^2}$

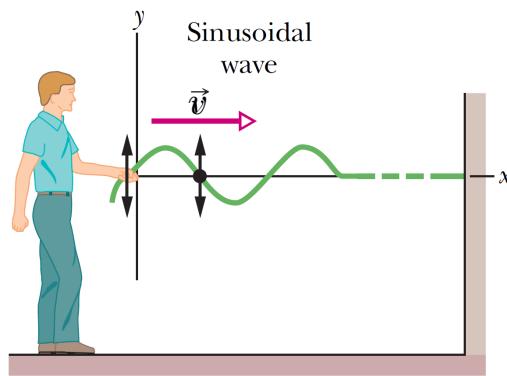
permits solutions of the form: $f(x - vt)$

All waves move at the same speed, regardless of wavelength



Elastic waves

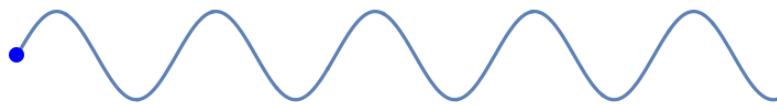
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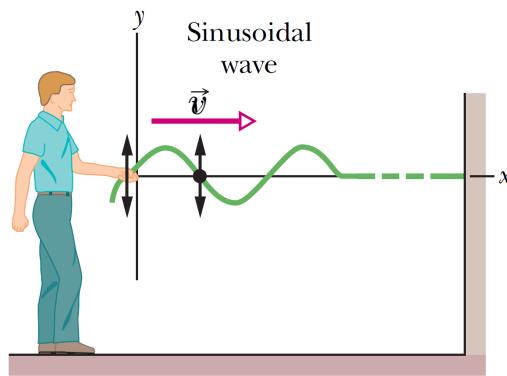
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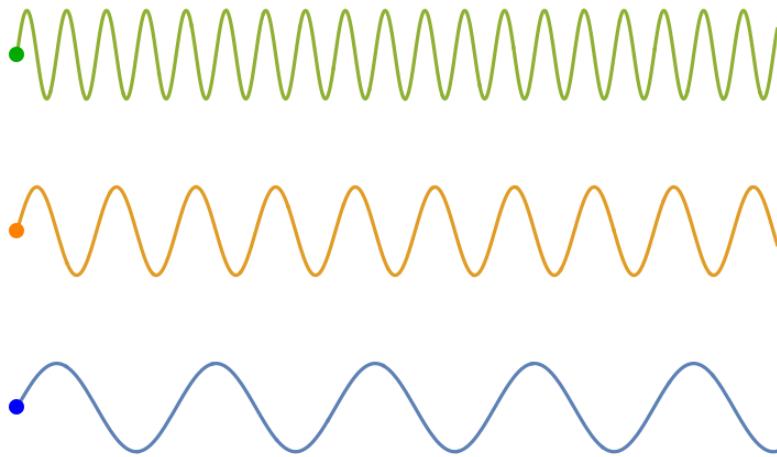
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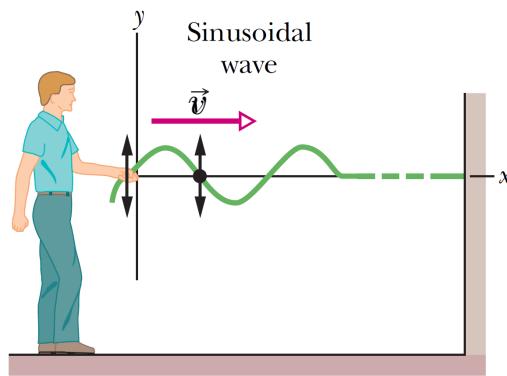
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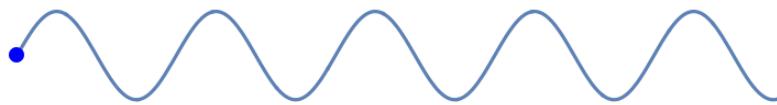
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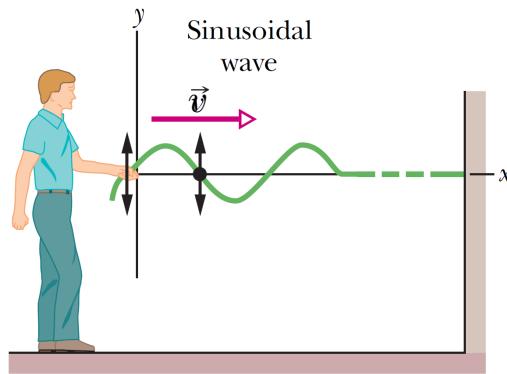
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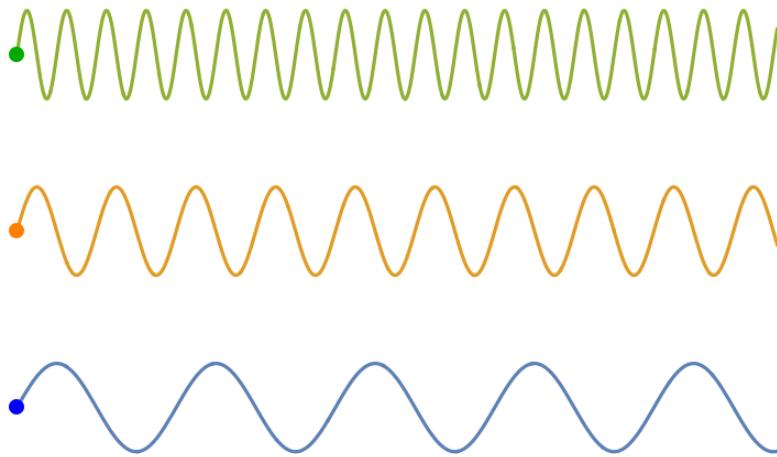
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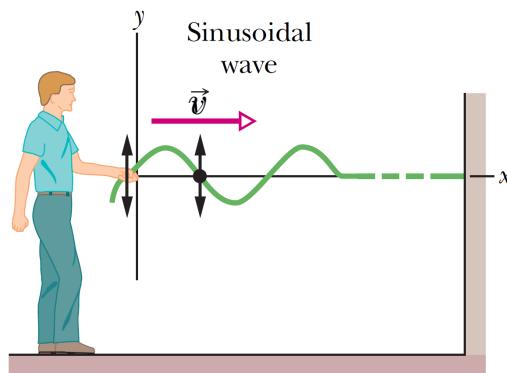
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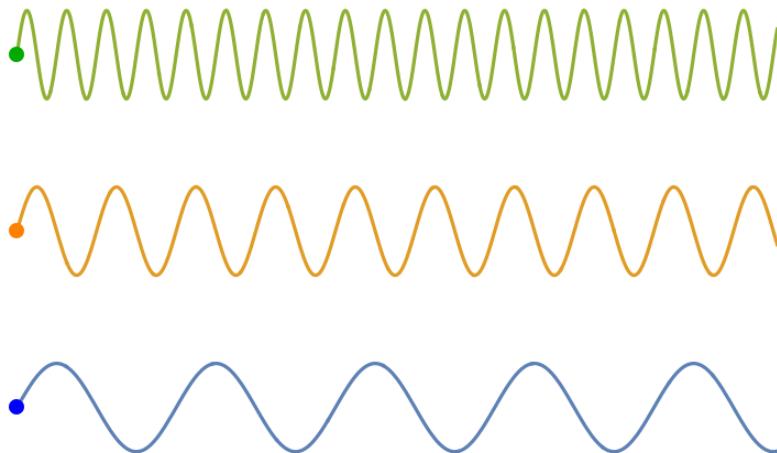
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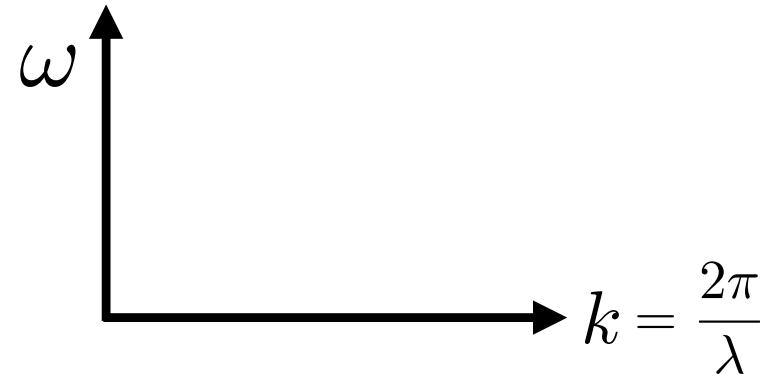
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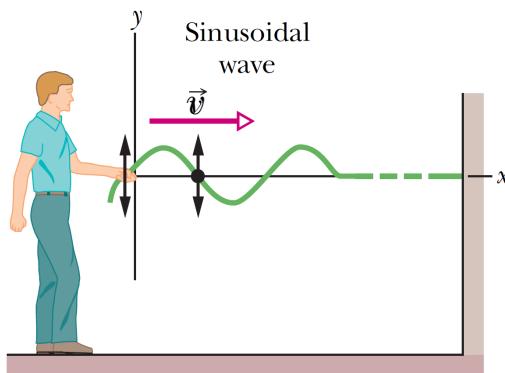


Dispersion relation for waves on a string



Elastic waves

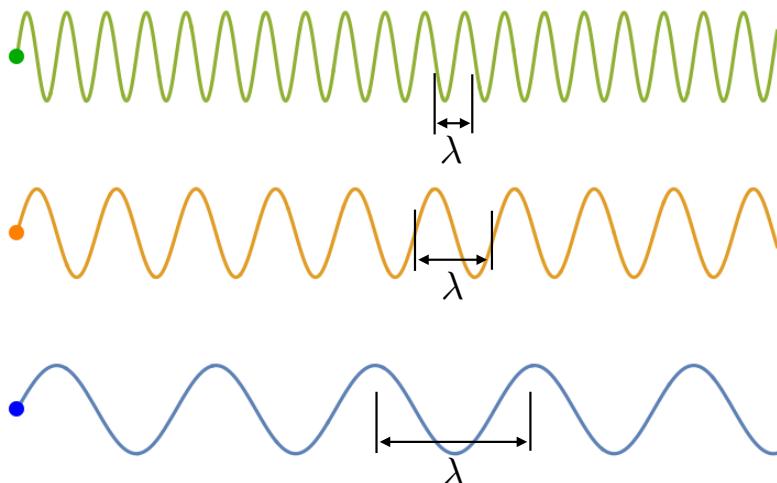
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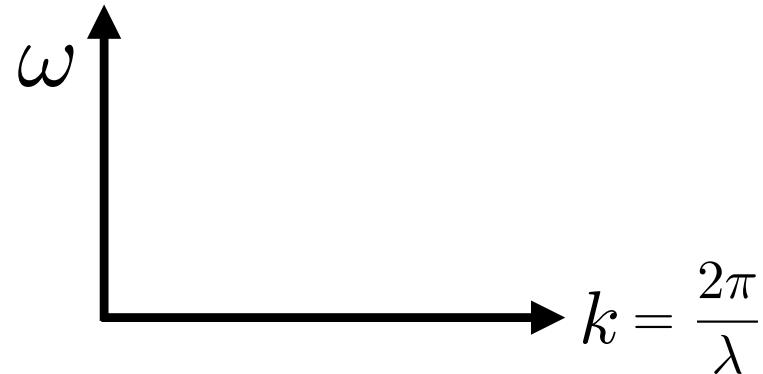
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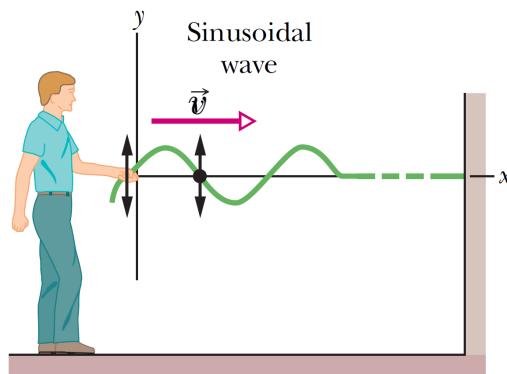


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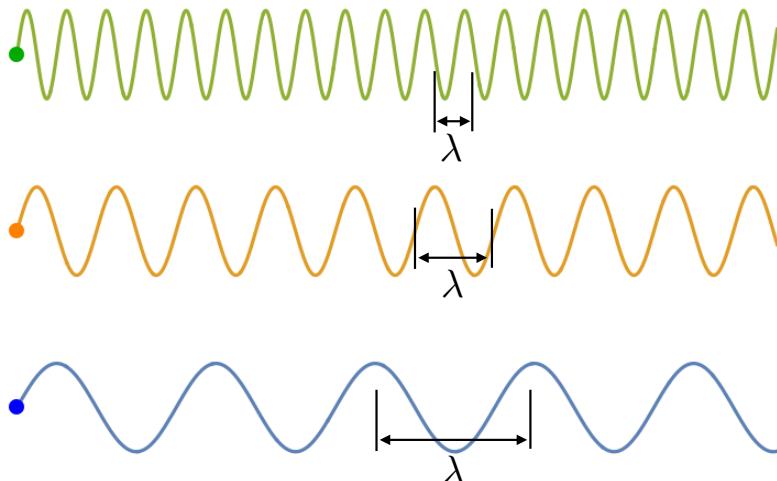
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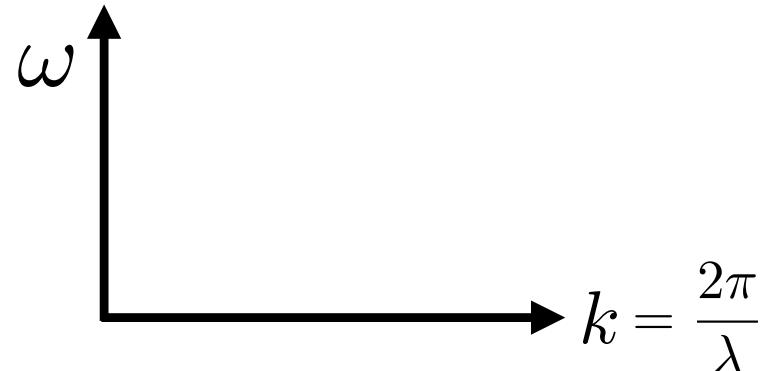
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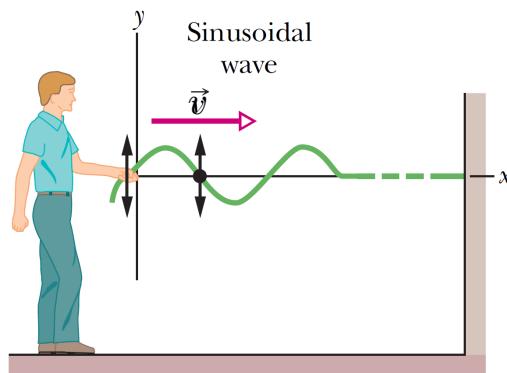


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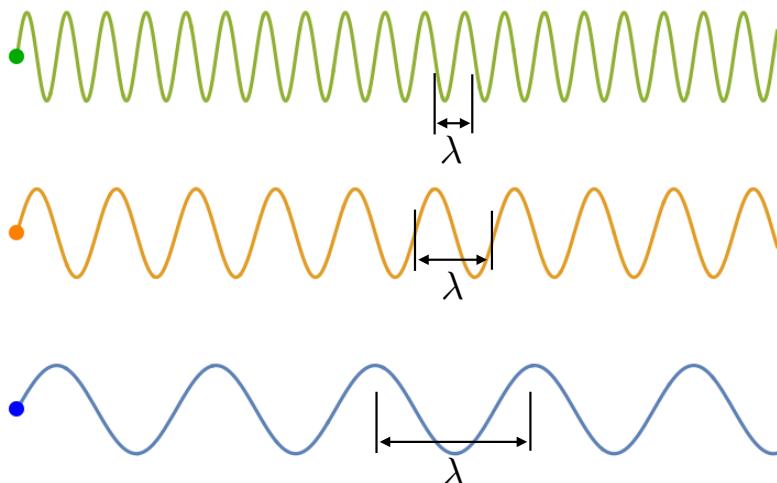
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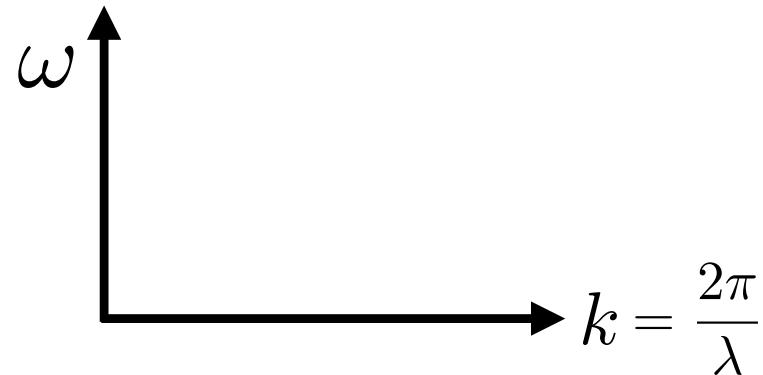
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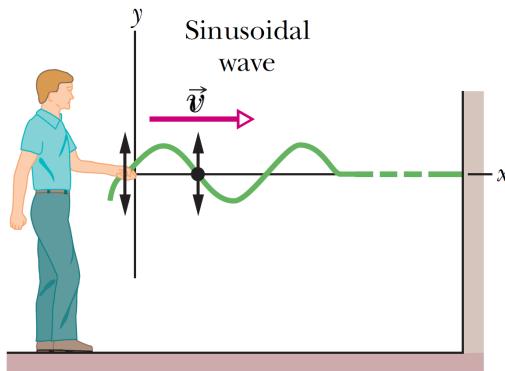


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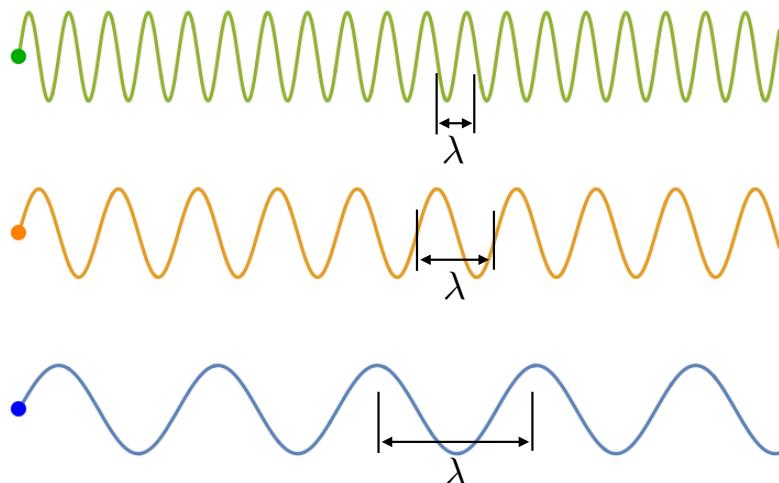
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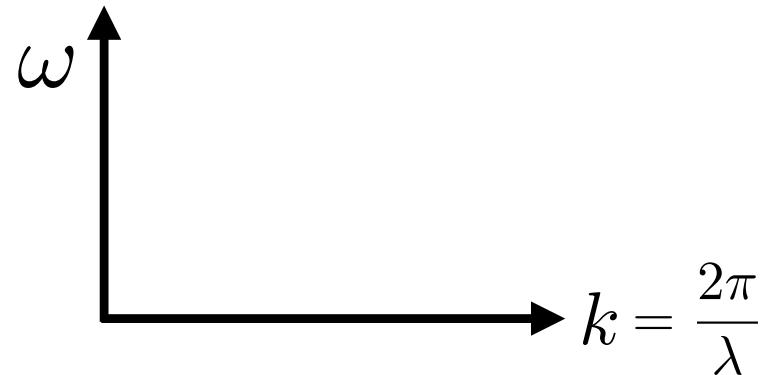
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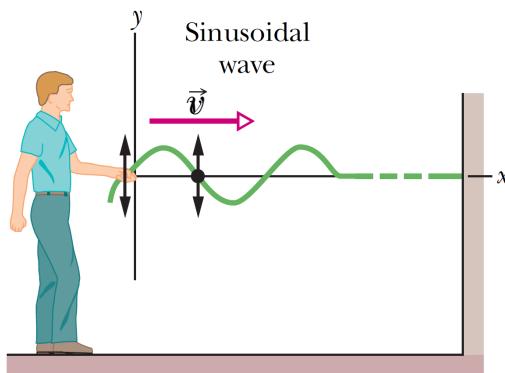


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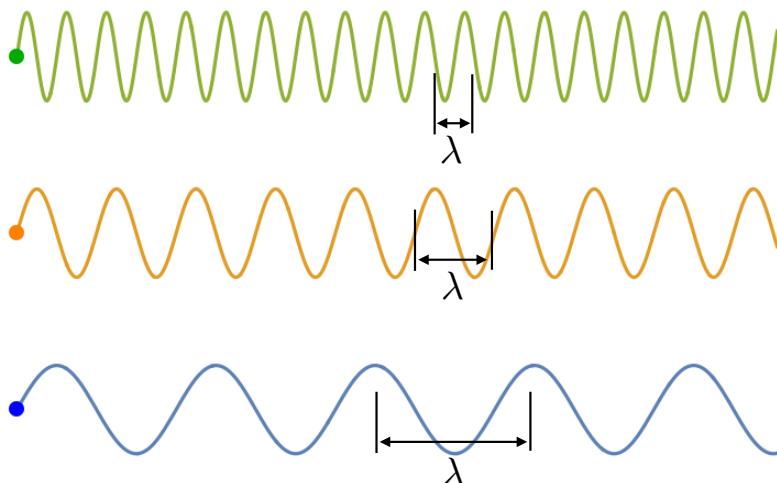
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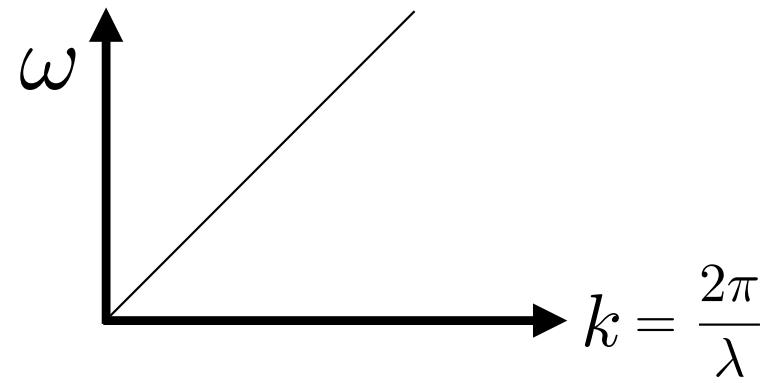
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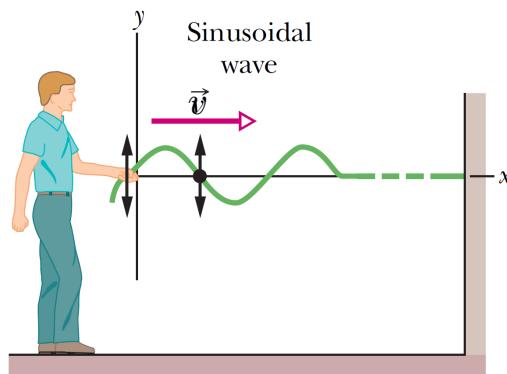


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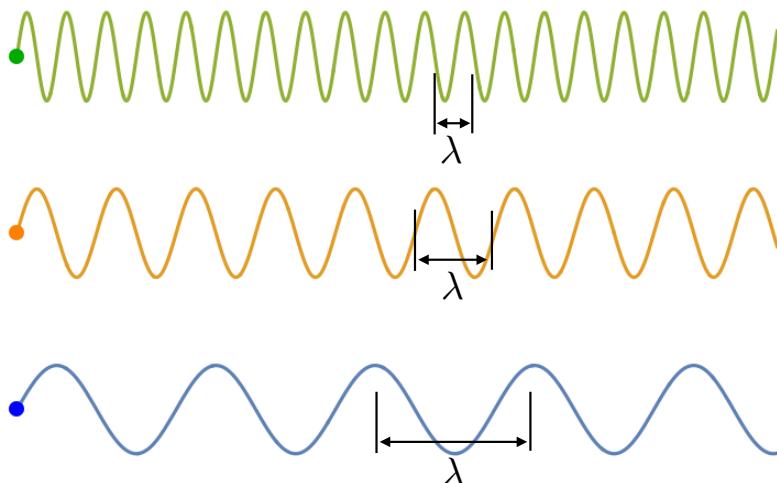
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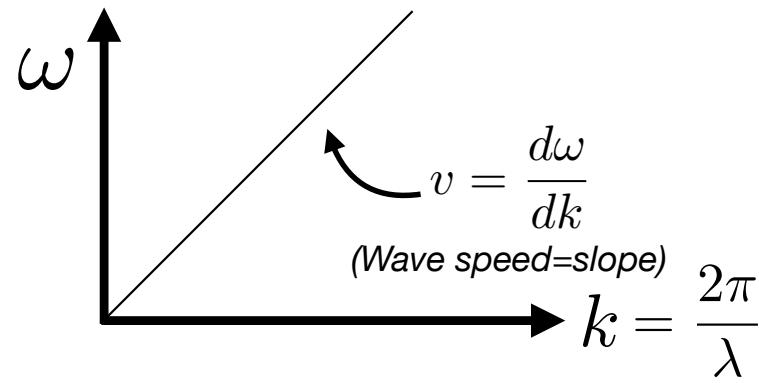
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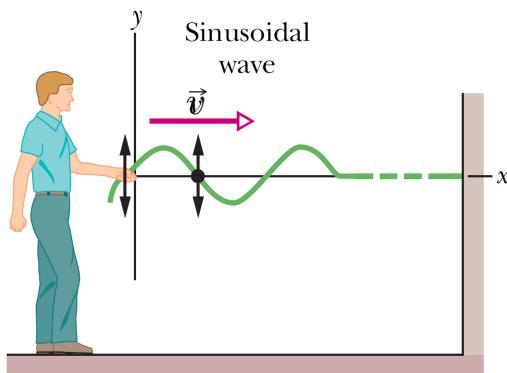


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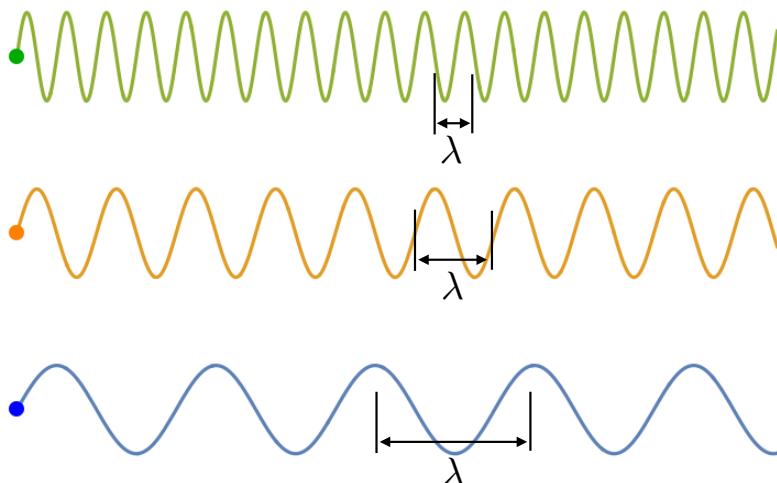
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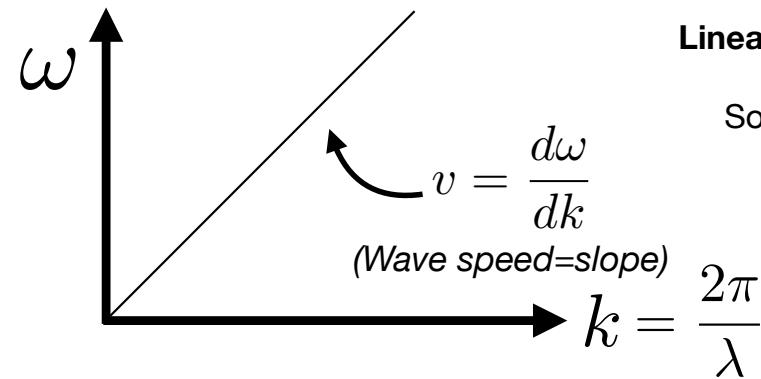
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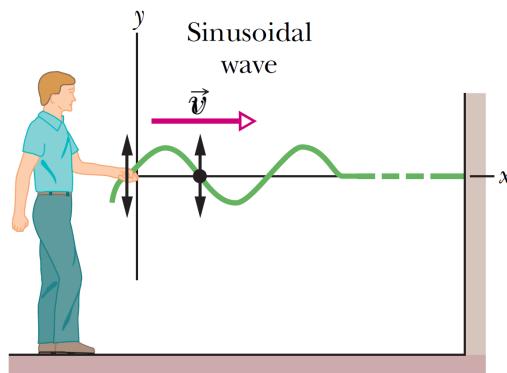
Dispersion relation for waves on a string



Linear dispersion also found in:
X-rays and light
Sound at long wavelengths
Shallow water waves
Gravitational waves

Elastic waves

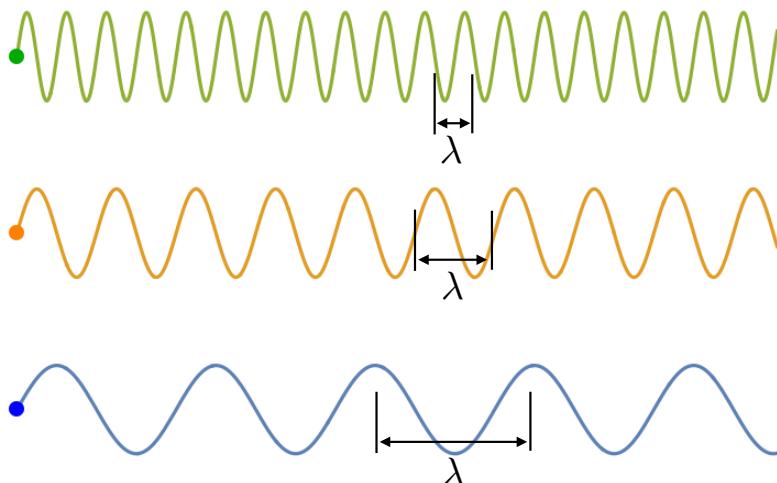
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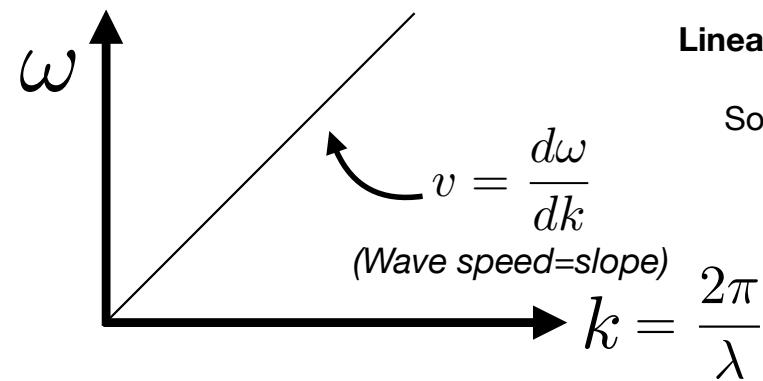
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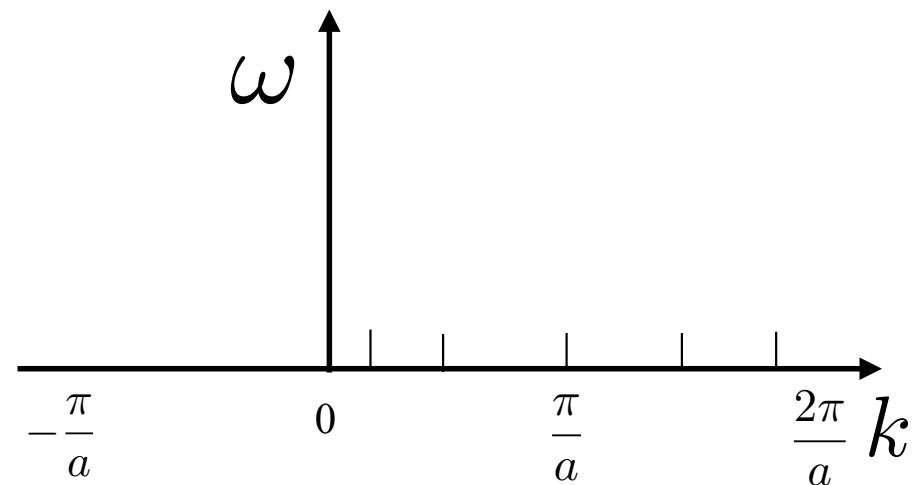
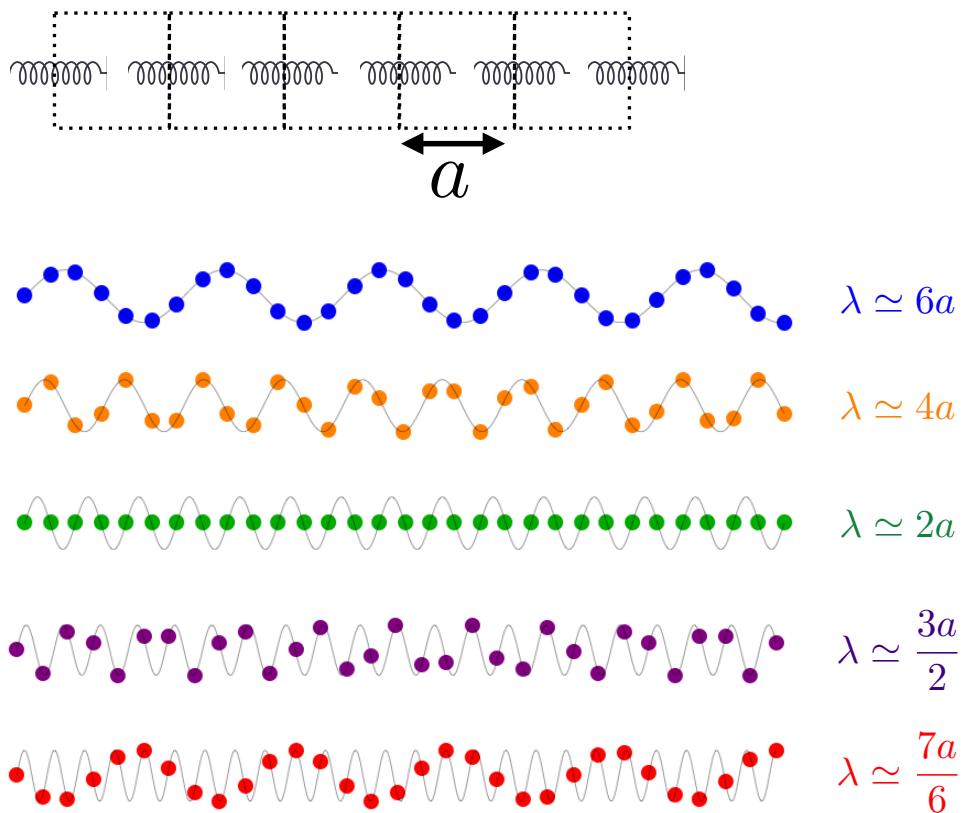
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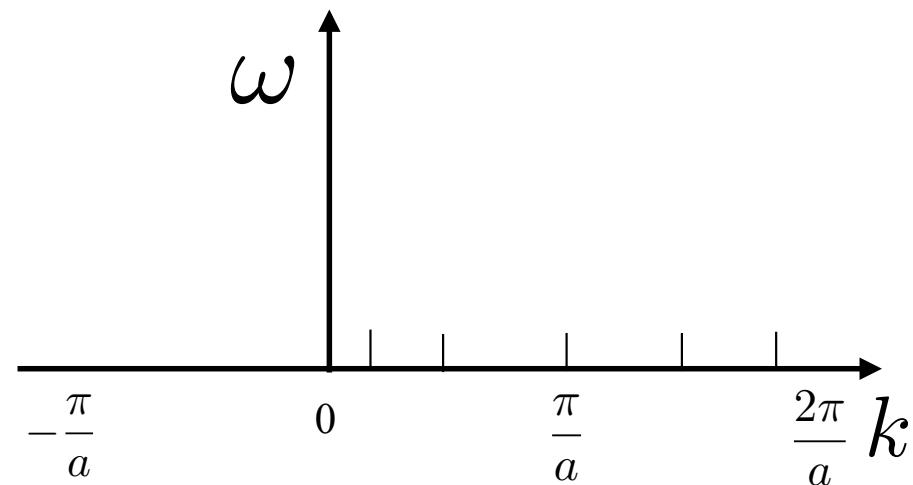
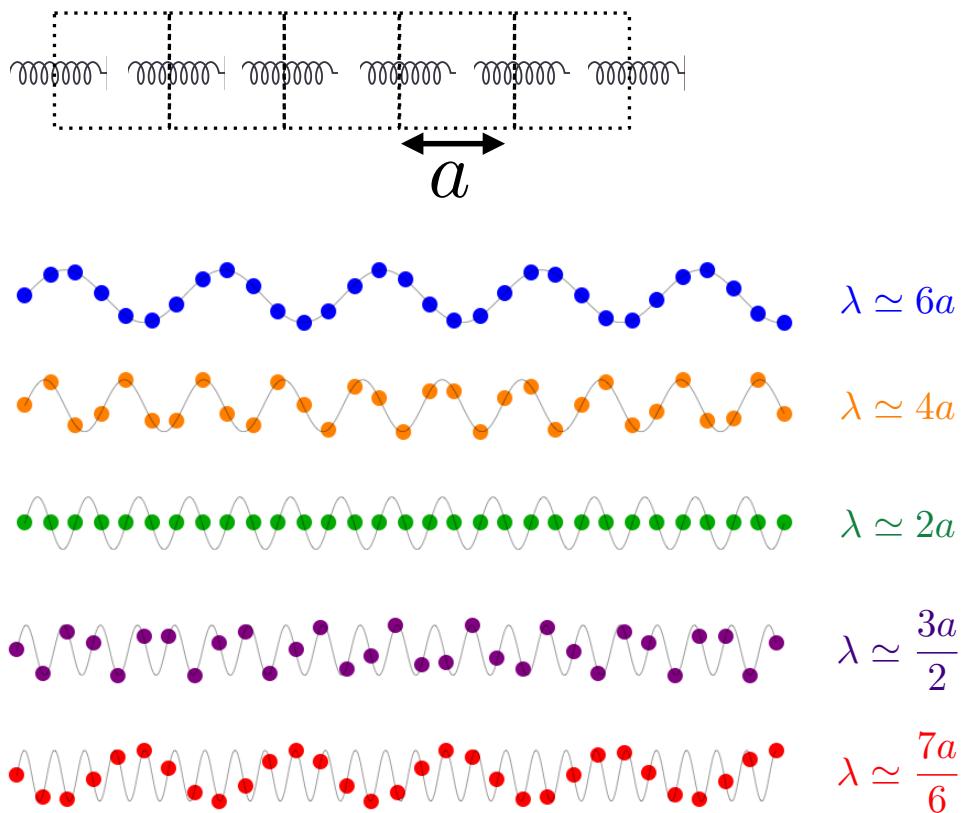
What about atomic scale disturbances?

- Consider monatomic 1D chain and wavelengths approaching lattice parameter a



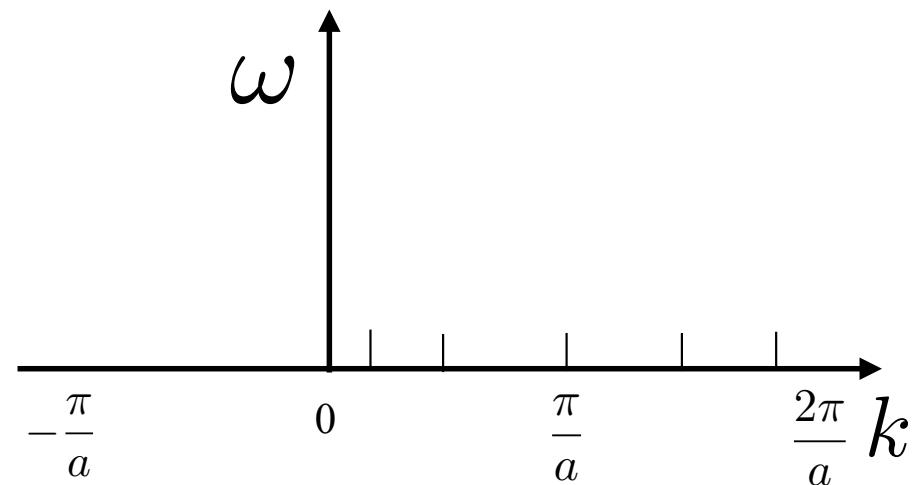
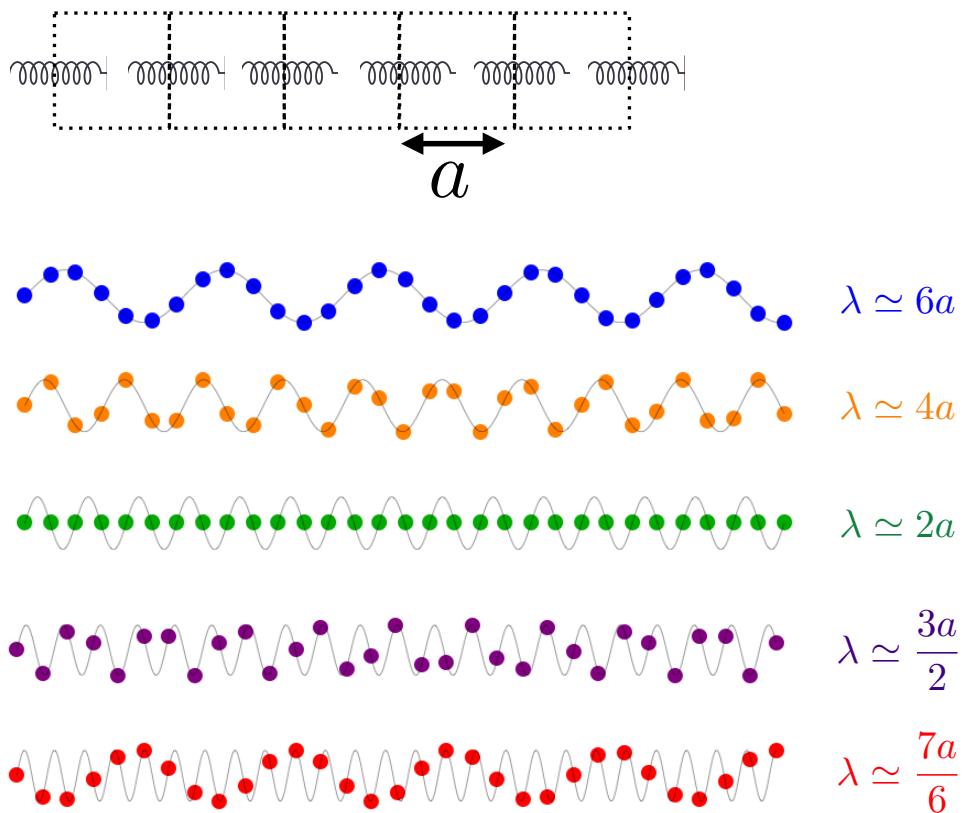
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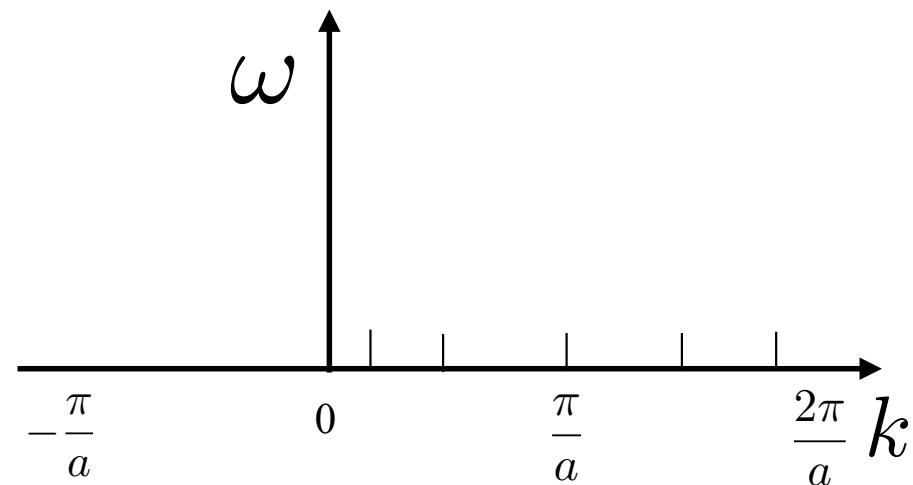
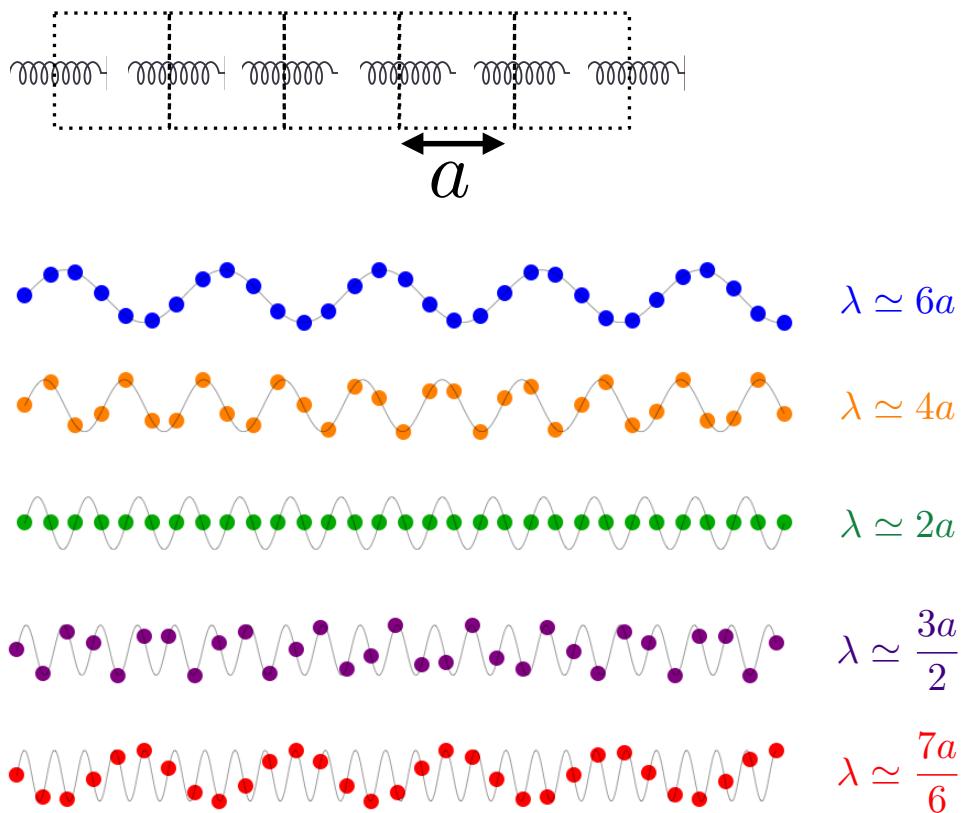
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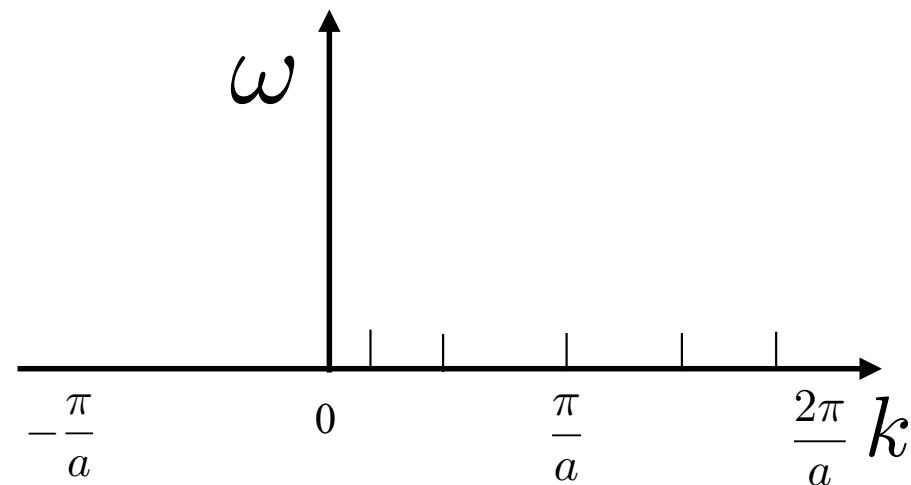
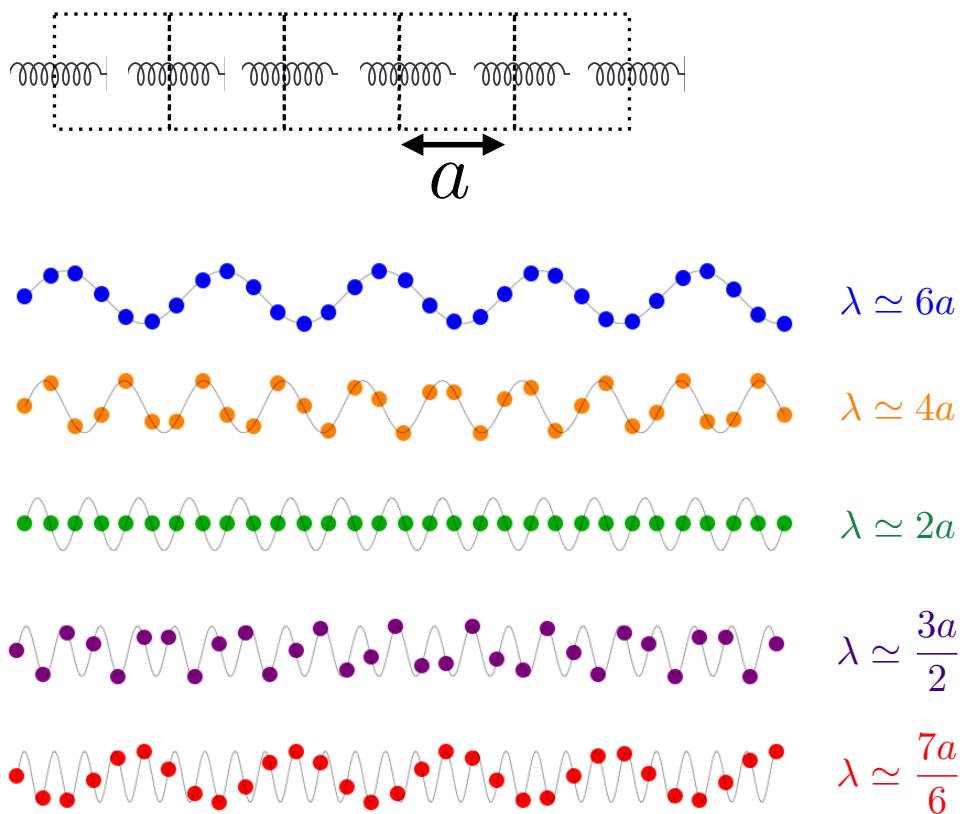
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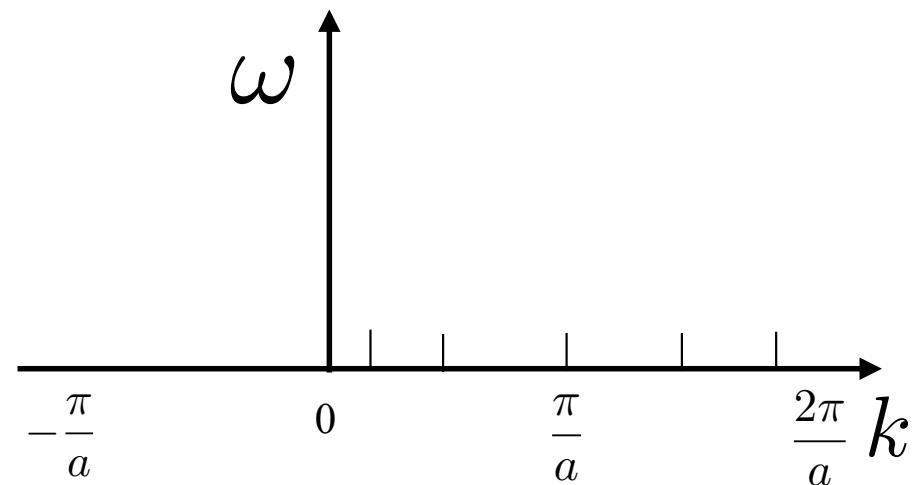
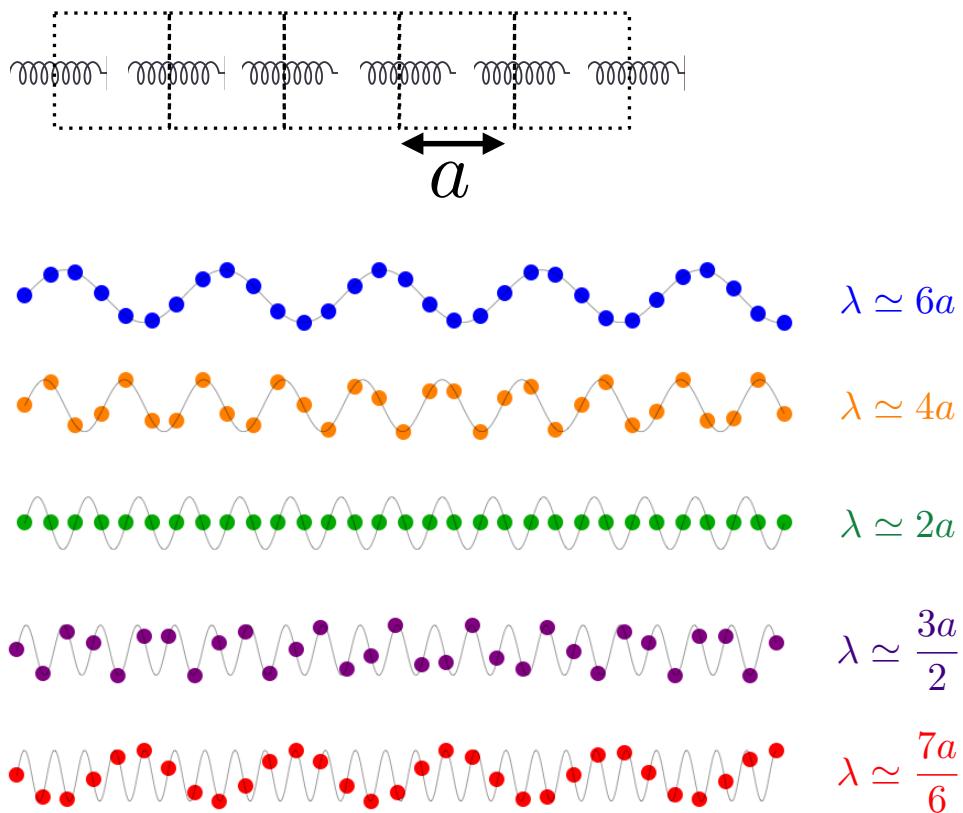
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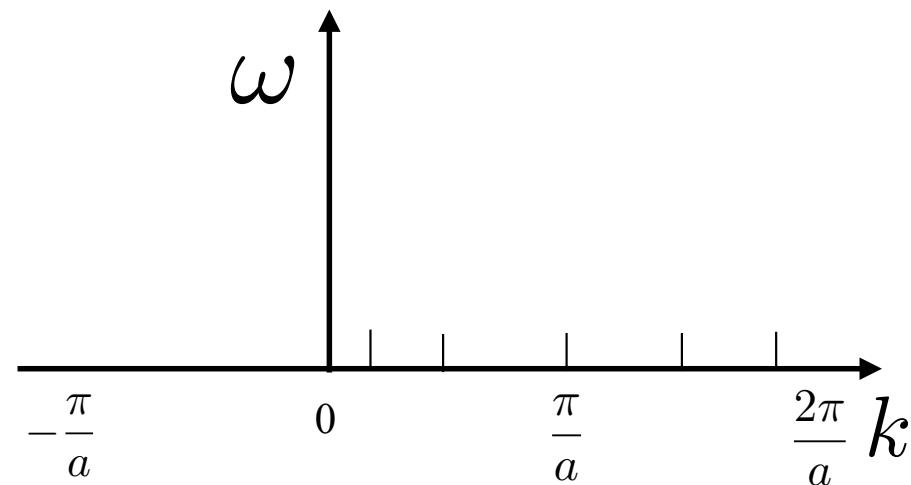
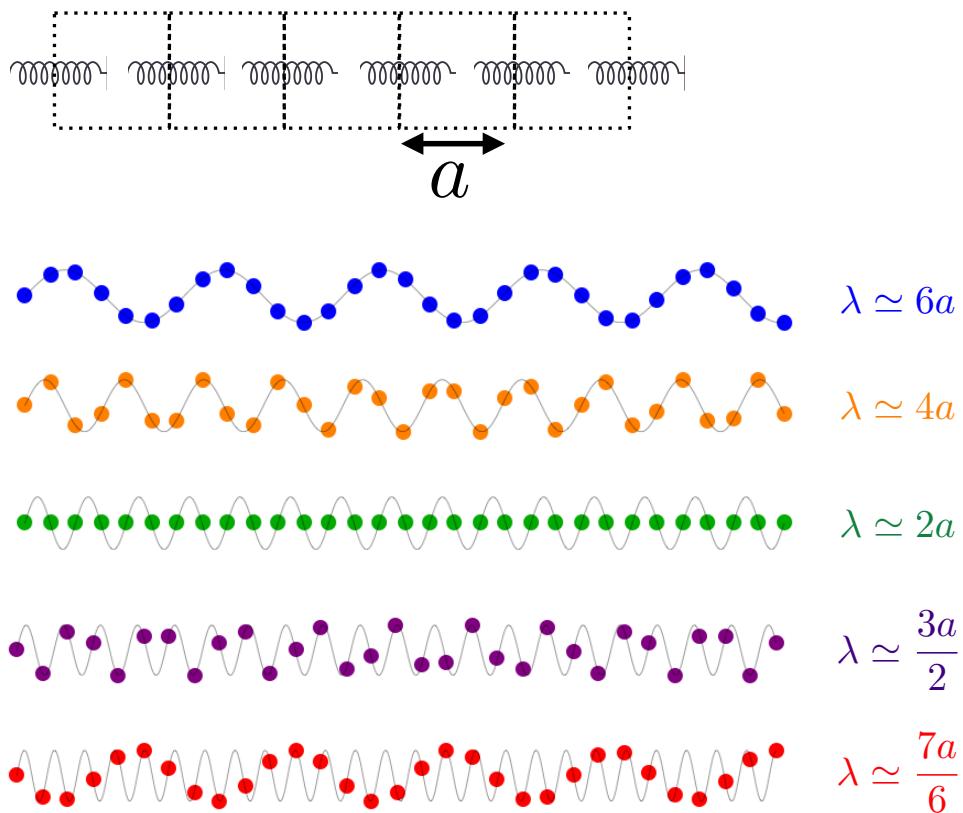
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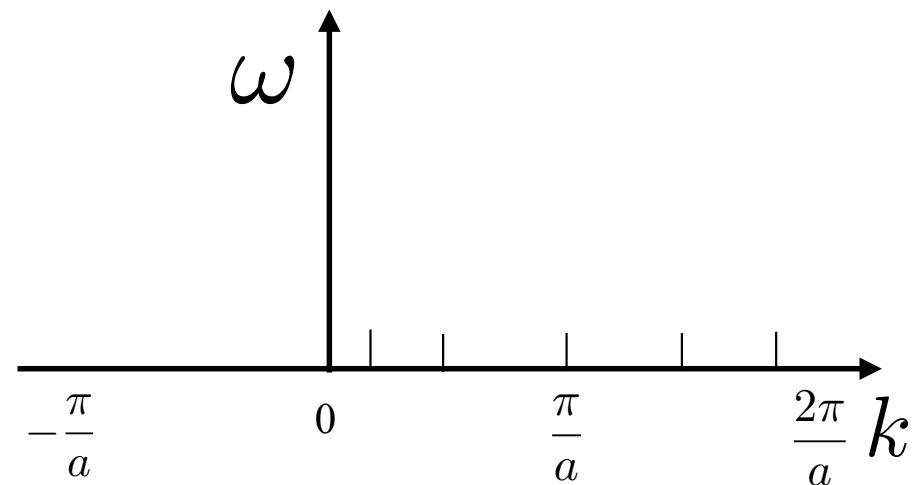
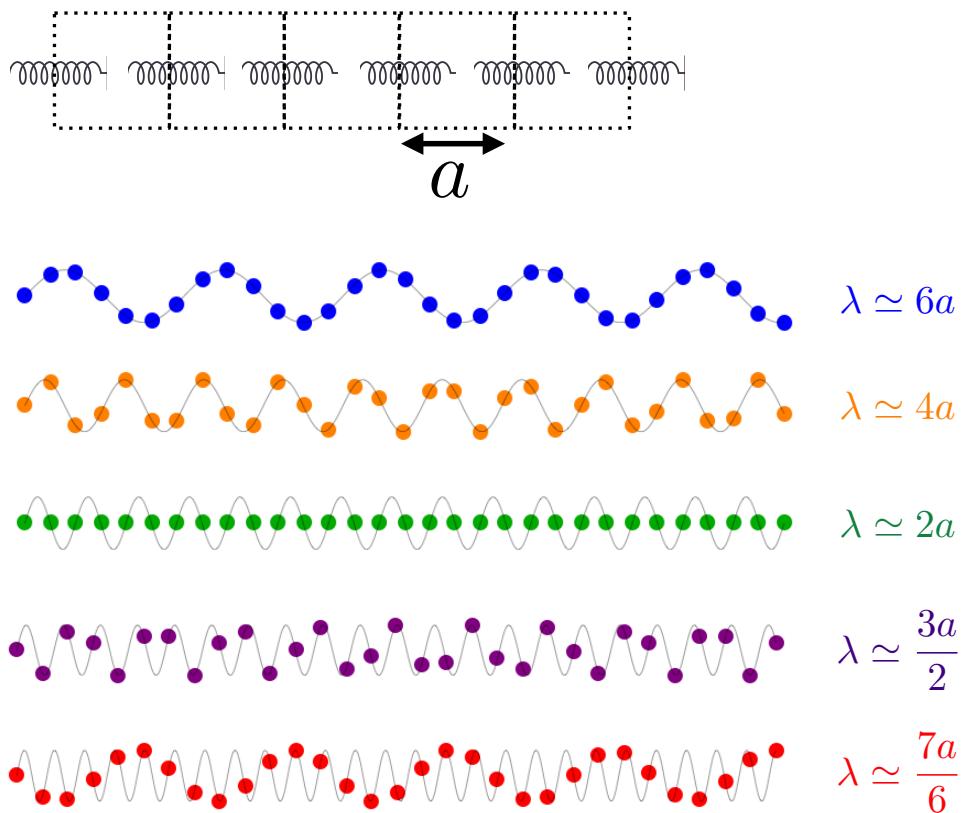
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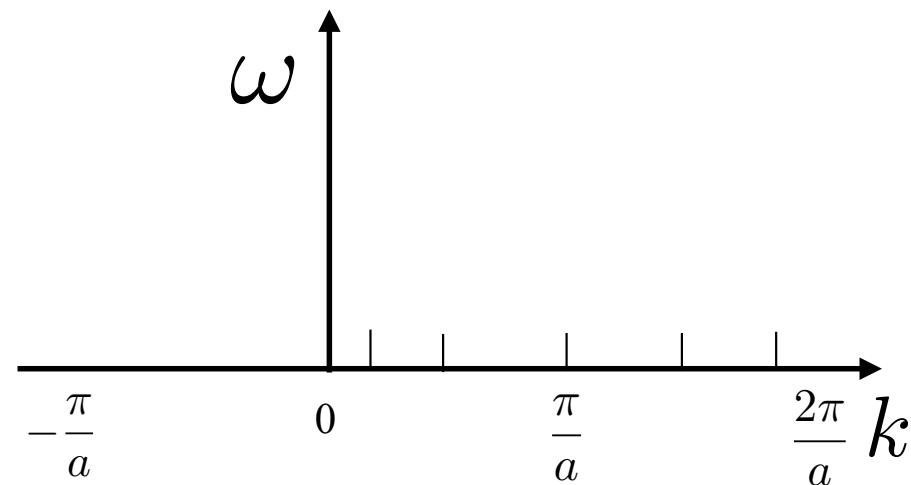
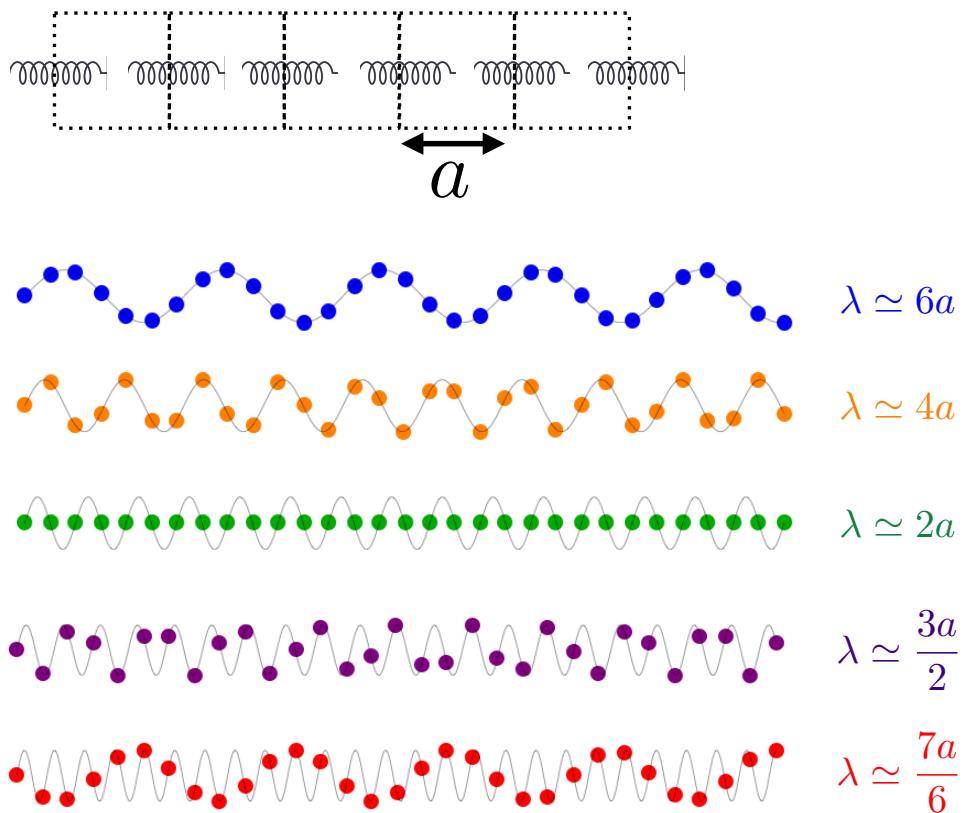
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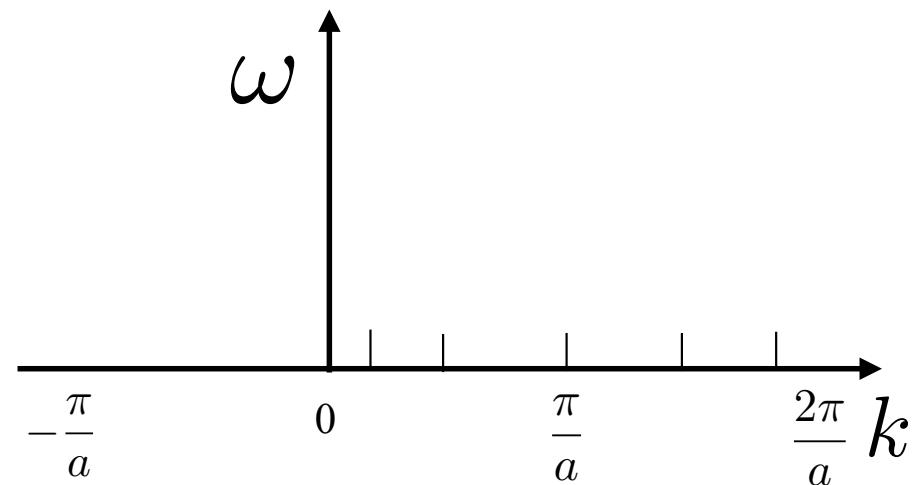
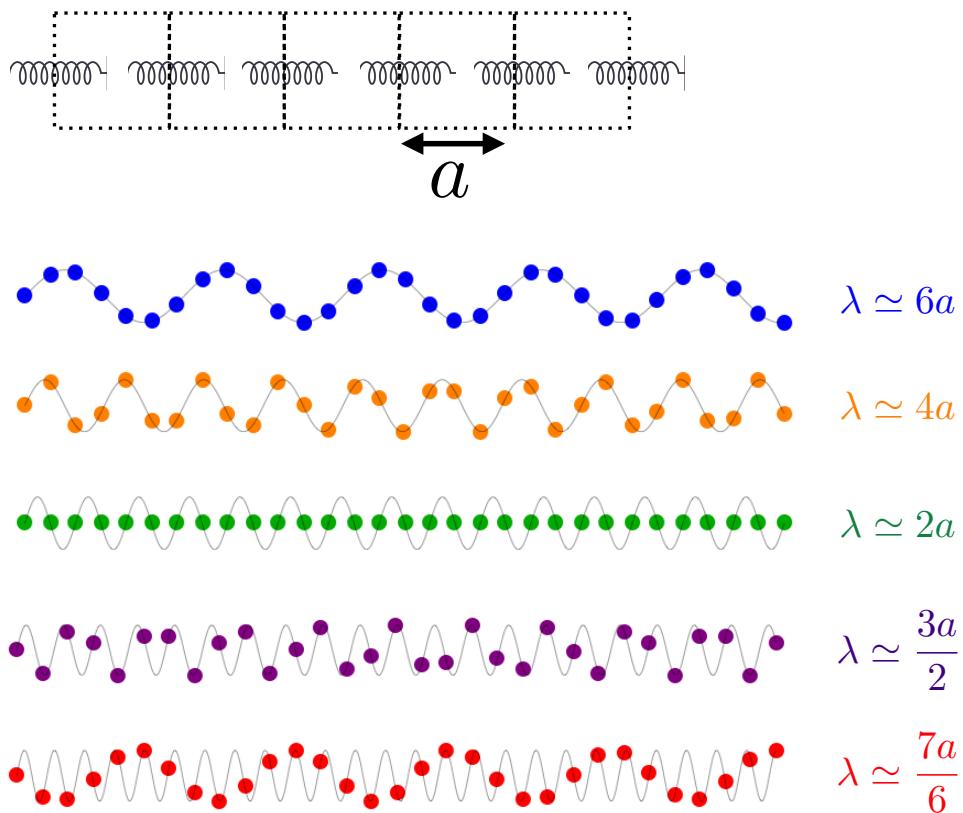
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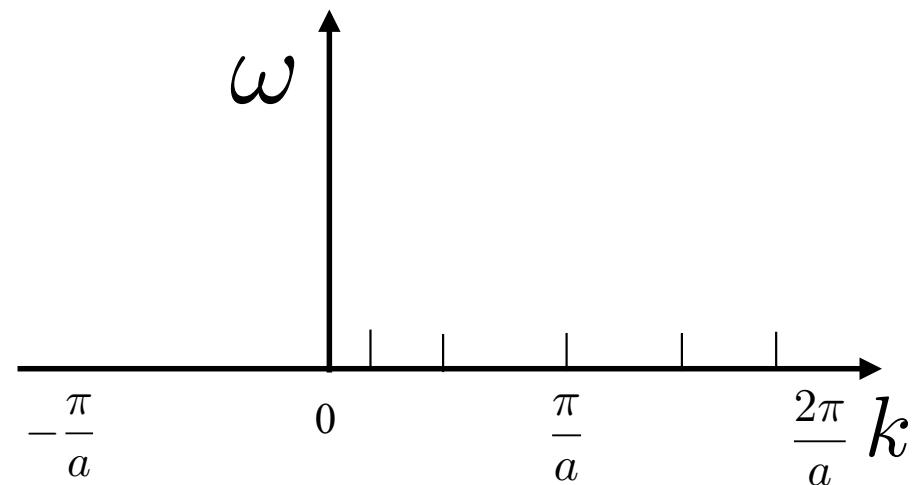
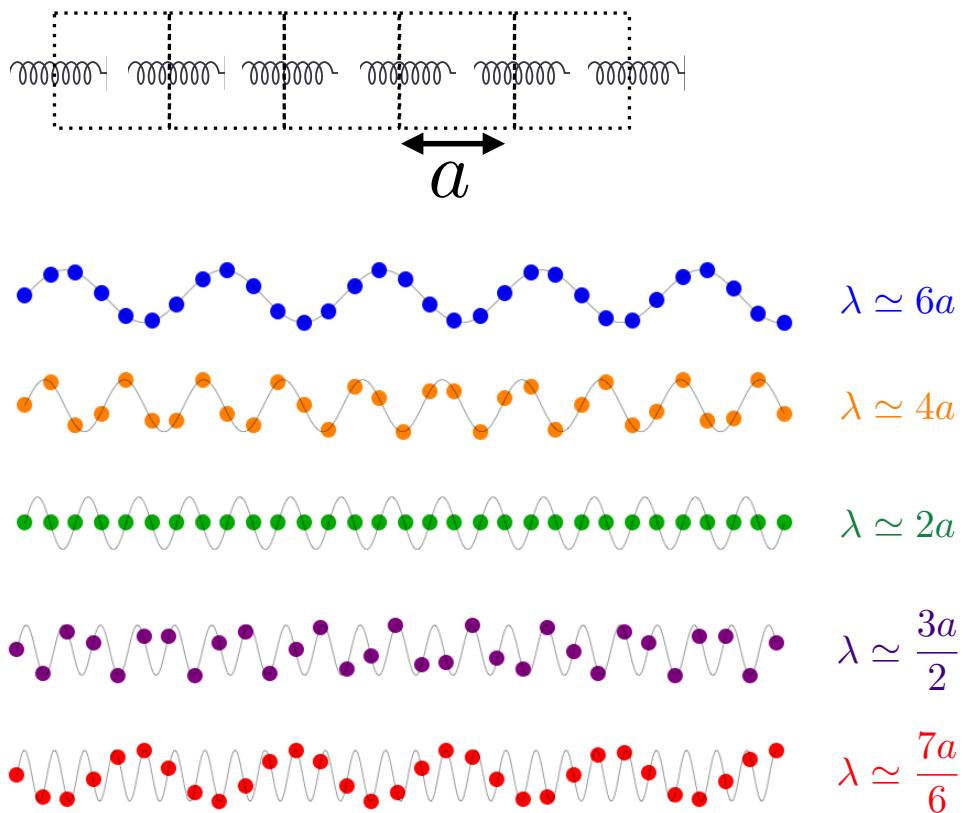
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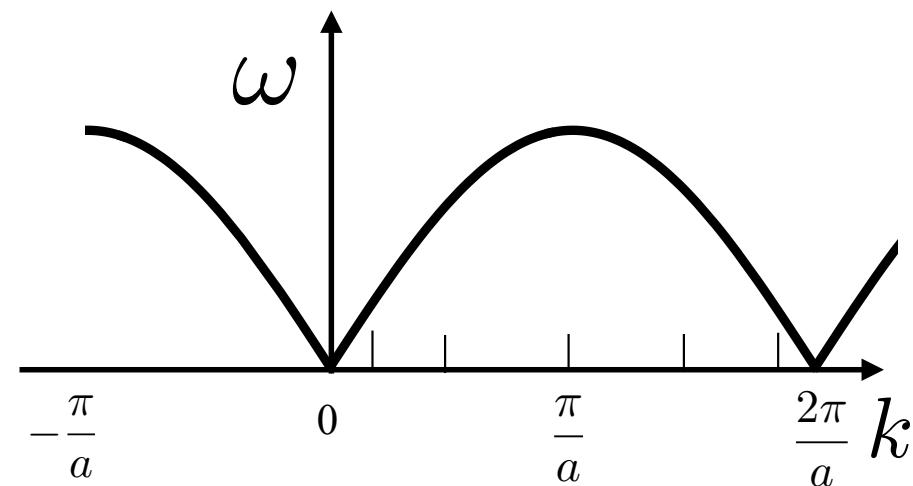
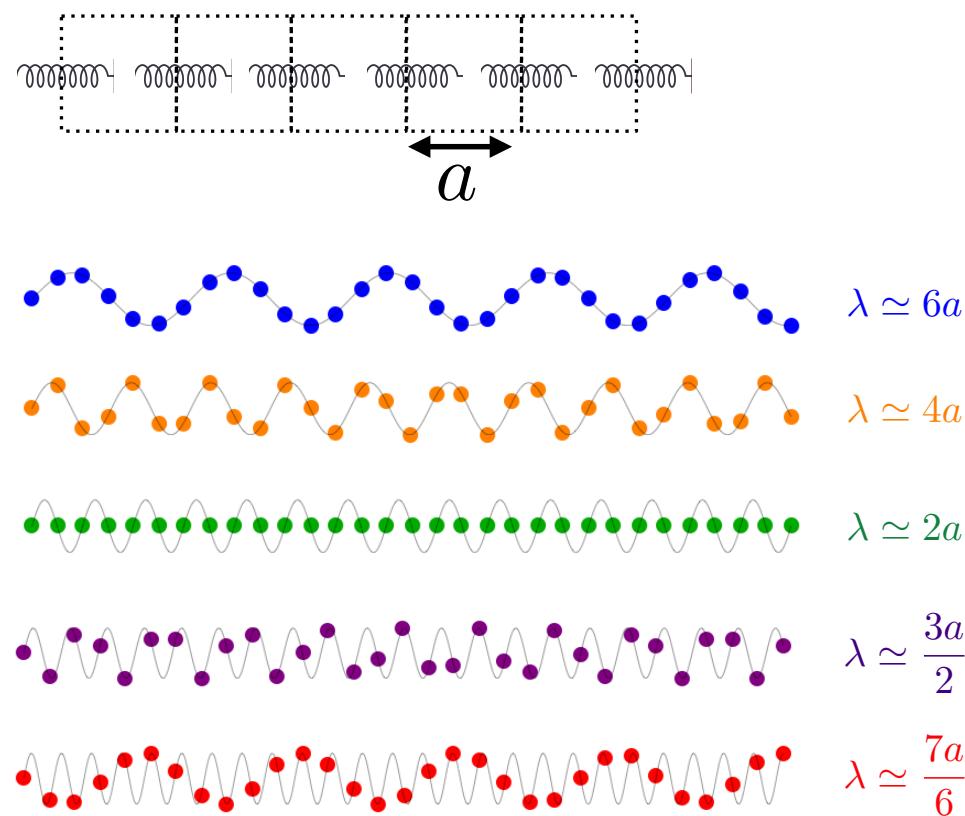
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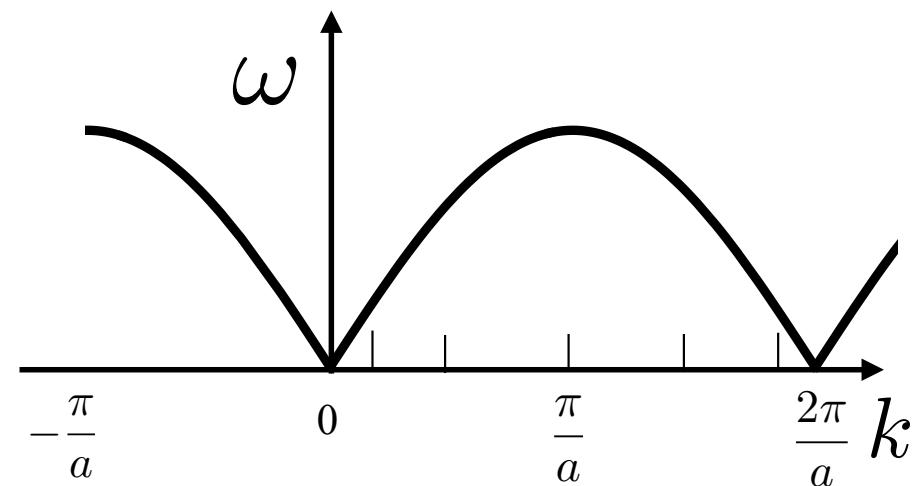
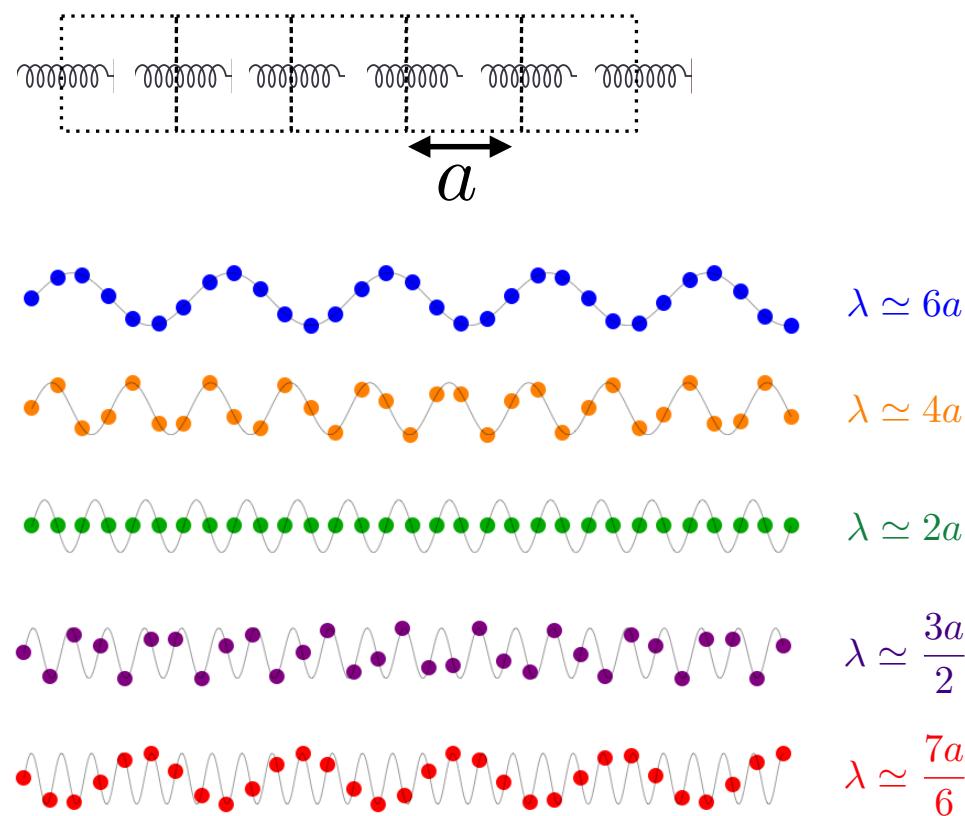
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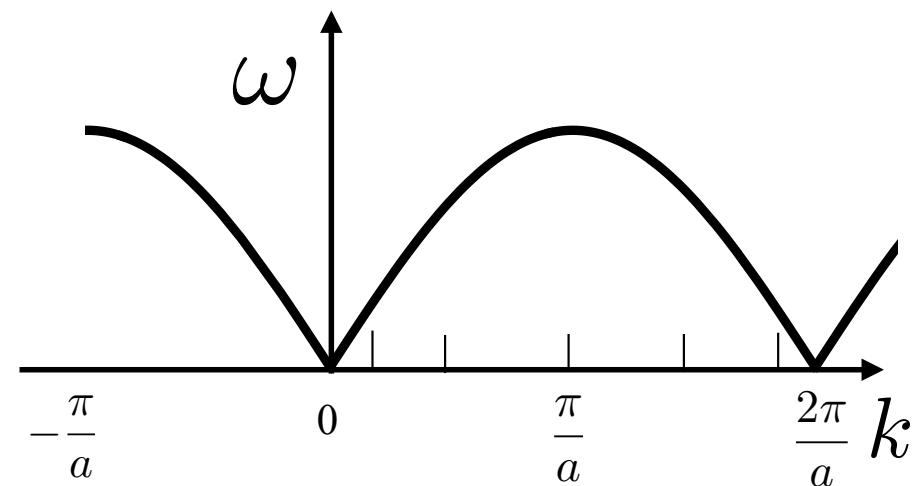
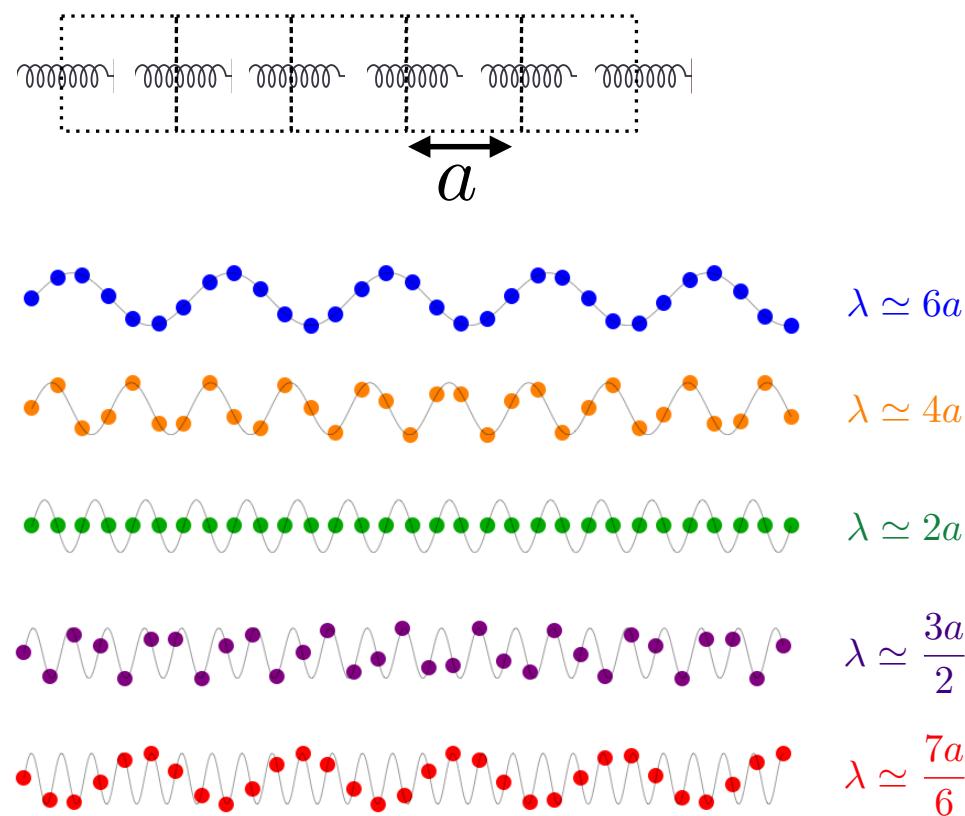
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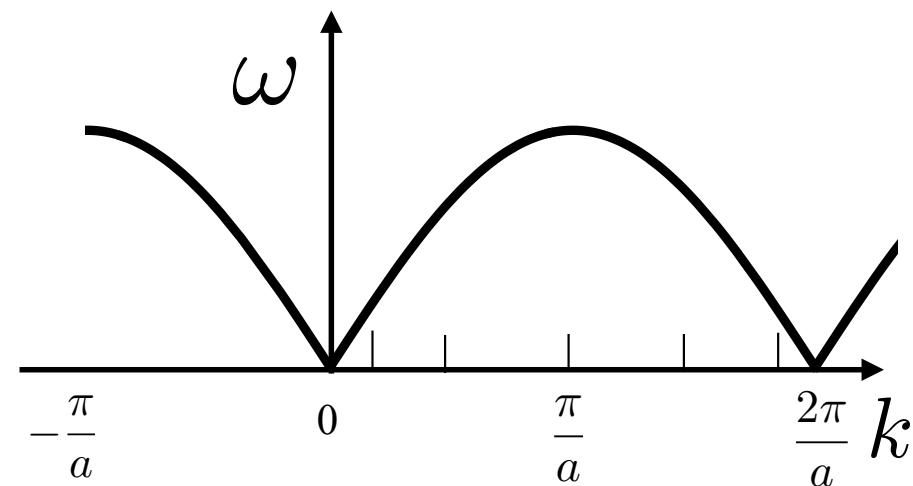
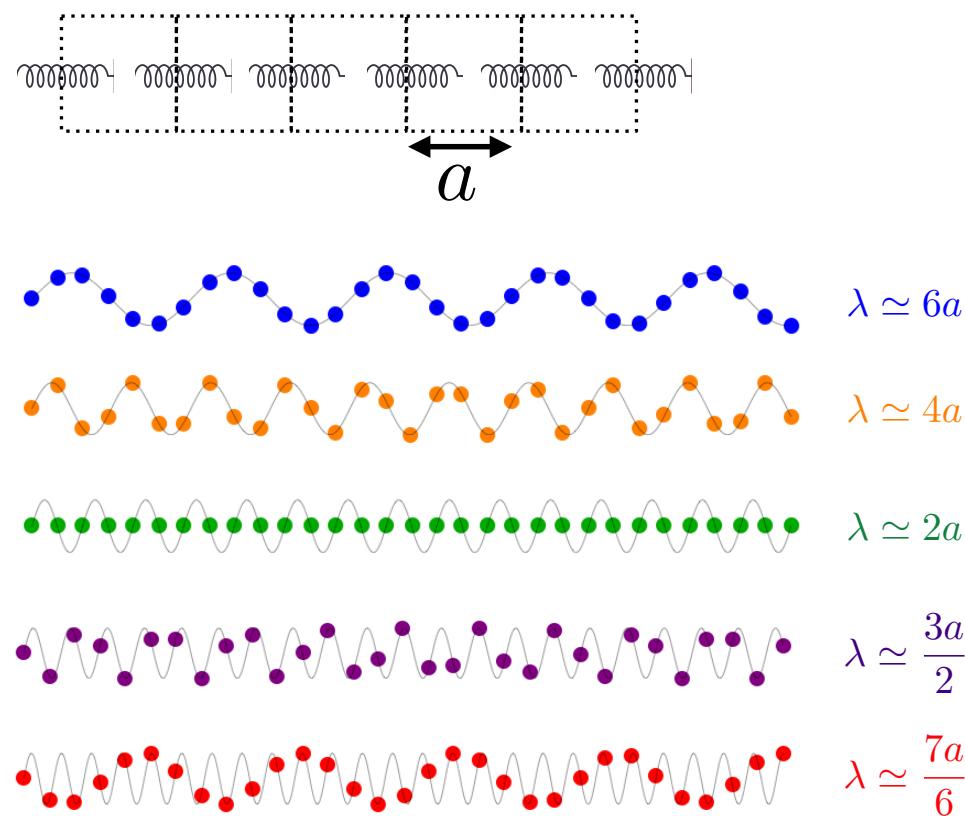
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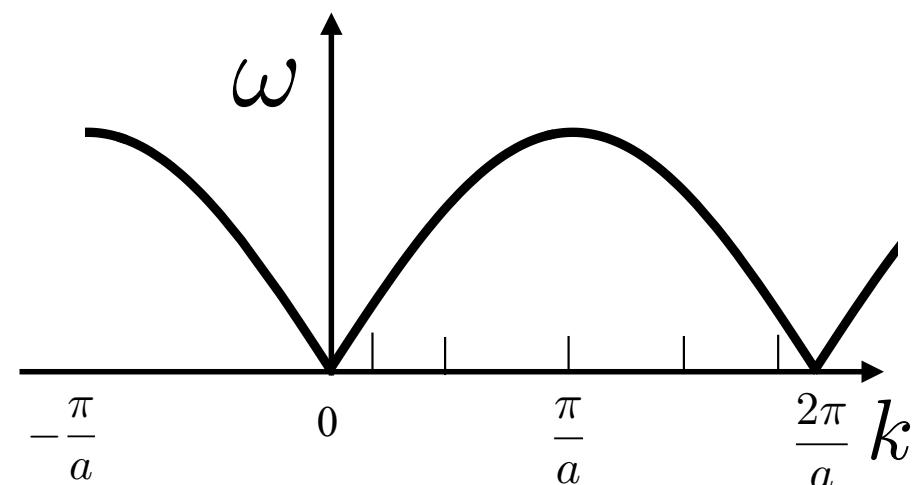
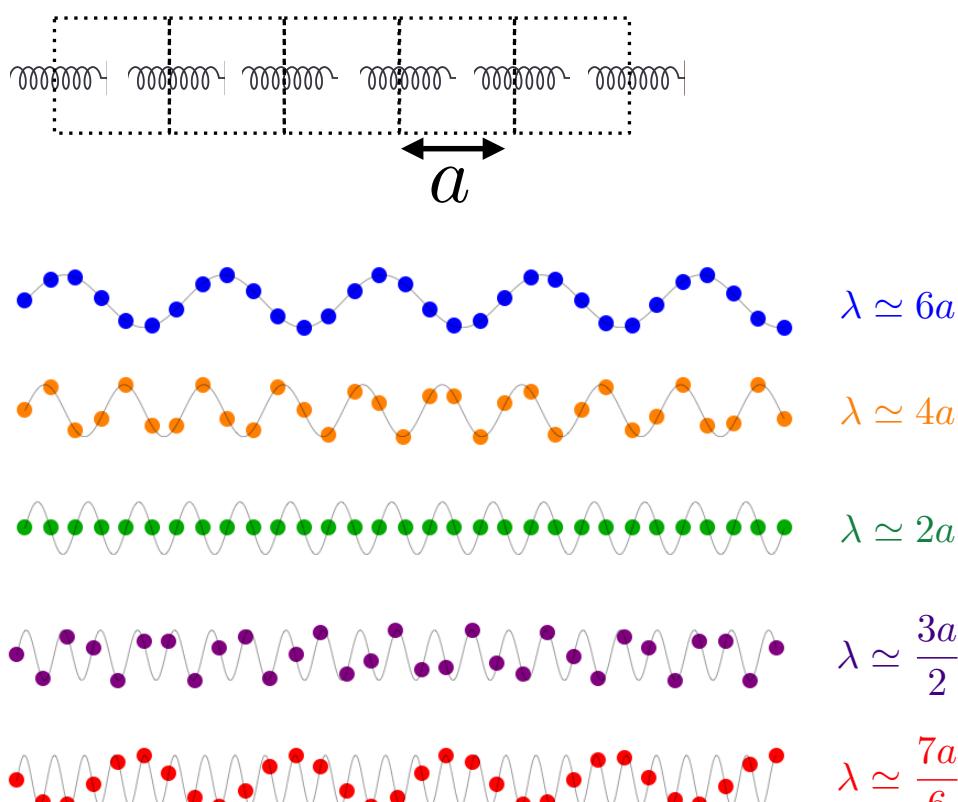
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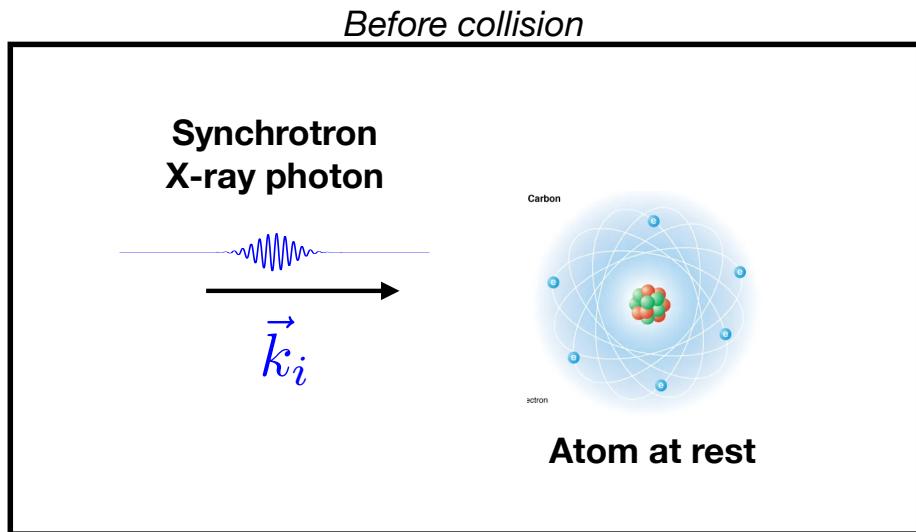
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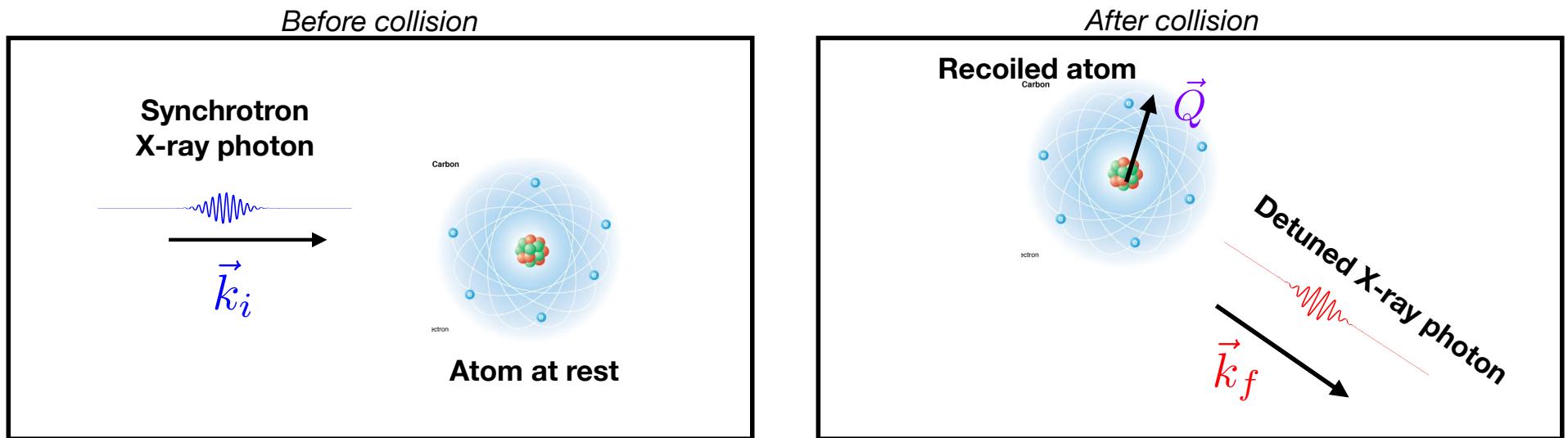
Lattice waves are periodic in momentum, just like the lattice
Quantum mechanics applied to lattice vibrations lead to the concept of “phonons”
Loosely, each k value refers to a different “phonon mode”

How do X-rays scatter from atoms?



- Off resonance, an X-ray can cause an atom to recoil in a momentum conserving collision

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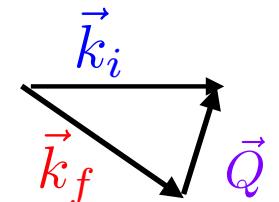


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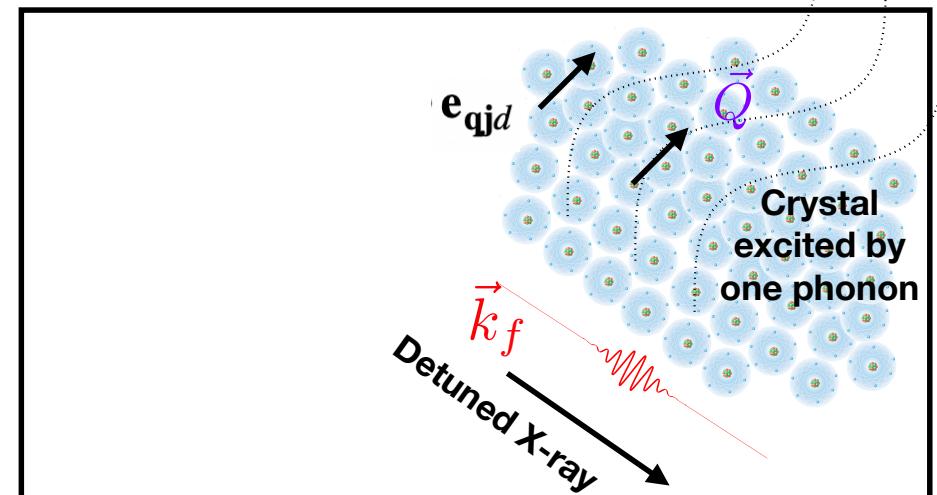
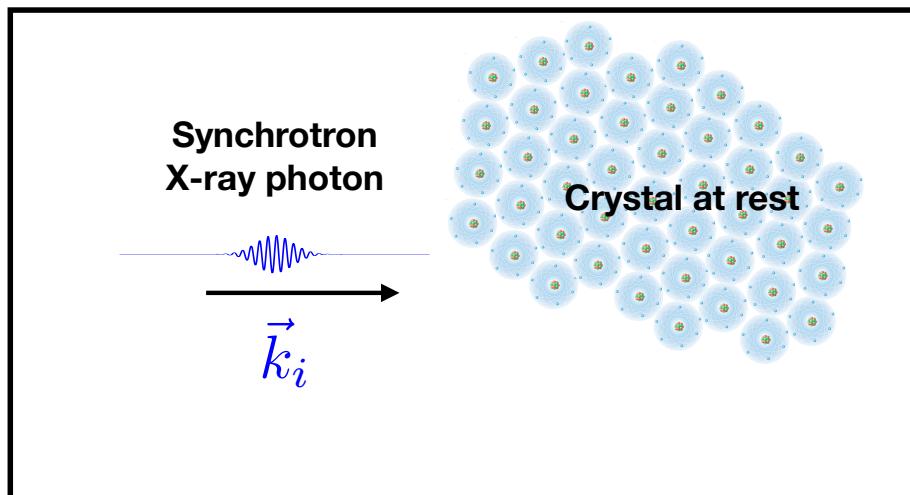


$$\vec{Q} = \vec{k}_f - \vec{k}_i \quad \text{is the momentum delivered to the atom}$$



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How do X-rays scatter from crystals?

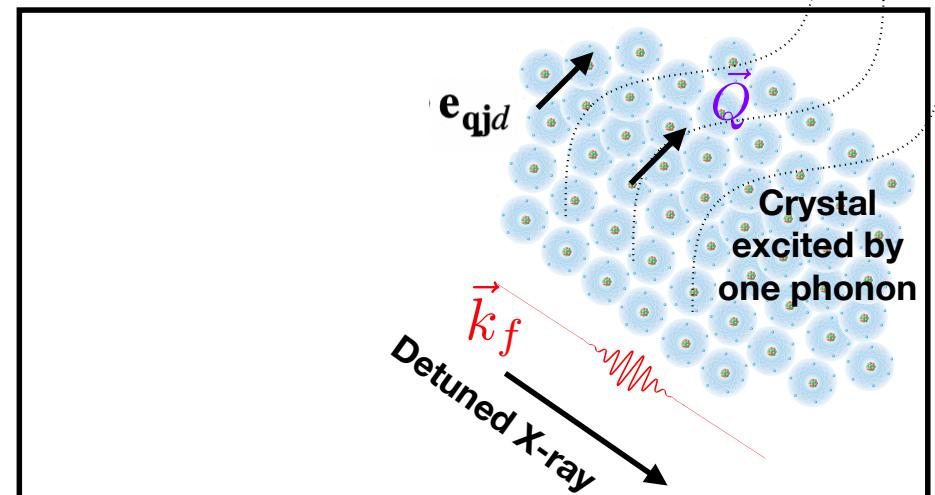
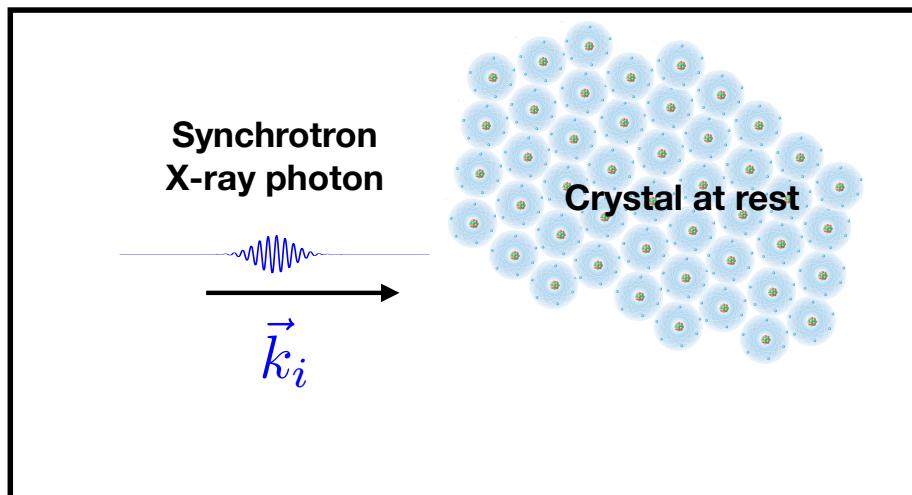


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$$|F_{1p}(\tau, \mathbf{qj})|^2 = \frac{1}{\omega_{\mathbf{qj}}} \left| \sum_d \frac{f_d(\mathbf{Q})}{\sqrt{2M_d}} e^{-W_d} \mathbf{Q} \cdot \mathbf{e}_{\mathbf{qjd}} e^{i\mathbf{Q} \cdot \mathbf{r}_d} \right|^2$$

- Probability of scattering off of a given phonon is strong when the displacements are along \mathbf{Q}

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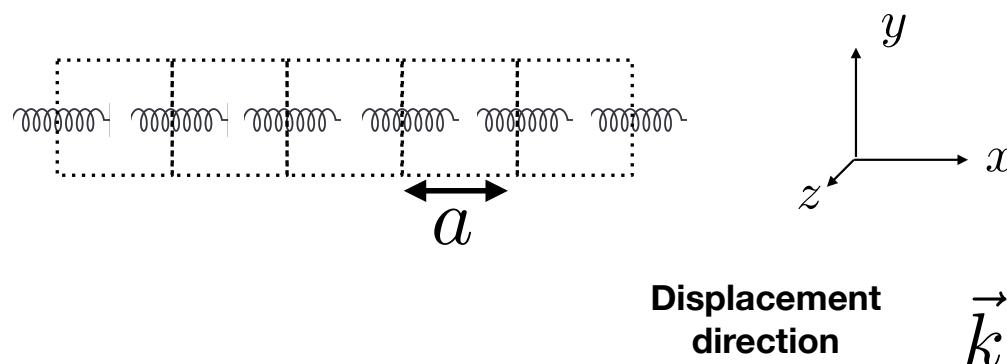
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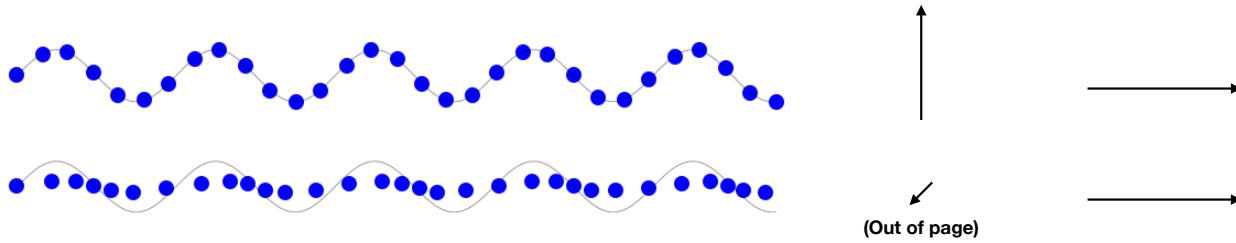
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Polarization of phonons

- Atoms can move in three directions, different dispersion for different directions



Y or Z directions: Transverse acoustic waves

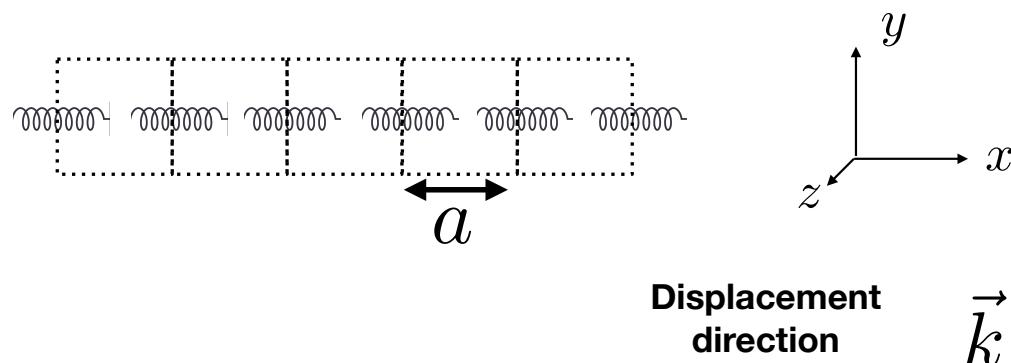


X-direction: Longitudinal acoustic wave (aka sound)

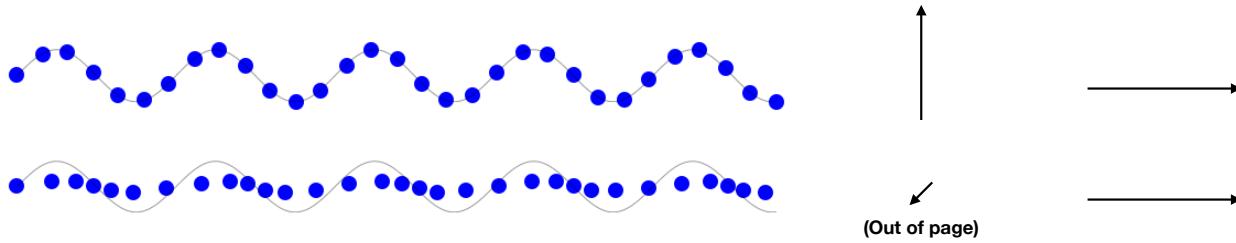


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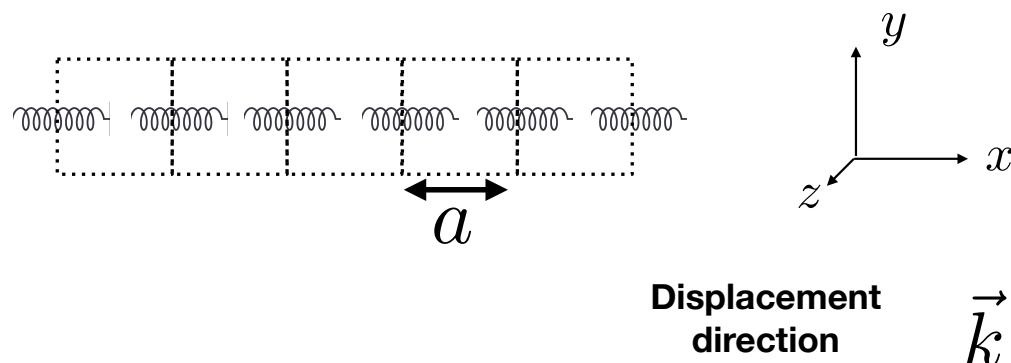


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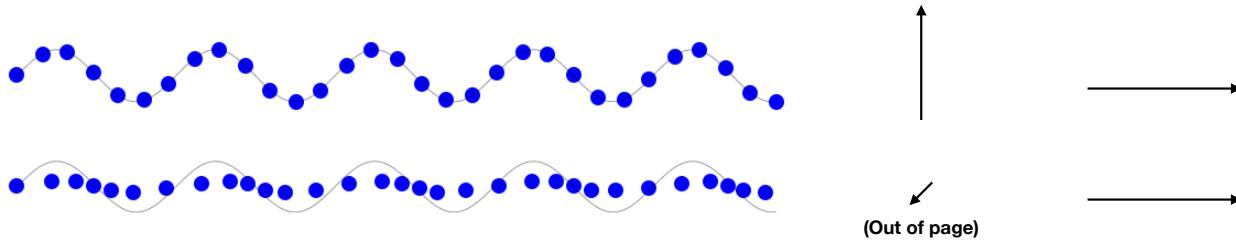


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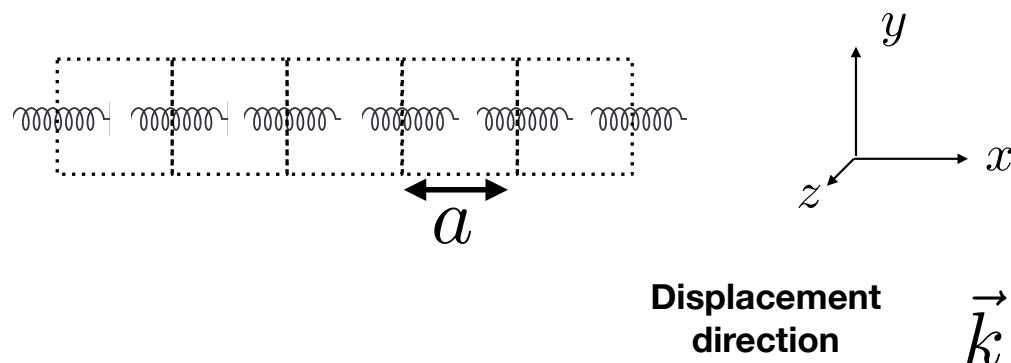


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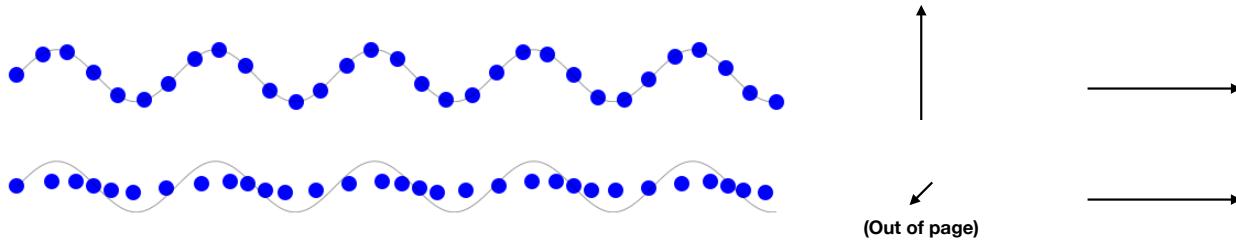


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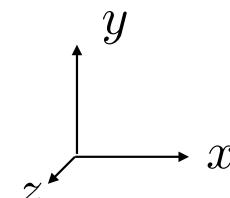
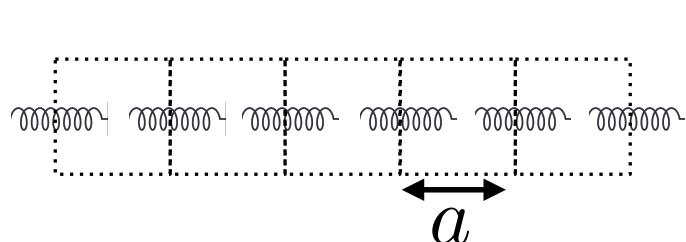


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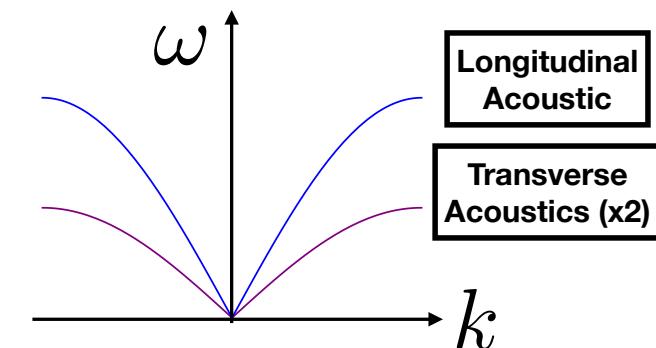
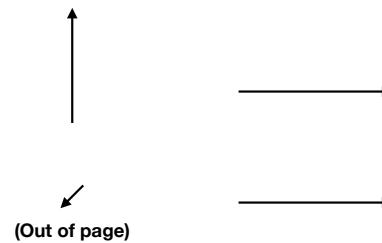
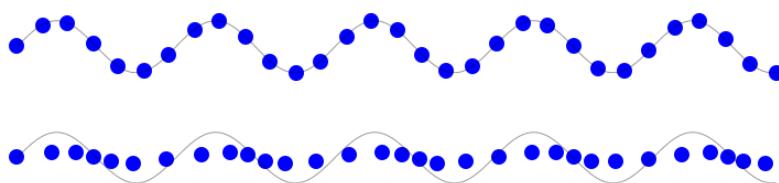
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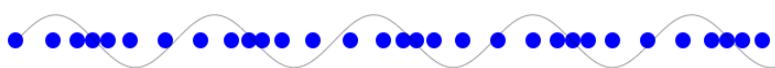


Displacement direction
 \vec{k}

Y or Z directions: Transverse acoustic waves

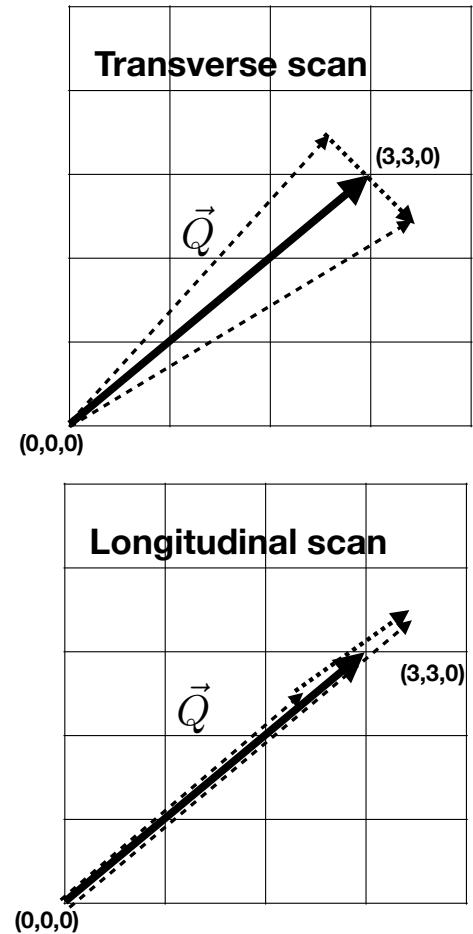
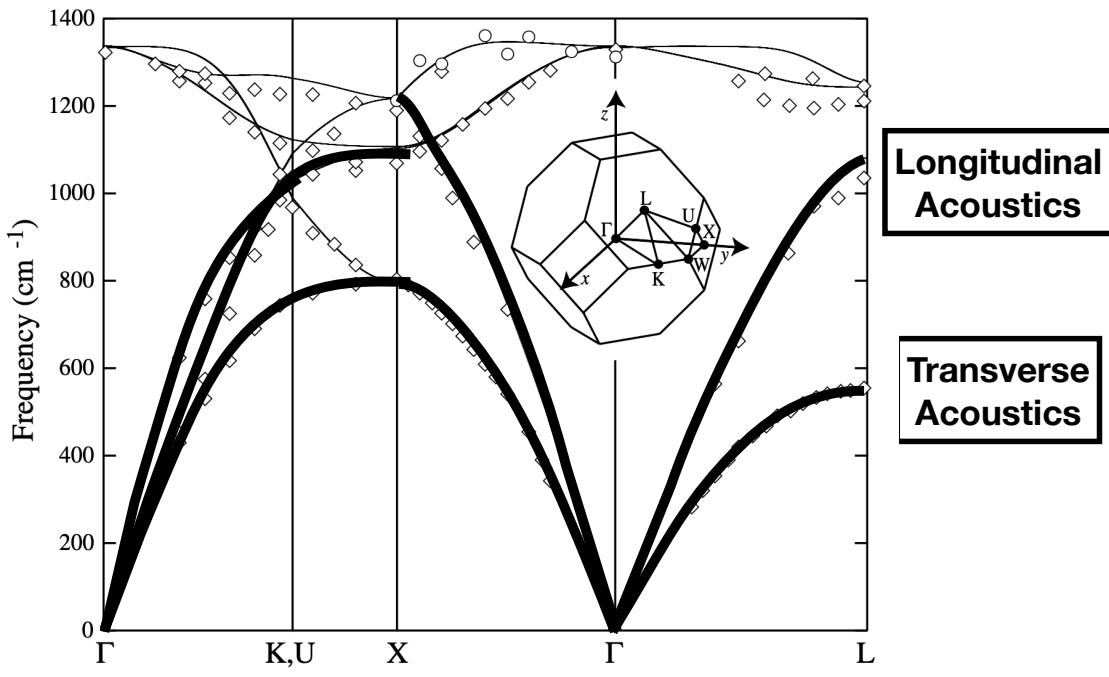


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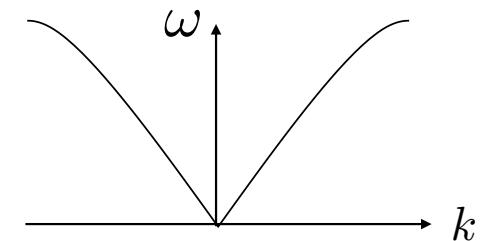
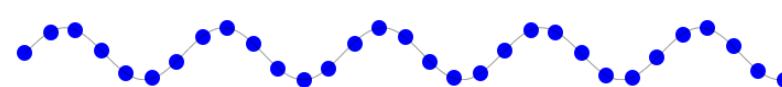
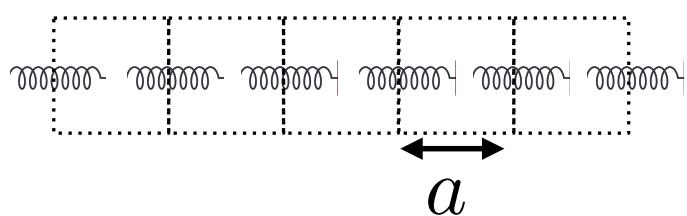
Measuring acoustic phonons with IXS

- \mathbf{Q} sets the direction of atomic displacements
- Polarization of atomic displacement wrt reduced \mathbf{q} can be changed by scanning different directions

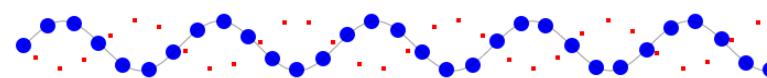
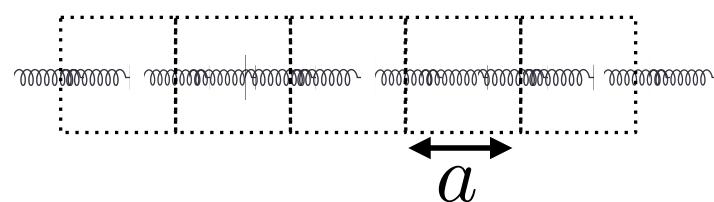


What about more complicated unit cells?

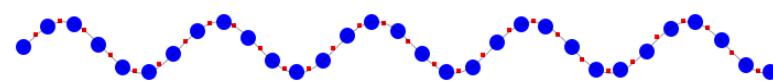
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- Diatomic 1D chain



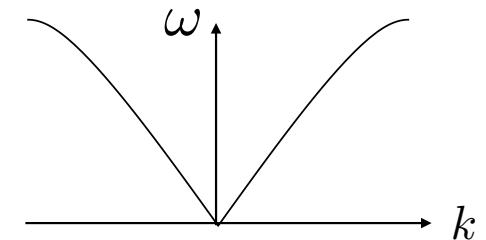
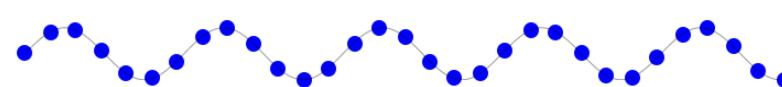
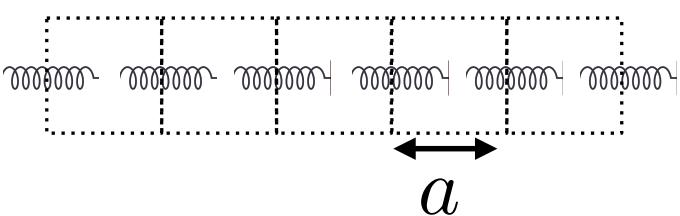
Red atom beats against blue atom



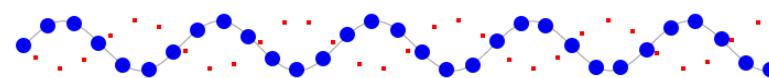
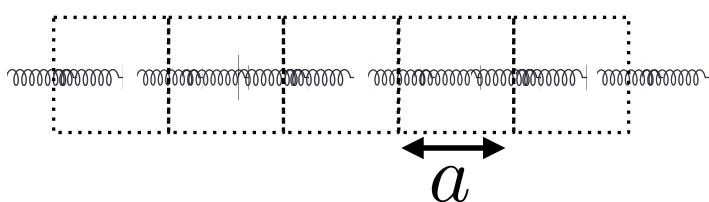
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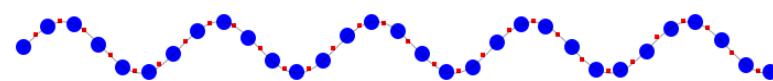
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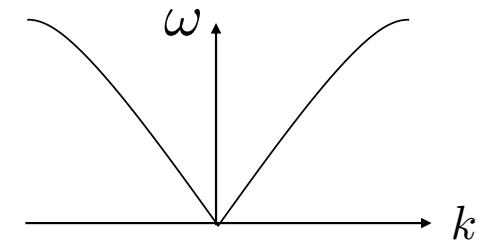
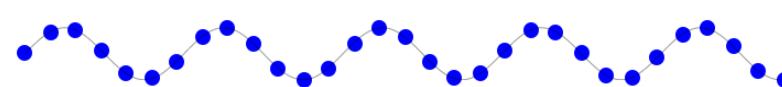
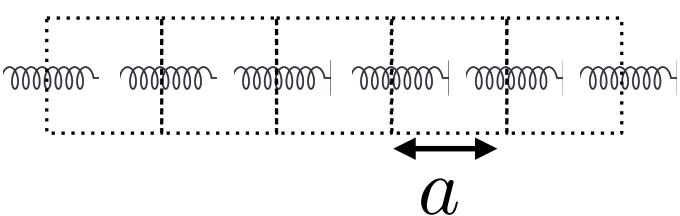
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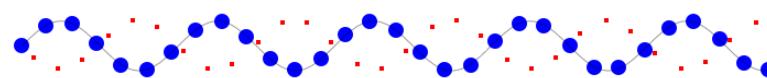
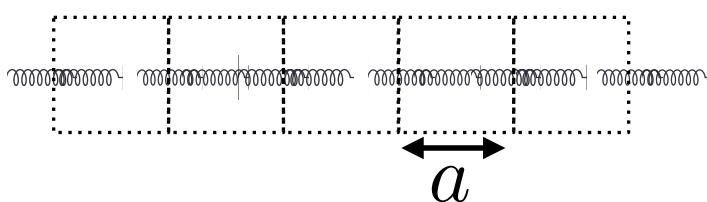
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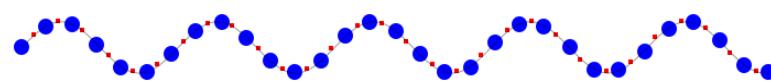
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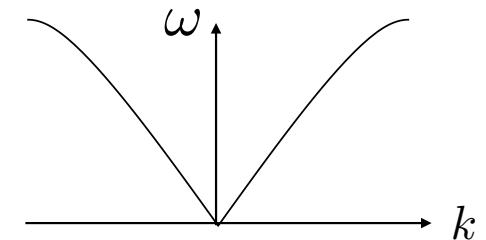
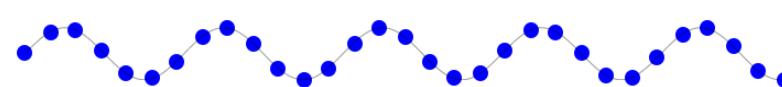
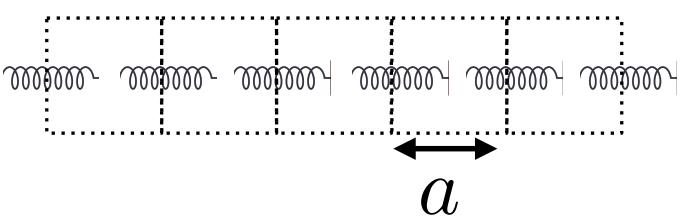
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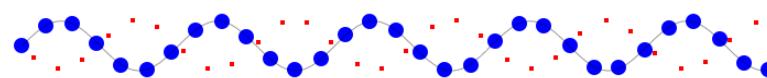
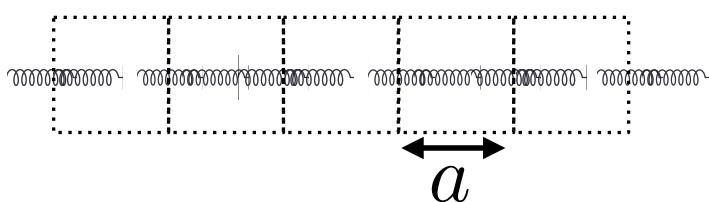
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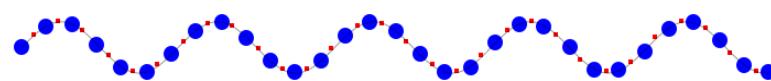
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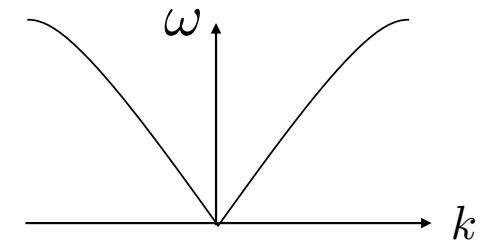
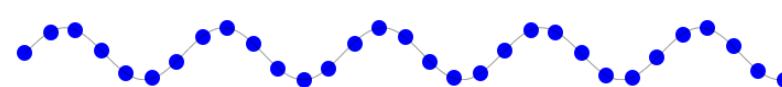
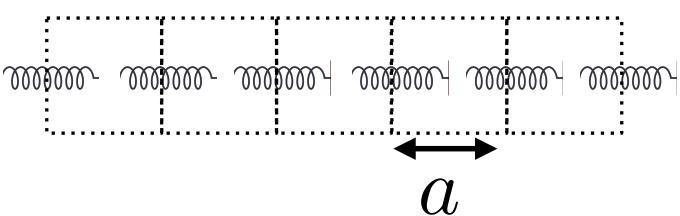
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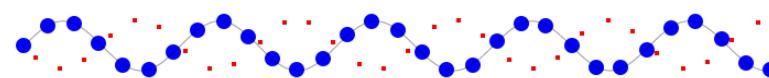
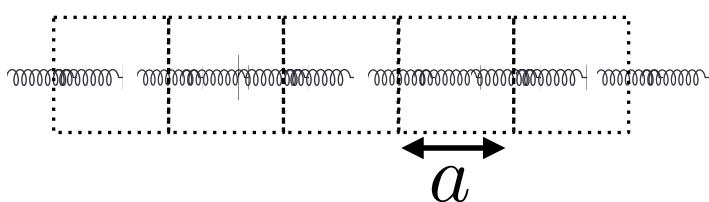
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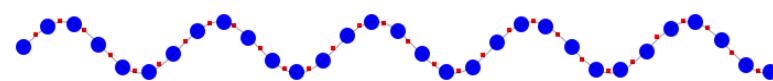
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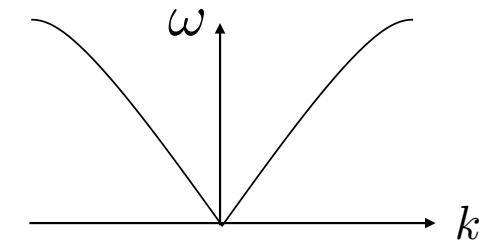
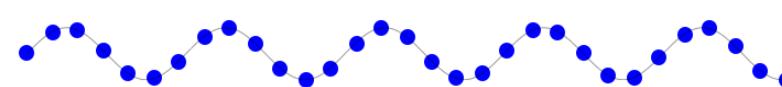
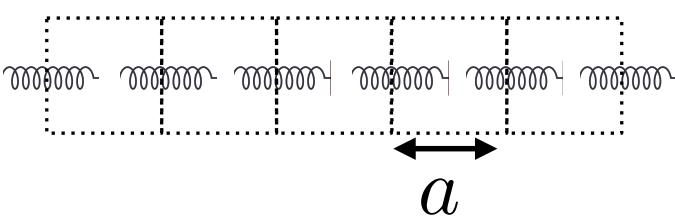
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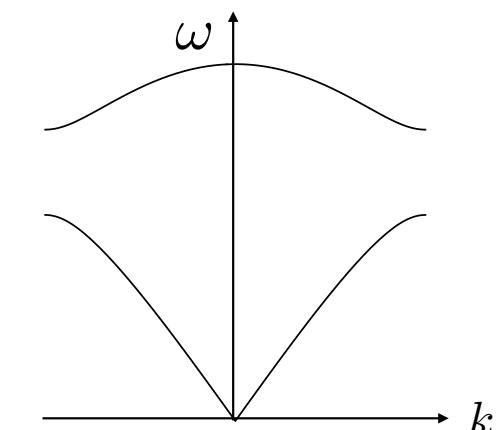
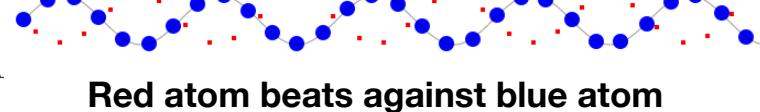
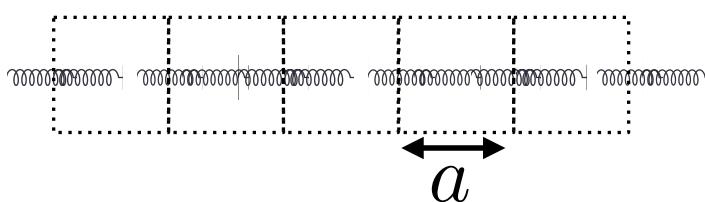
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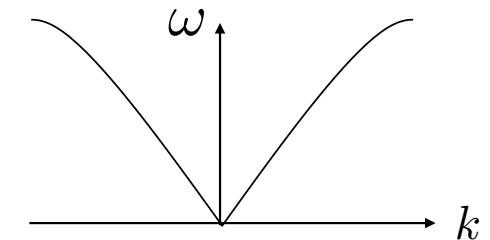
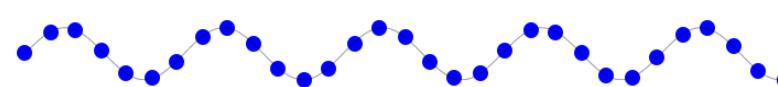
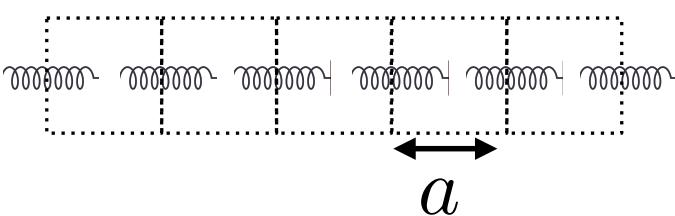


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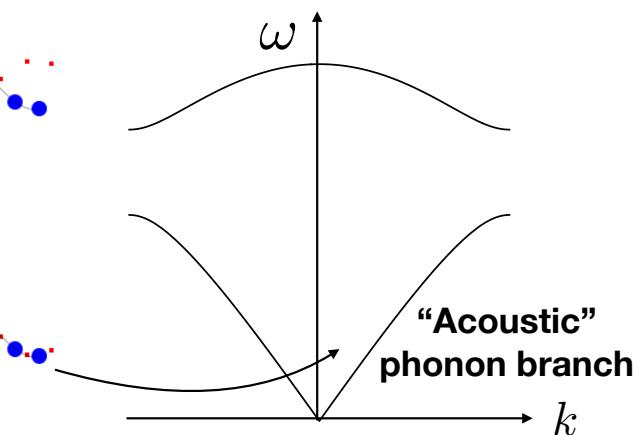
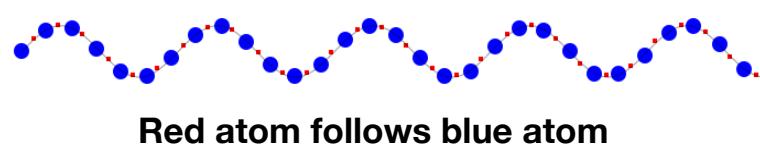
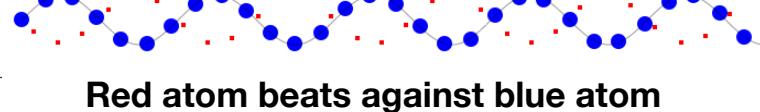
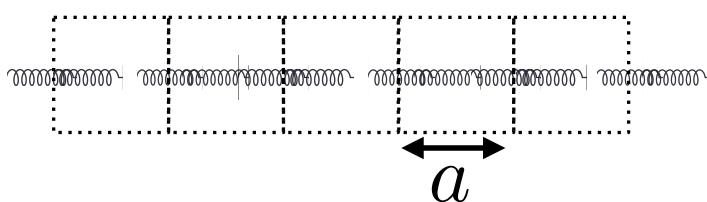


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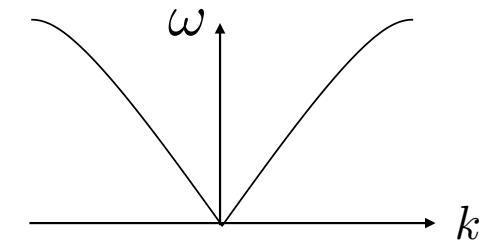
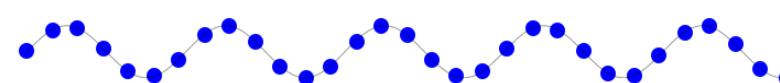
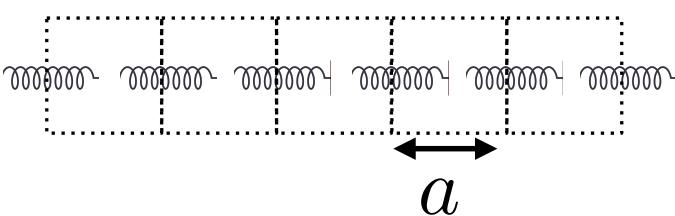


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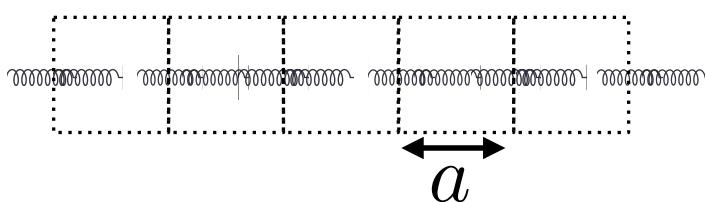


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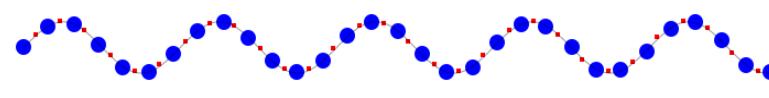
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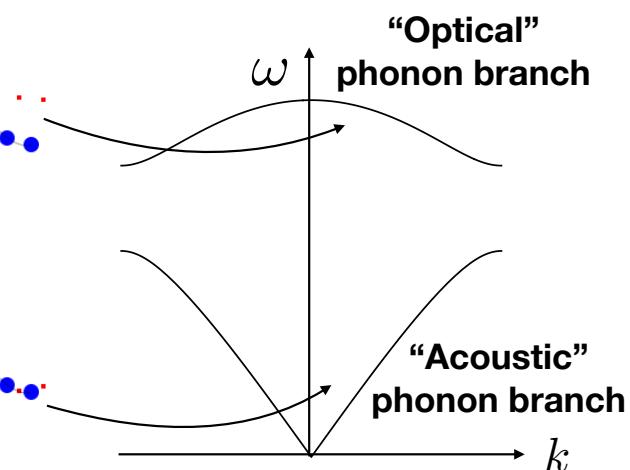
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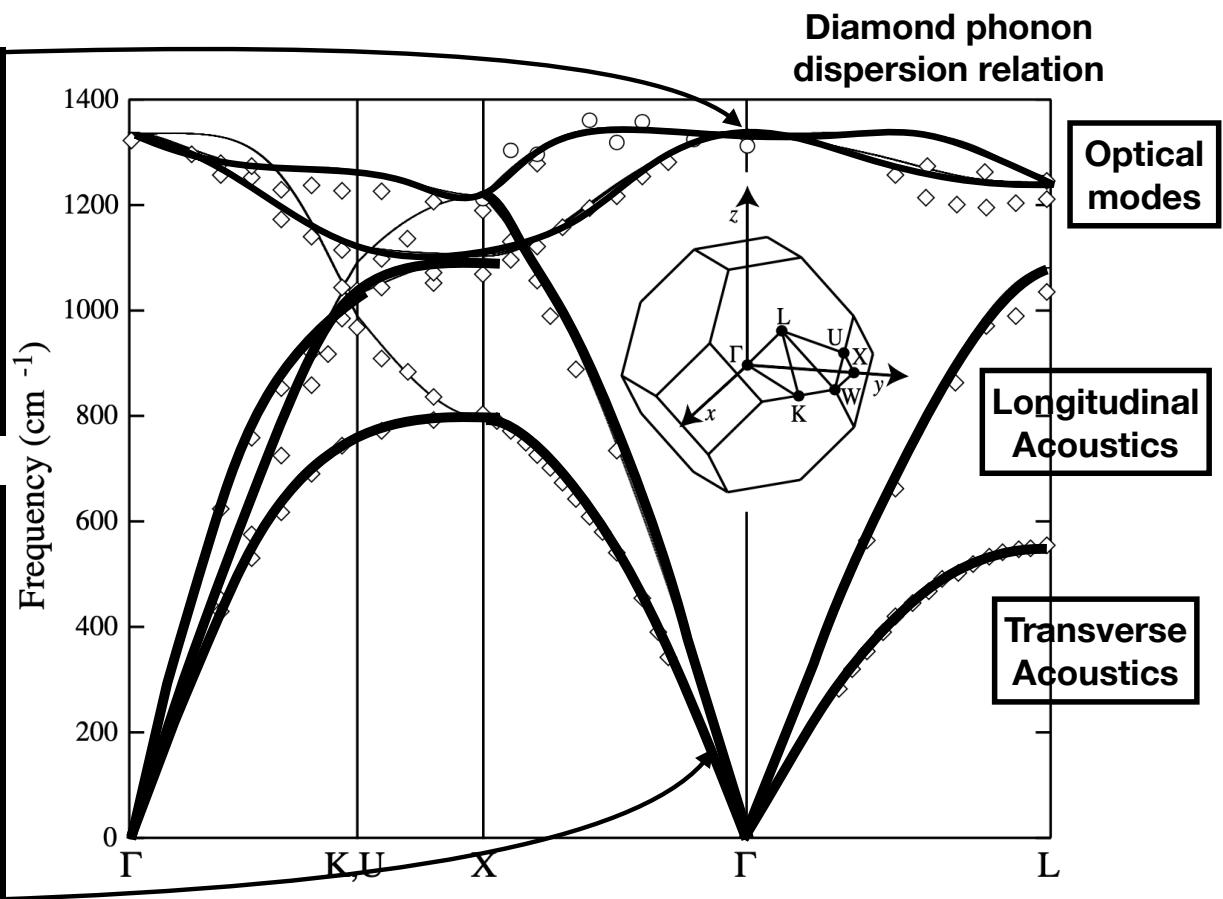
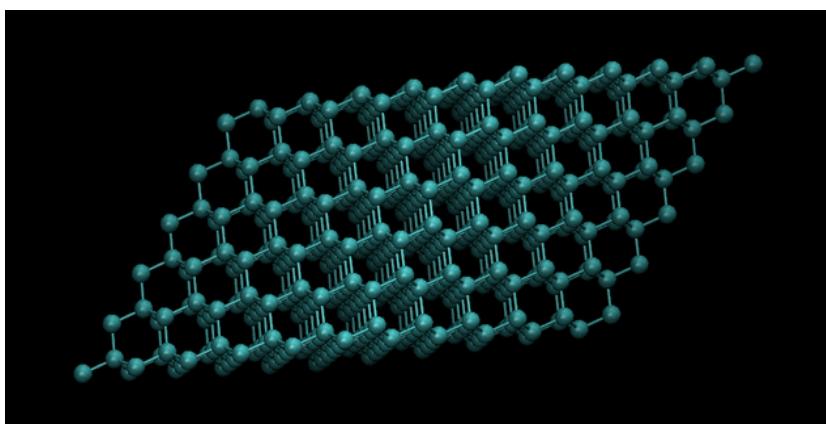
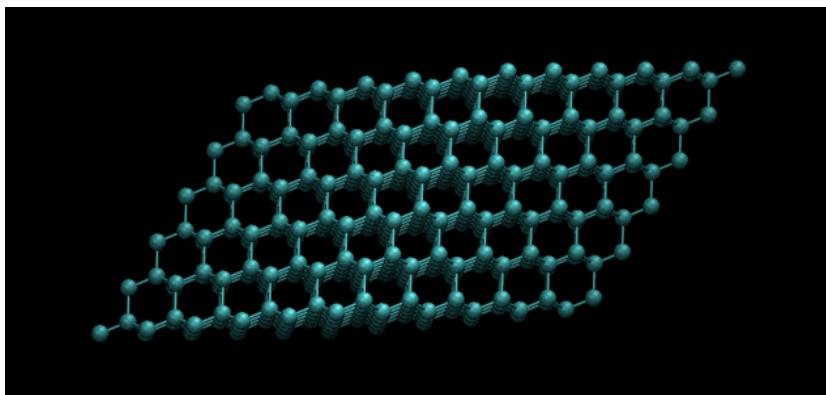
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Red atom follows blue atom

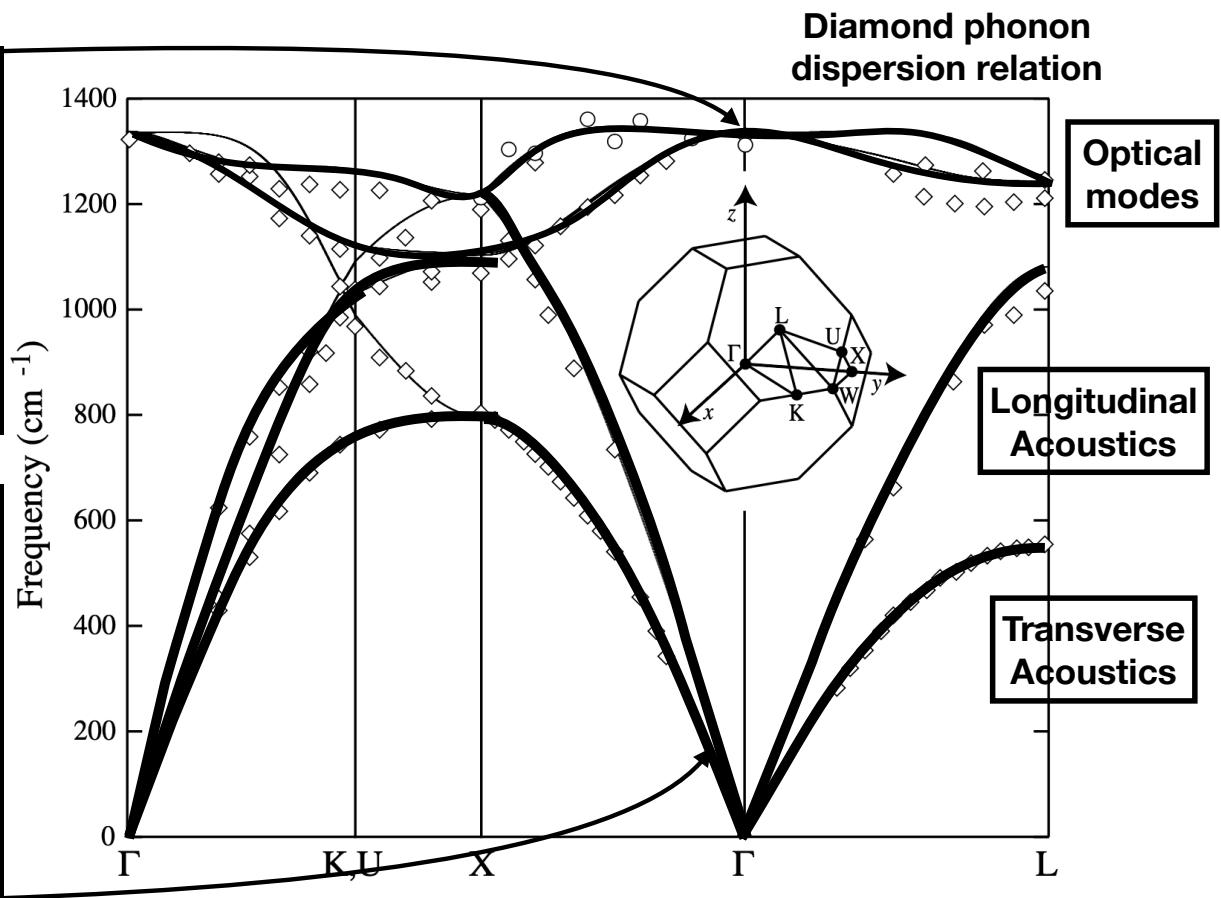
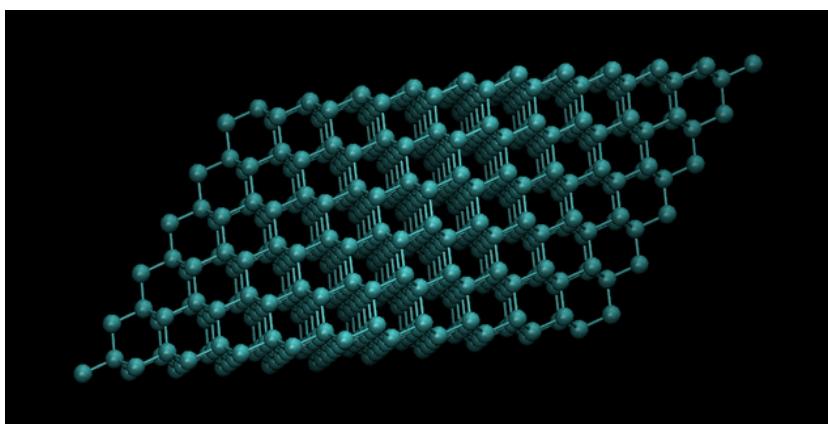
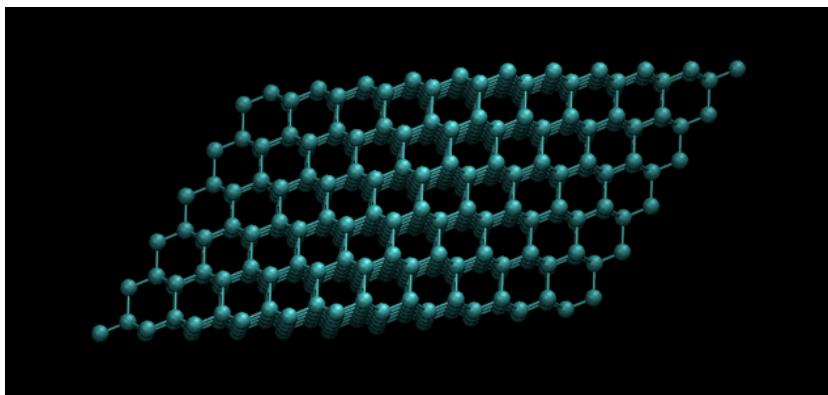


Phonons throughout the Brillouin zone



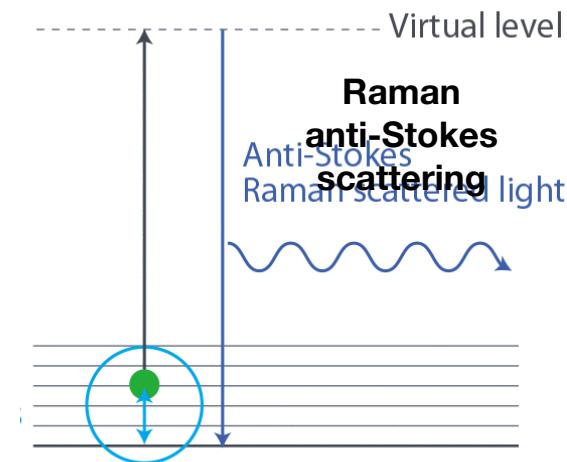
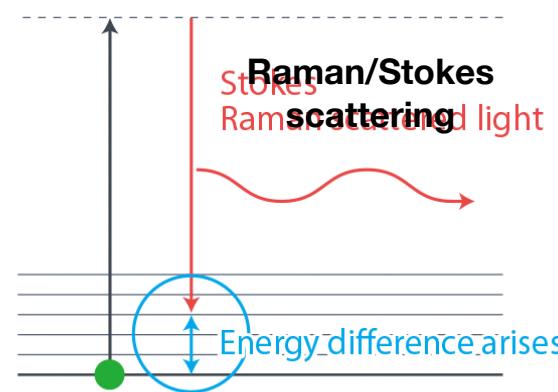
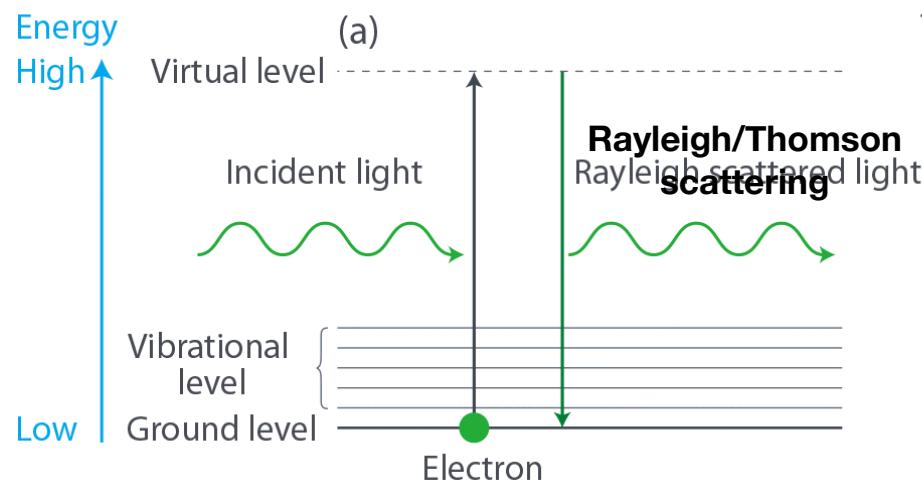
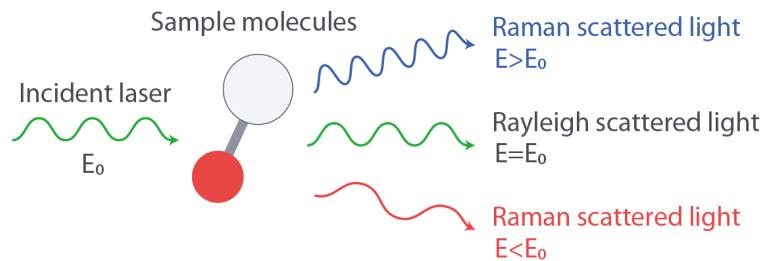
Gonze, et al, Z. Krist. 220 (2005) 458

Phonons throughout the Brillouin zone



Gonze, et al, Z. Krist. 220 (2005) 458

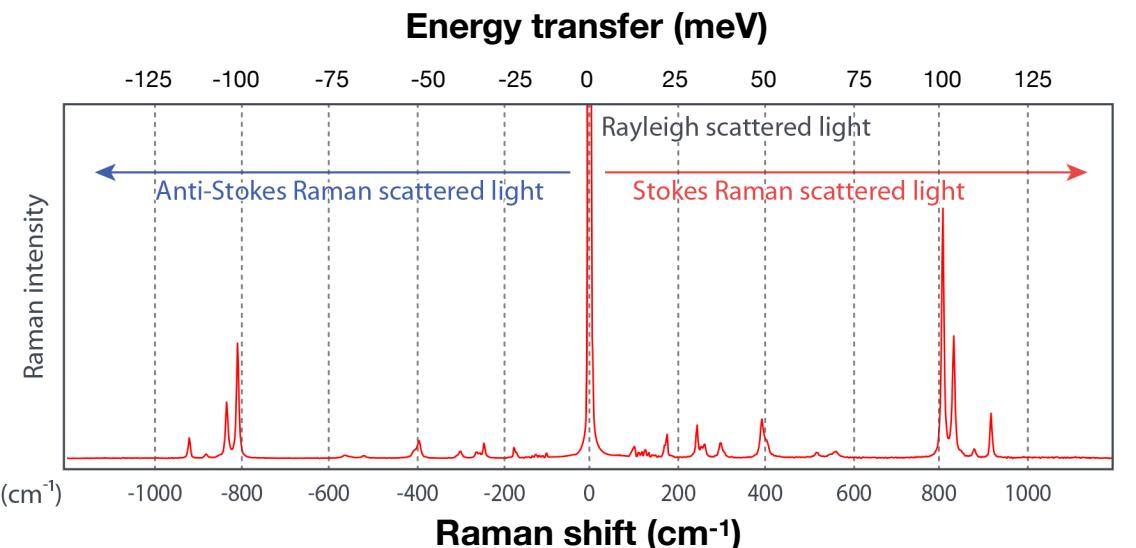
Raman=inelastic light scattering



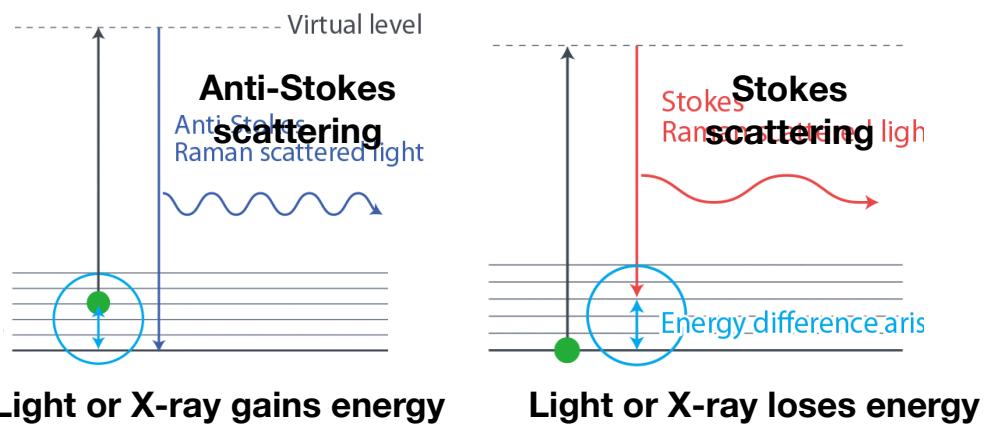
Discovered 1928
Nobel 1930

**Strong limitation:
momentum of light
is very small**

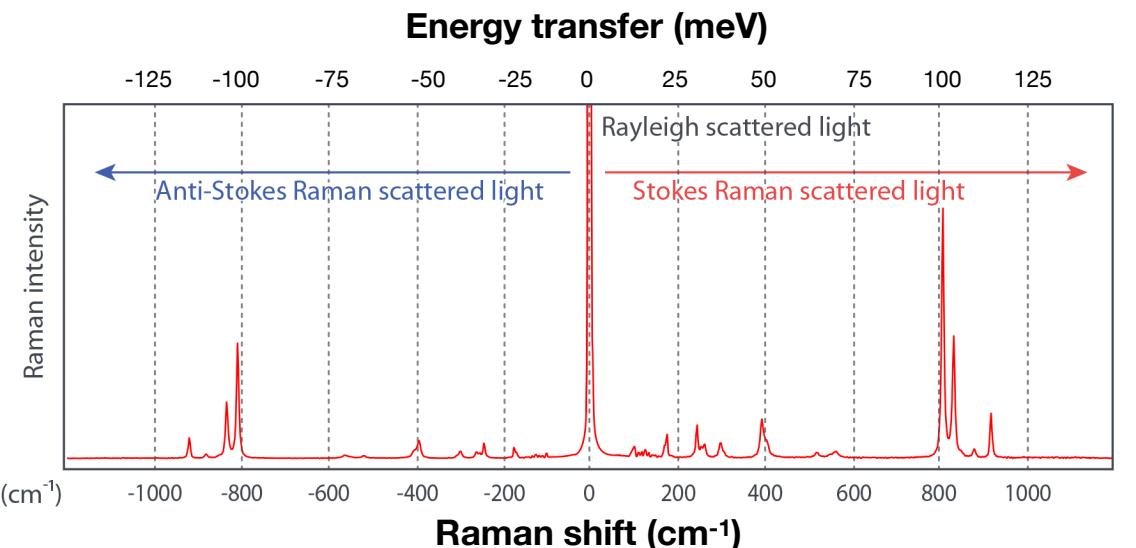
Energy transfer and Raman shift



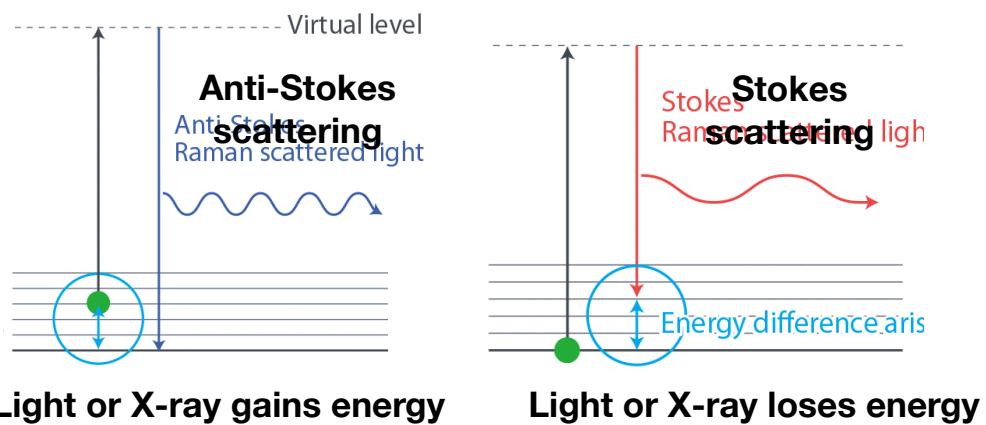
- Laser source=high resolution
- Lattice excitations easy to see
- Electronic and magnetic excitations possible



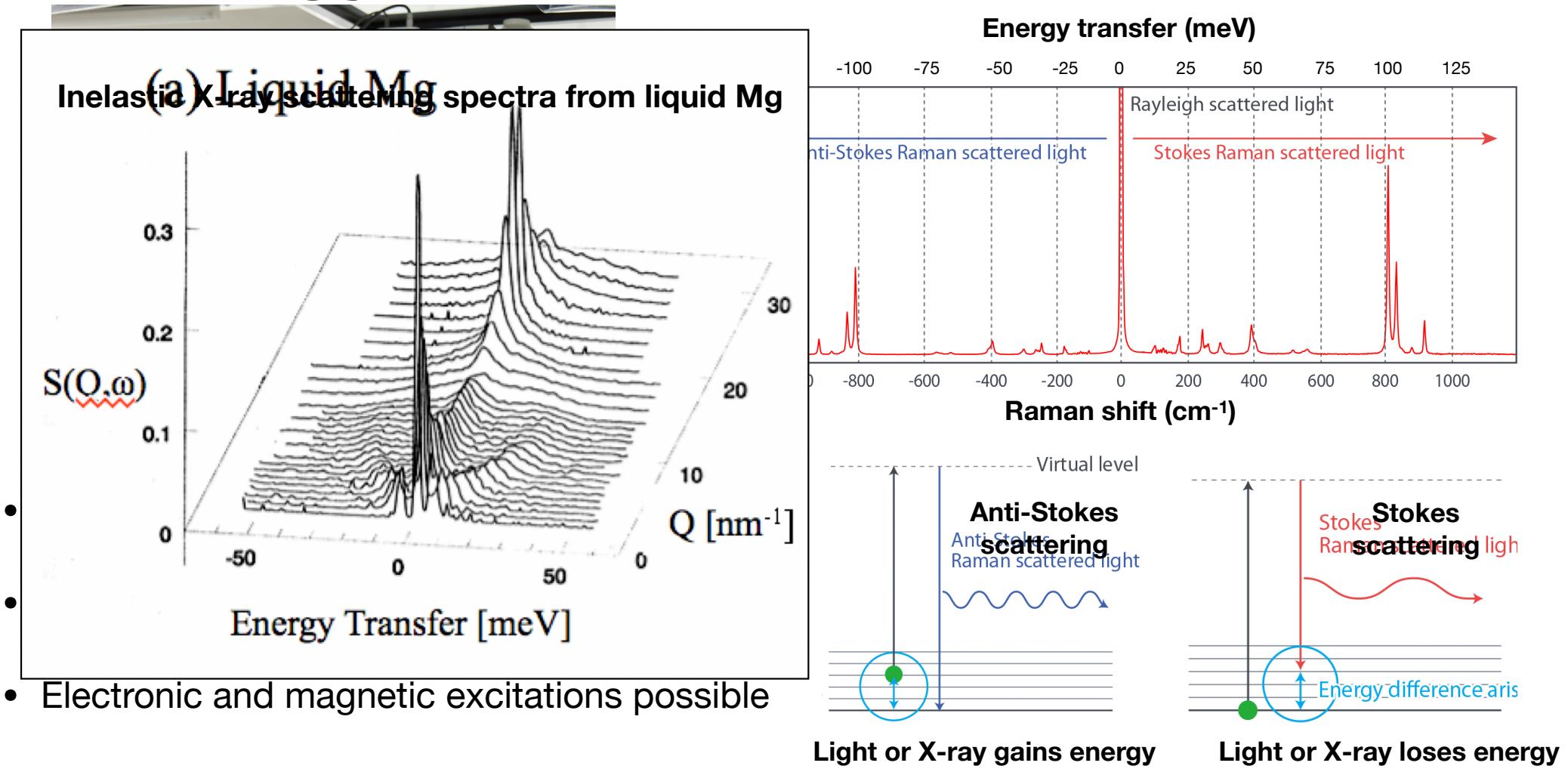
Energy transfer and Raman shift



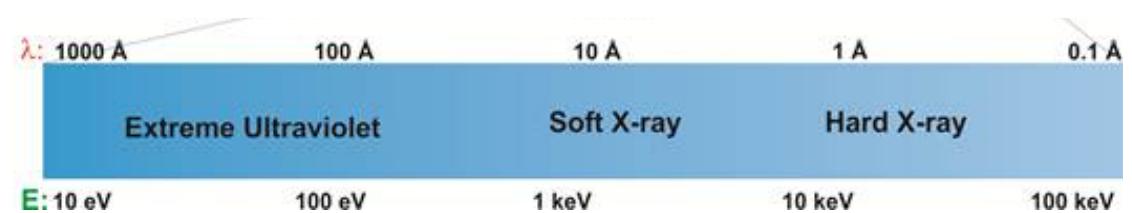
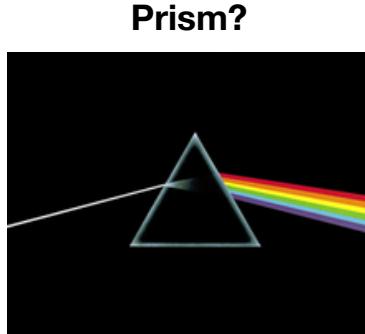
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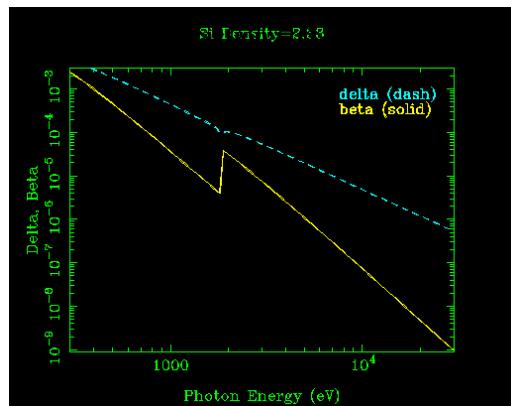
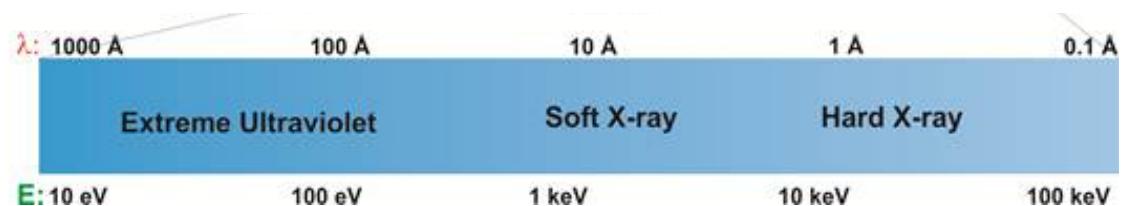
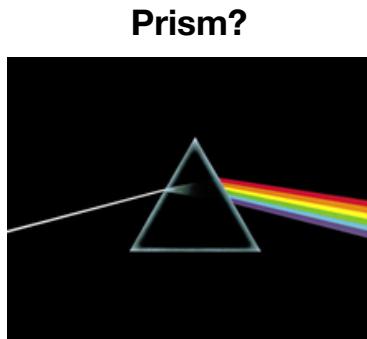
Energy transfer and Raman shift



How can we disperse X-rays according to energy/wavelength?



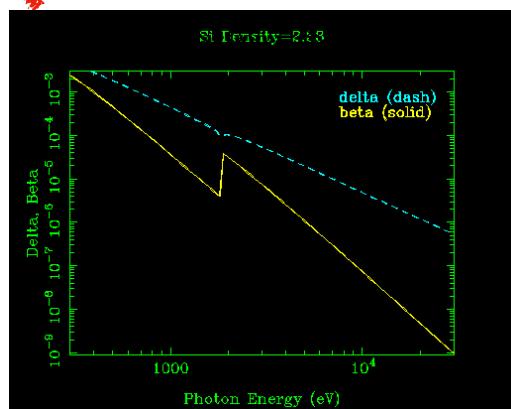
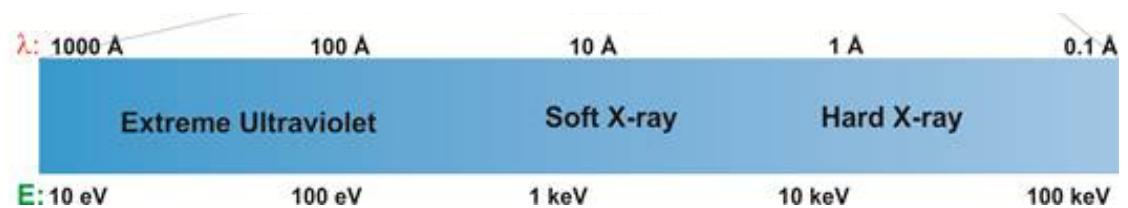
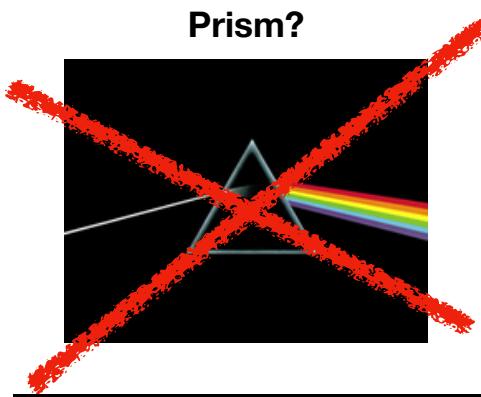
How can we disperse X-rays according to energy/wavelength?



Inefficient

Index changes too small to be effective

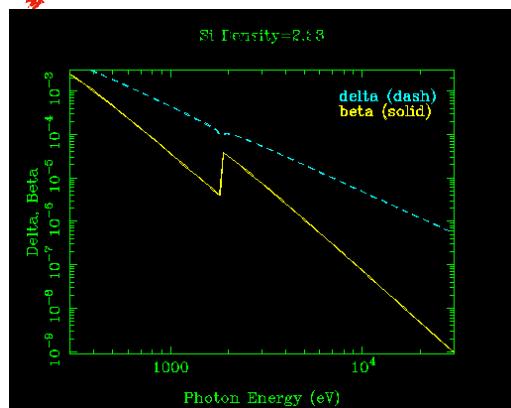
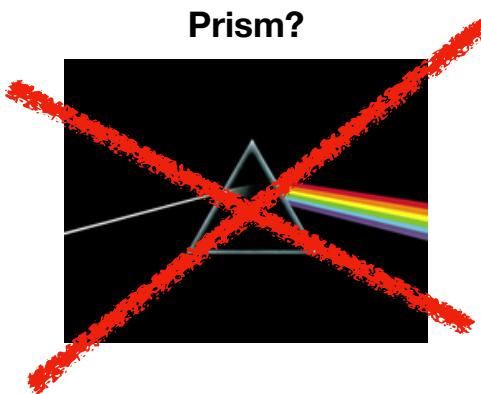
How can we disperse X-rays according to energy/wavelength?



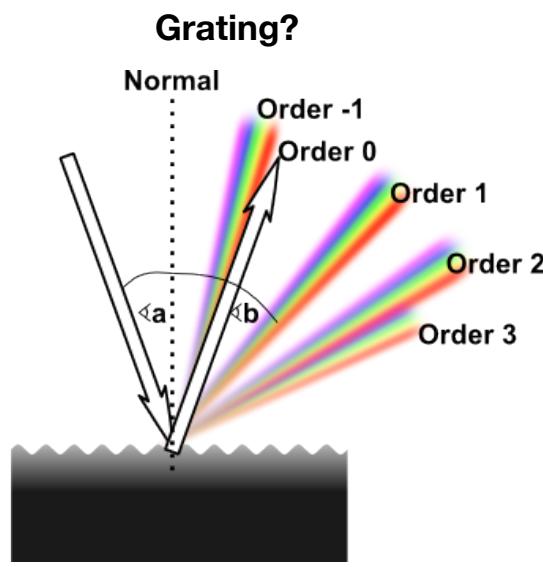
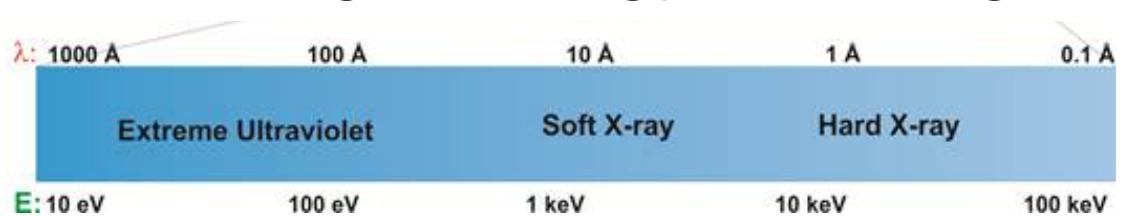
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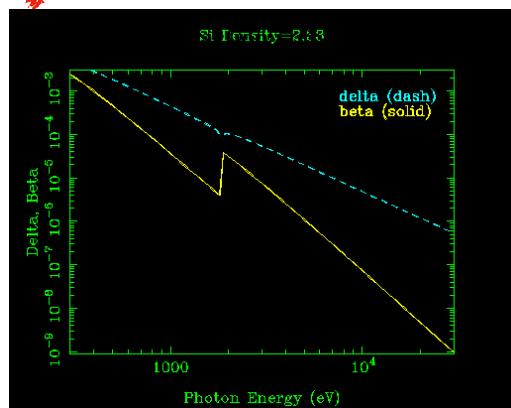
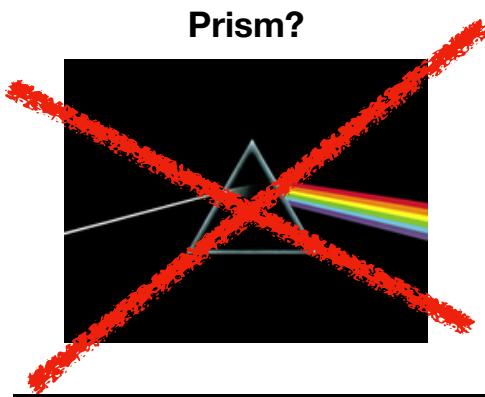
Inefficient
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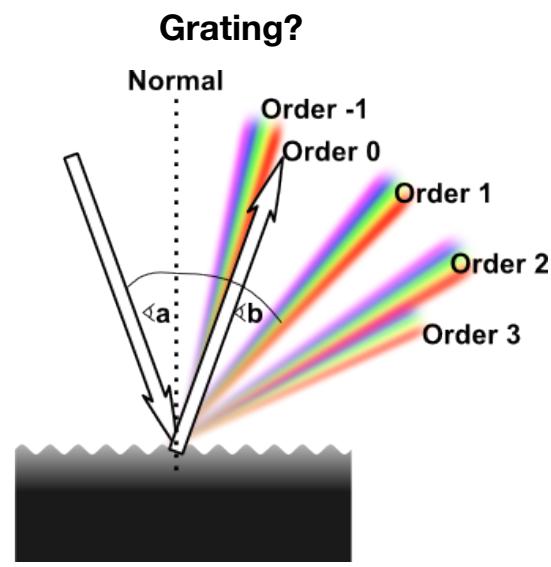
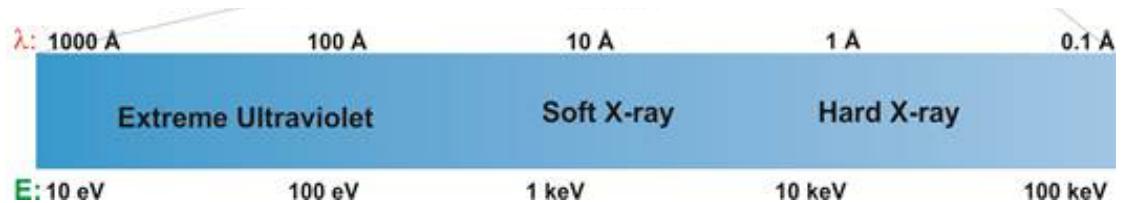
If feature sizes can
be on order wavelength

Good choice for lasers up to
soft X-ray ~ 1 nm

How can we disperse X-rays according to energy/wavelength?

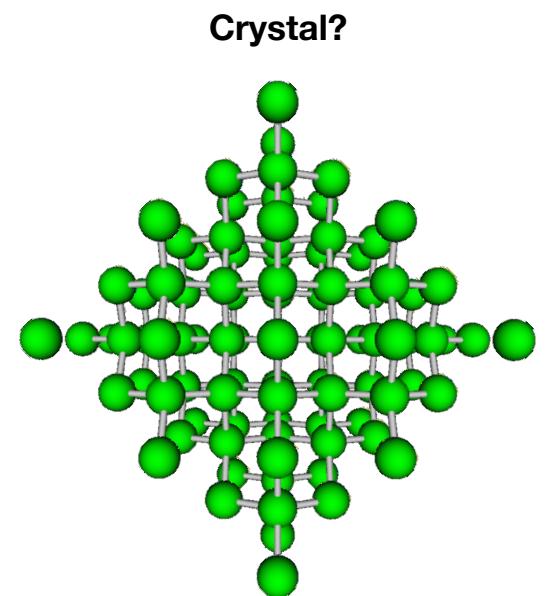


Inefficient
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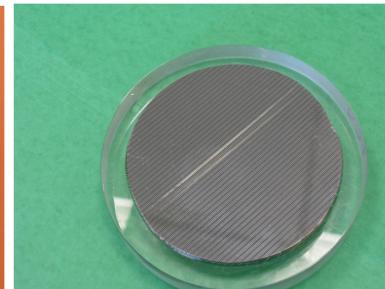


Hard X-ray << 1nm

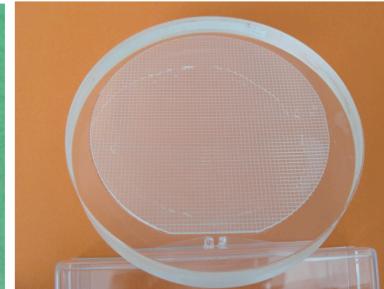
Diffracting X-rays crystals



Si



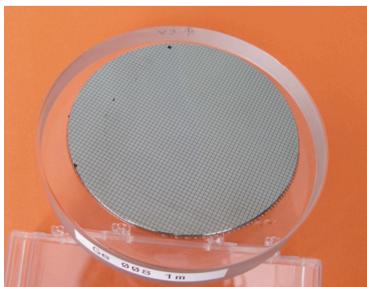
Ge



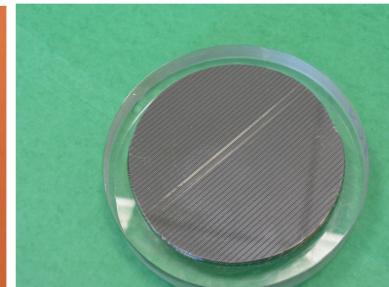
LiNbO₃

- Collecting inelastically scattered photons requires high quality, large, diced, bent crystal analyzers
- APS is a world leader in analyzer development

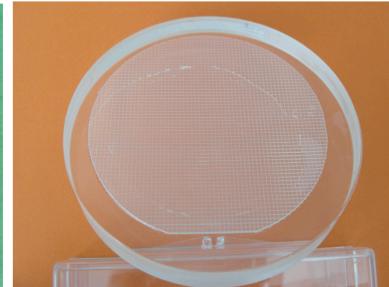
Diffracting X-rays crystals



Si

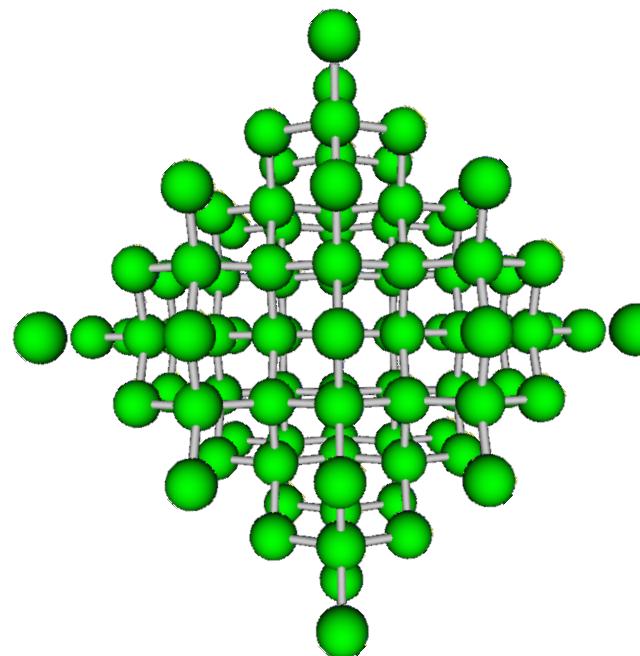


Ge



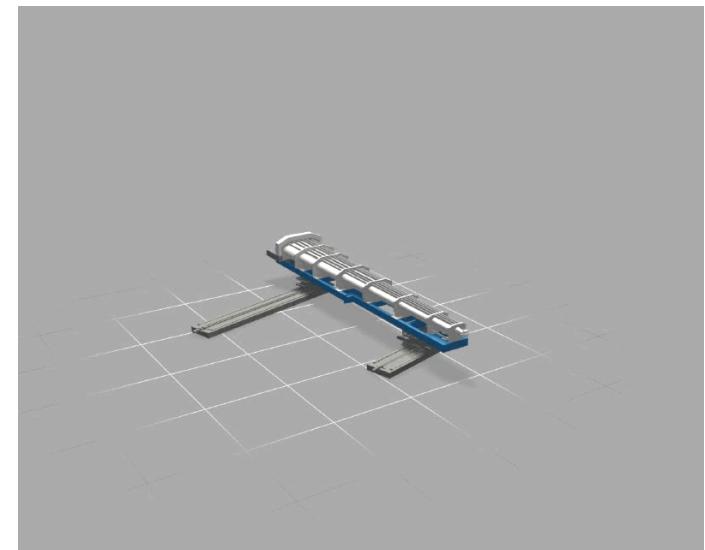
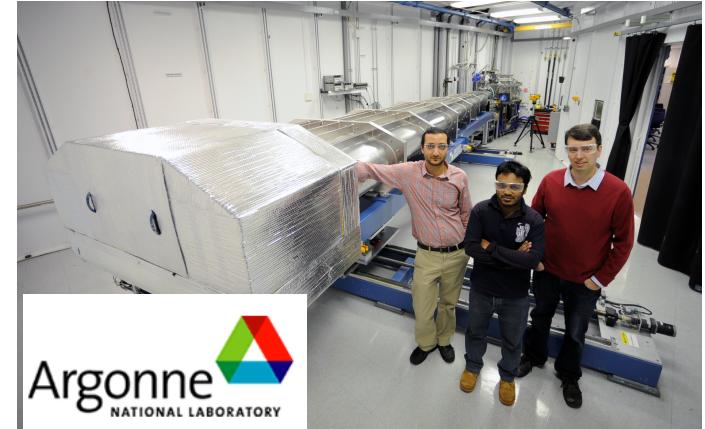
LiNbO_3

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- APS is a world leader in analyzer development



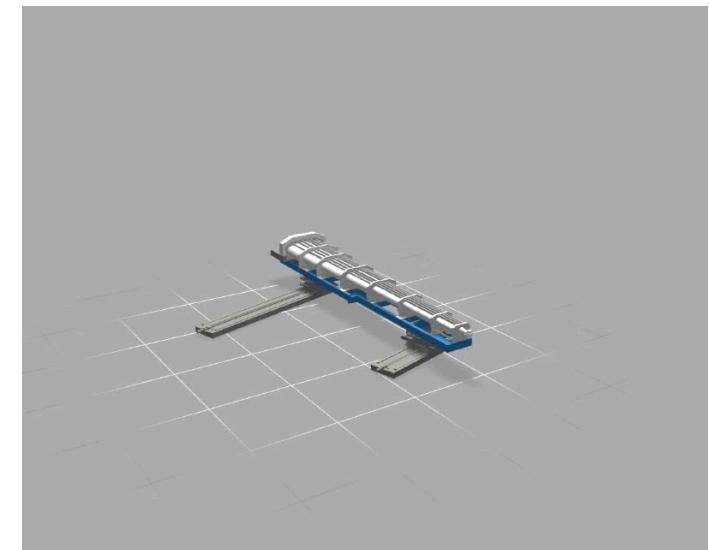
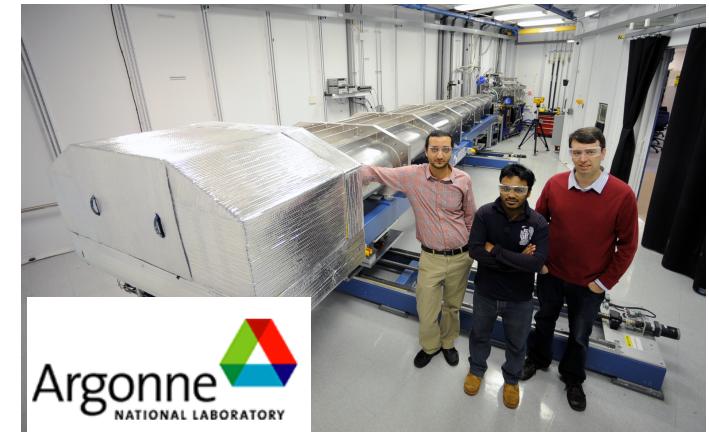
High Energy Resolution Inelastic X-ray Scattering

- Sector 30, Advanced Photon Source, Argonne National Lab
- 23,724 eV incident energy
- <1 meV incident bandwidth
- Resolving power $E_i/\Delta E_i = 2 \times 10^7$
- 9 analyzers sample, 9 momenta transfers simultaneously, 9m arm
- 1.5 meV energy resolution
- $20\mu\text{m} \times 5\mu\text{m}$ spot size
- Measures energy and momentum distribution of lattice vibrations $S(\vec{q}, \omega)$



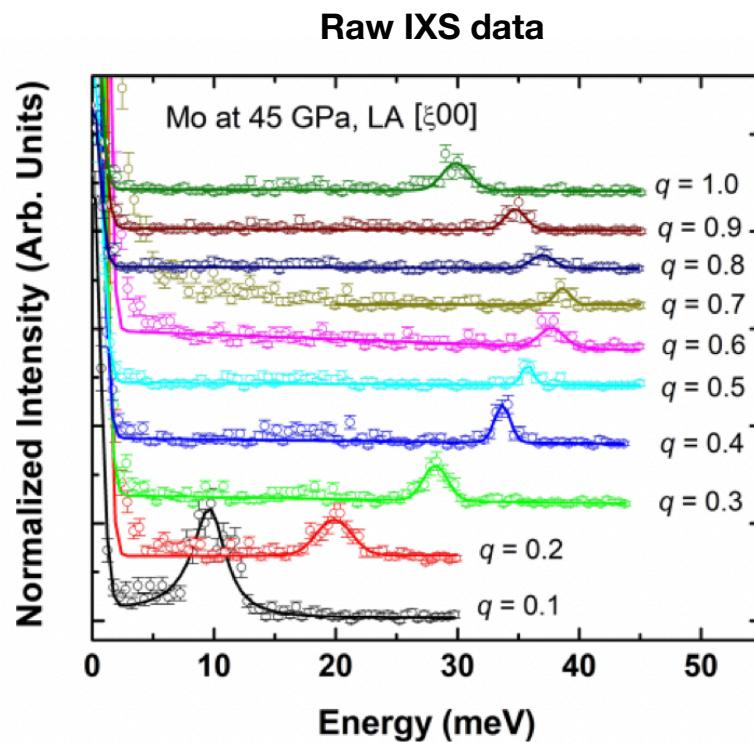
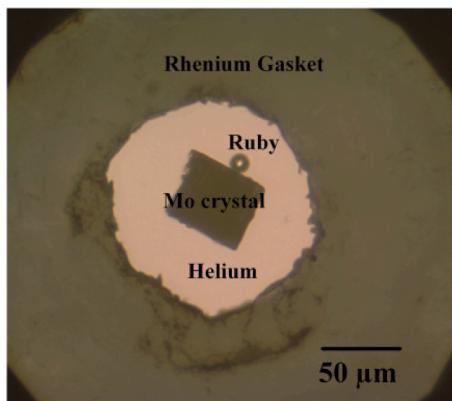
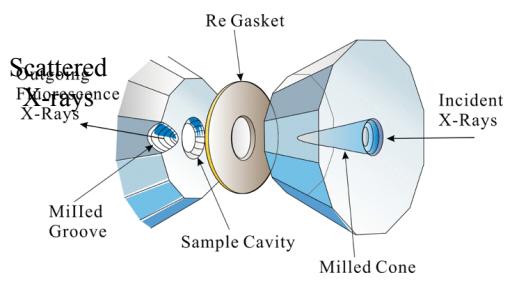
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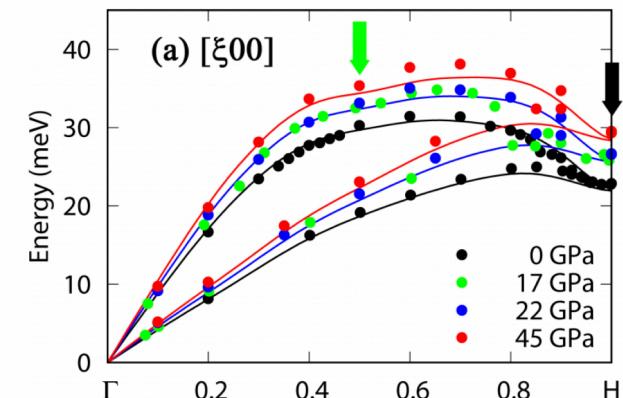


Kohn anomaly and elastic softening in body-centered cubic molybdenum at high pressure

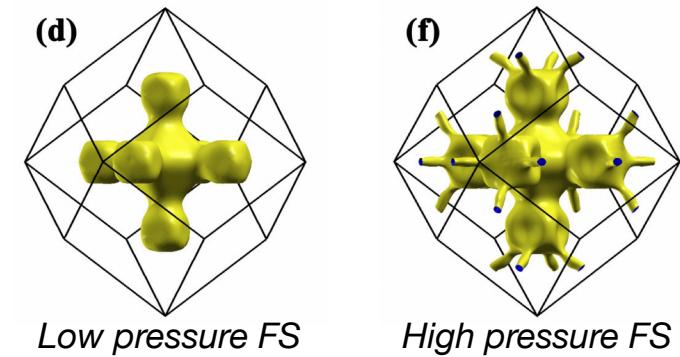
Sample and environment



Pressure-dependent dispersion

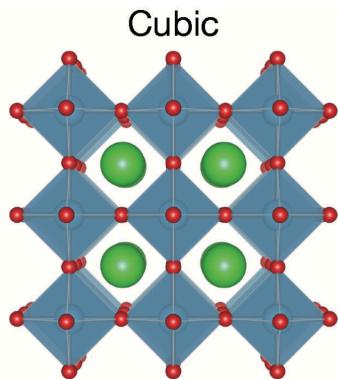


Interpretation - fermi surface changes at high pressure

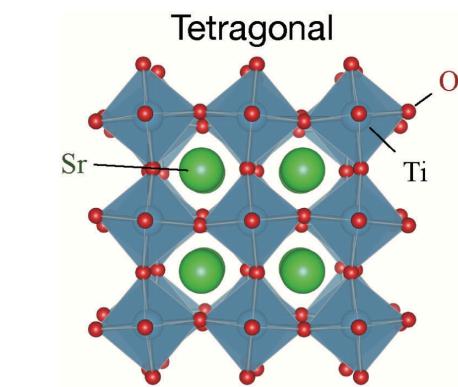
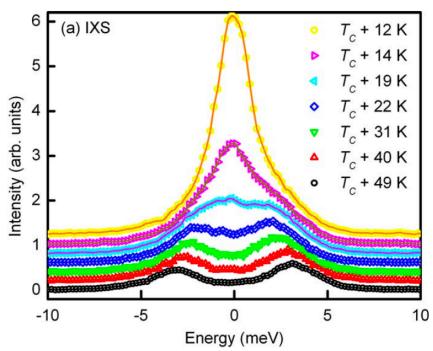


Central peak and narrow component in x-ray scattering measurements near the displacive phase transition in SrTiO_3

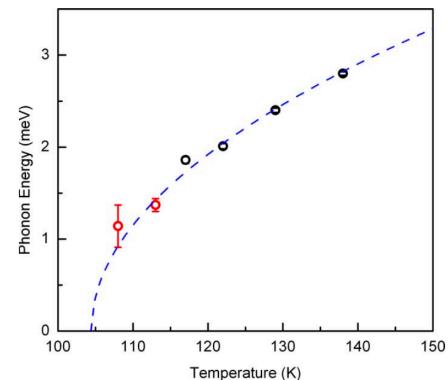
- Structural transitions can be described as the “freezing” of a phonon mode



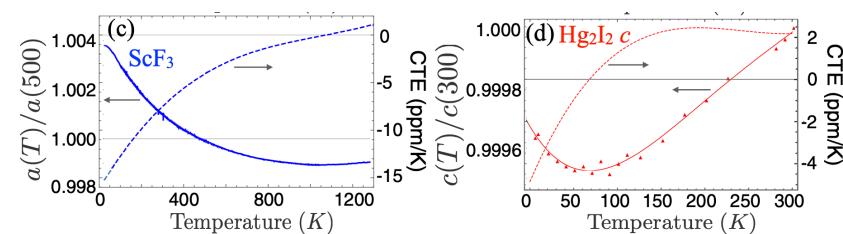
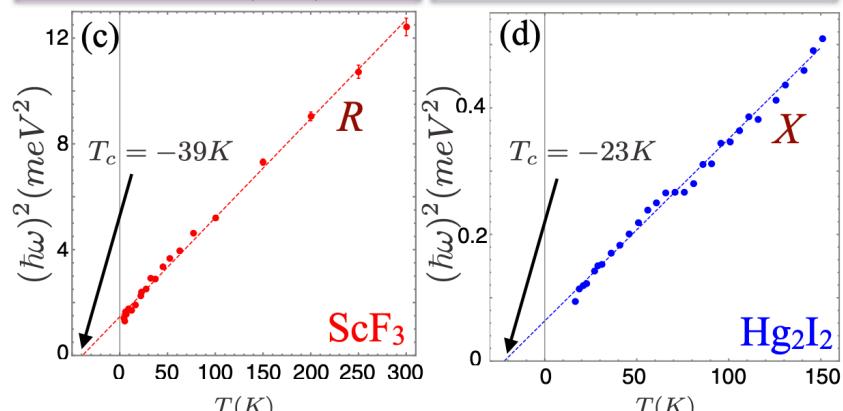
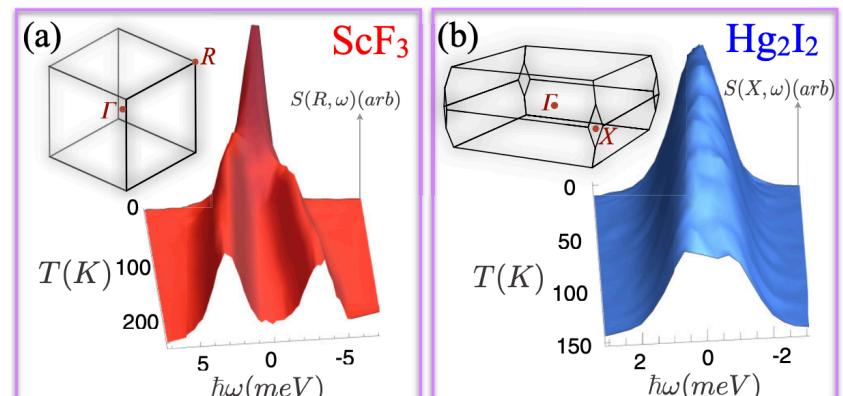
High temperature structure



Low temperature structure



Negative thermal expansion near two structural quantum phase transitions



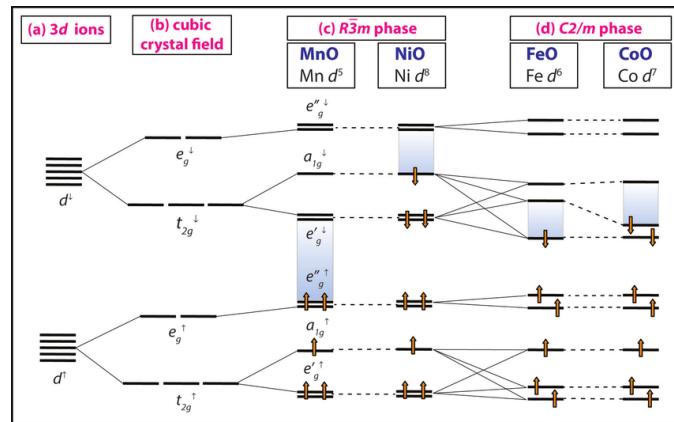
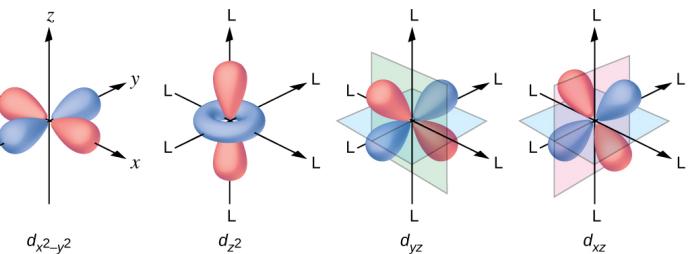
Non-phonon IXS-active excitations

PRL 99, 026401 (2007)

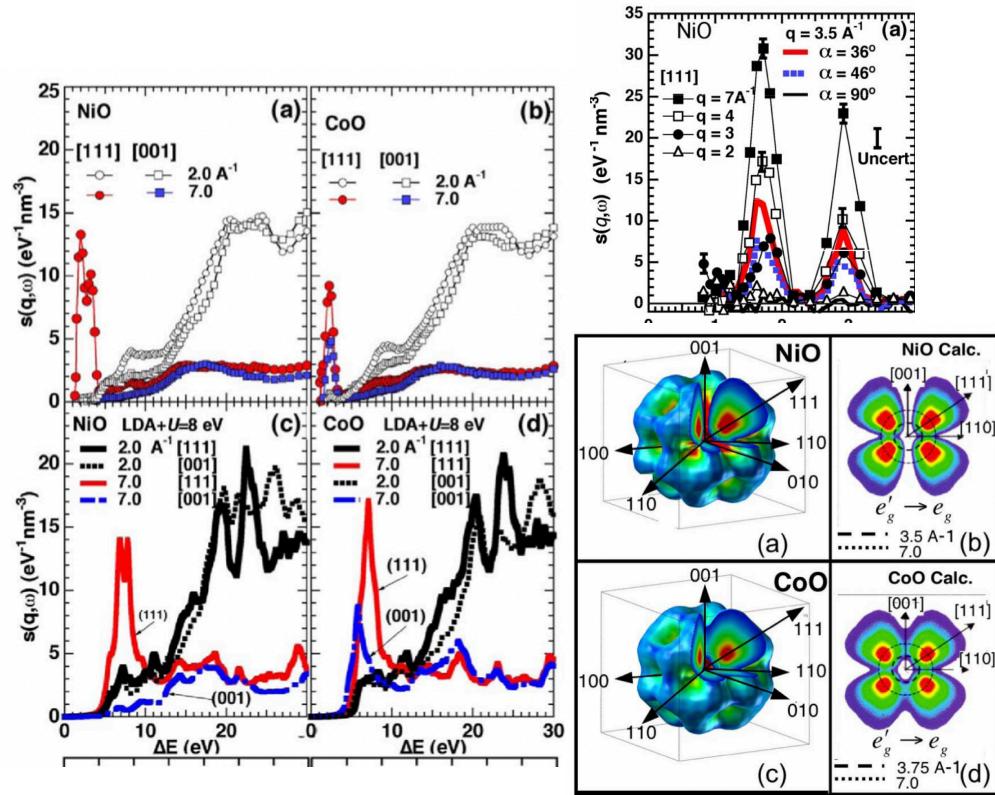
PHYSICAL REVIEW LETTERS

week ending
13 JULY 2007

- Different d electron orbitals have different energy when placed in a crystal
- Gives rise to “crystal field excitations”



Nonresonant Inelastic X-Ray Scattering and Energy-Resolved Wannier Function Investigation of *d-d* Excitations in NiO and CoO



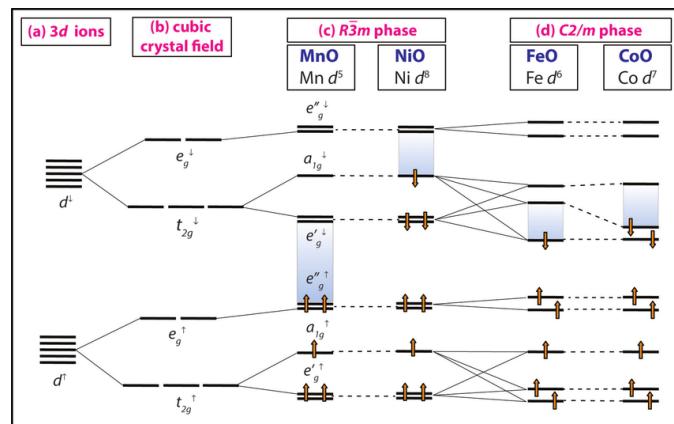
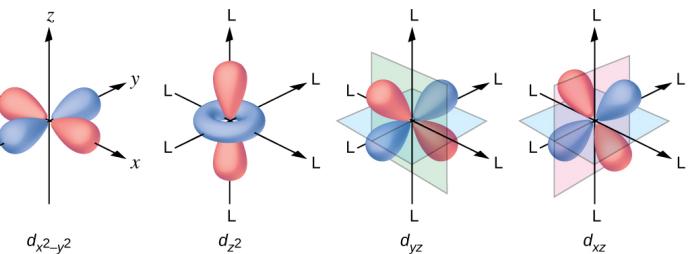
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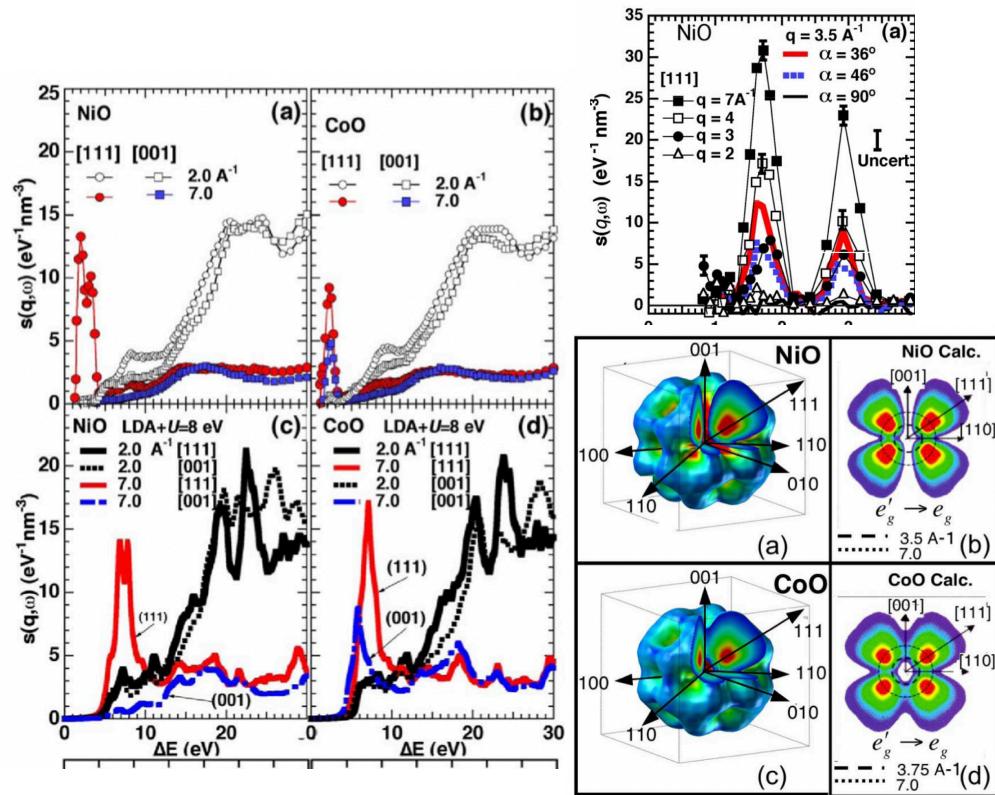
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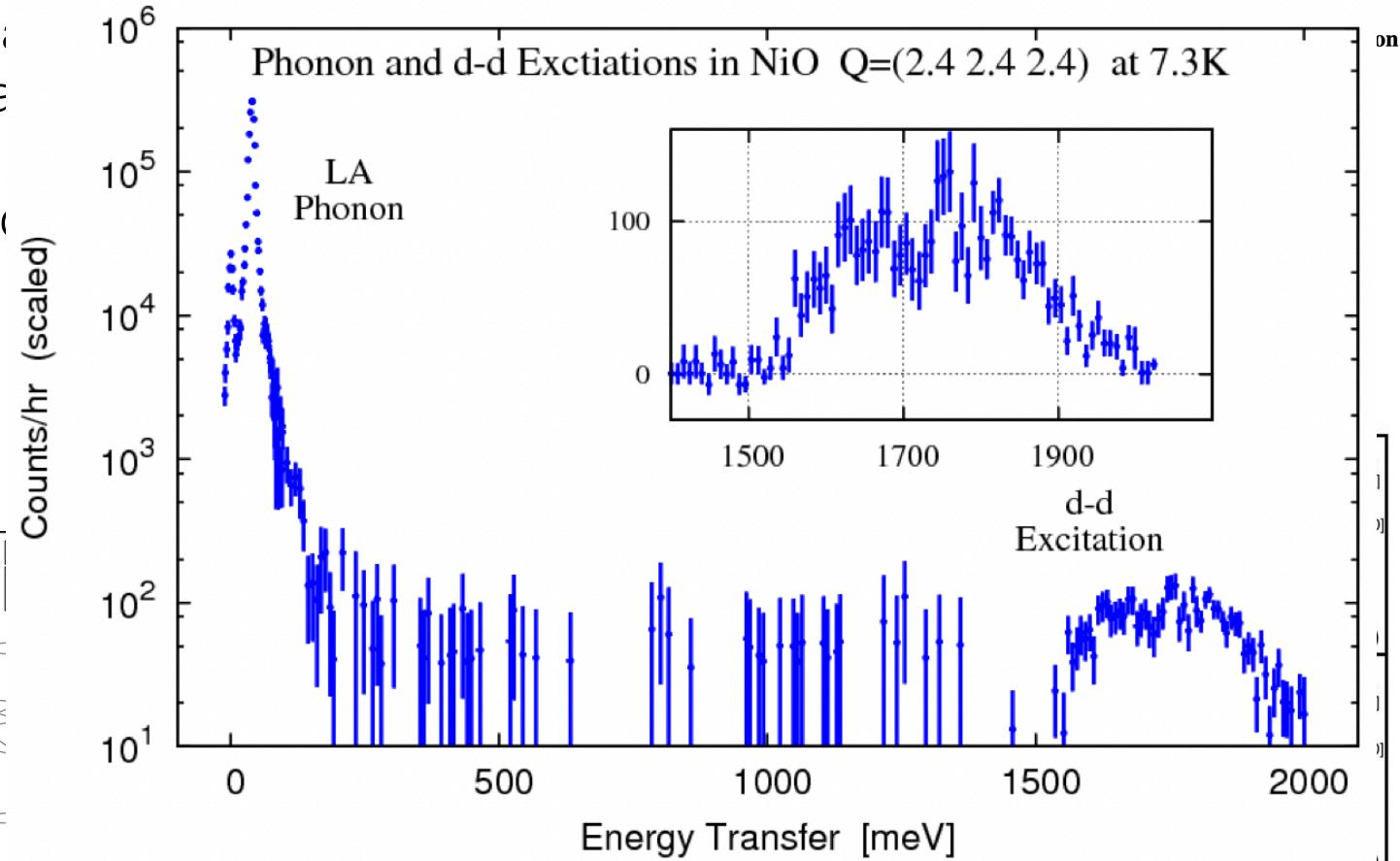
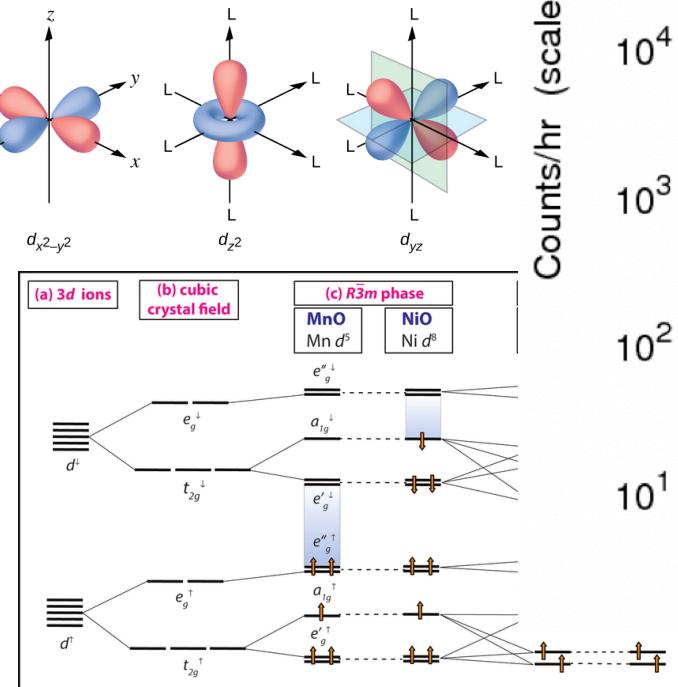
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PHYSICAL REVIEW LETTERS

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arXiv:1504.01098v7

Thank you!

Feedback

Lecture – 8:30 – 9:30

Inelastic X-ray Scattering - Jason Hancock

<https://forms.office.com/g/mXiC92R1B4>

